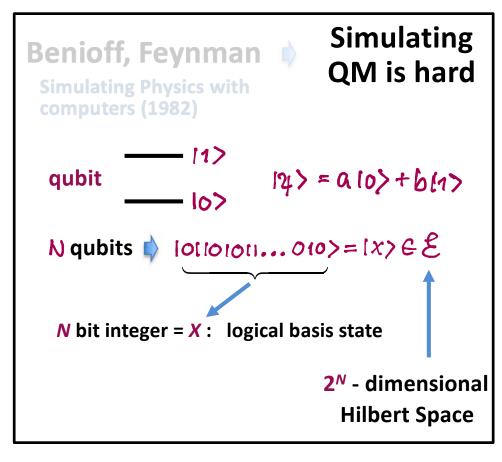
Quantum Complexity



General State:

$$|z|_{x=0} = \sum_{x=0}^{2^{N-1}} a_x |x\rangle$$

Simulating Physics with Computers

Richard P. Feynman

Department of Physics, California Institute of Technology, Pasadena, California 91107

Received May 7, 1981

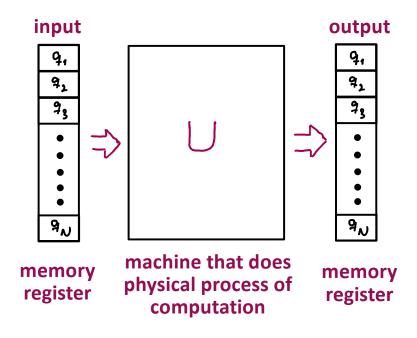
1. INTRODUCTION

On the program it says this is a keynote speech—and I don't know what a keynote speech is. I do not intend in any way to suggest what should be in this meeting as a keynote of the subjects or anything like that. I have my own things to say and to talk about and there's no implication that anybody needs to talk about the same thing or anything like it. So what I want to talk about is what Mike Dertouzos suggested that nobody would talk about. I want to talk about the problem of simulating physics with computers and I mean that in a specific way which I am going to explain. The reason for doing this is something that I learned about from Ed Fredkin, and my entire interest in the subject has been inspired by him. It has to do with learning something about the possibilities of computers, and also something about possibilities in physics. If we suppose that we know all the physical laws perfectly, of course we don't have to pay any attention to computers. It's interesting anyway to entertain oneself with the idea that we've got something to learn about physical laws; and if I take a relaxed view here (after all I'm here and not at home) I'll admit that we don't understand everything.

The first question is, What kind of computer are we going to use to simulate physics? Computer theory has been developed to a point where it realizes that it doesn't make any difference; when you get to a *universal computer*, it doesn't matter how it's manufactured, how it's actually made. Therefore my question is, Can physics be simulated by a universal computer? I would like to have the elements of this computer *locally interconnected*, and therefore sort of think about cellular automata as an example (but I don't want to force it). But I do want something involved with the

OK – Plausible QM can do more Where does the QC's power come from?

Visualization of Computation



Classical: Register is in one of the logical states

Reversible transformation

Quantum: Register can be in any coherent superposition of logical states [x>

Unitary transformation U: I×>→ ly>

Maps basis to basis $U: \{ |x \rangle \} \rightarrow \{ |y \rangle \}$

Quantum Parallelism

$$|\mathcal{A}_{in}\rangle \rightarrow \sum_{x} a_{x}|_{x}\rangle \rightarrow |_{x}$$

$$\rightarrow |\mathcal{A}_{out}\rangle = \bigcup |\mathcal{A}_{in}\rangle = \sum_{x} a_{x}|_{y}\rangle = \sum_{x} b_{x}|_{x}\rangle$$

Machine processes 2^N inputs "in parallel"!

Beware: measurement collapses Q. Register into a single basis state at random

We get one random result out of 2^N

Quantum: Register can be in any coherent superposition of logical states (x>

Unitary transformation U: <a>!५> <a>!५>

Maps basis to basis $U: \{ |x \rangle \} \rightarrow \{ |y \rangle \}$

Quantum Parallelism

Quantum Sampling

Problem

$$|\mathcal{Y}_{in}\rangle \rightarrow \sum_{x} a_{x}|_{x}\rangle \rightarrow |\mathcal{Y}_{in}\rangle = \sum_{x} a_{x}|_{\eta}\rangle = \sum_{x} b_{x}|_{x}\rangle$$

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Quantum Algorithms look for global properties of functions — symmetry, periodicity, etc.

- * Classical -> requires many function evaluations
- Quantum -> design U so measurement gives answer with high probability
- * \exists classes of problems (sampling problems)

 which are classically hard but quantum "easy"

 Google "Quantum Supremacy"

Expert insight into current research

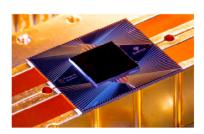
News & views

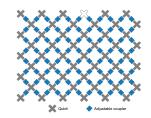
Quantum information

Quantum computing takes flight

William D. Oliver

A programmable quantum computer has been reported to outperform the most powerful conventional computers in a specific task – a milestone in computing comparable in importance to the Wright brothers' first flights. See p.505





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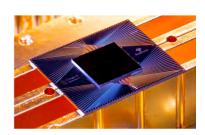
News & views

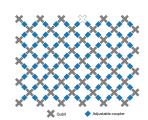
Quantum information

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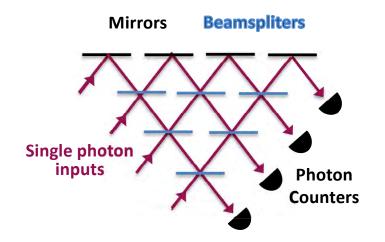
A programmable quantum computer has been reported to outperform the most powerful conventional computers in a specific task — a milestone in computing comparable in importance to the Wright brothers' first flights. See p.505

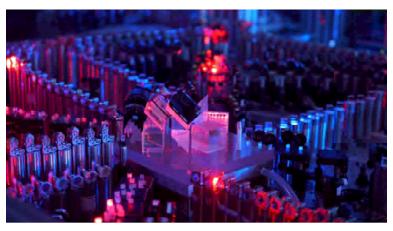




Boson Sampling

An example from Optics/Photonics Setup

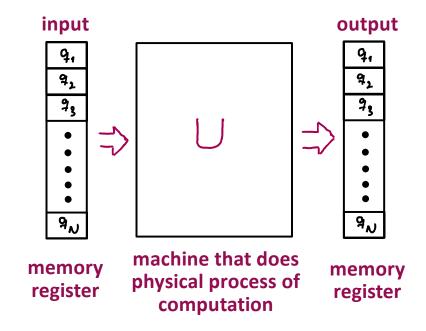




An optical quantum computer developed by a team of Chinese researchers including those from the University of Science and Technology of China. (courtesy of Han-Sen Zhong of the research group)

Back to Universal Computation

Visualization of Computation

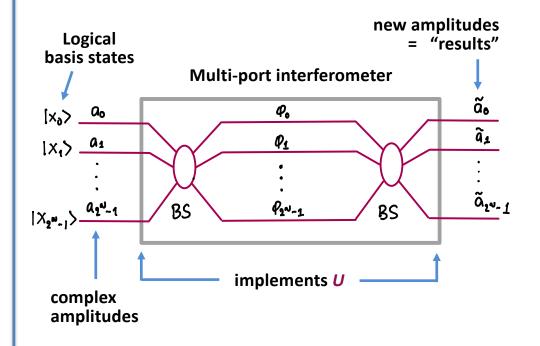


Classical: Register is in one of the logical states

Reversible transformation U: ×

What might be inside the machine?

Wave interference w/classical fields?



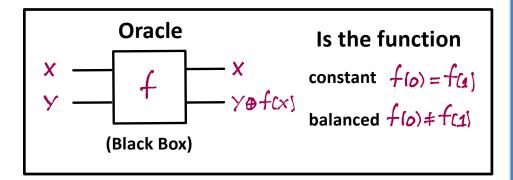
Note: N - qubit register | 2^N "paths"

Beware of Resource Scaling!

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Quantum Advantage

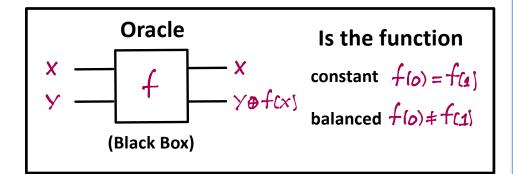
David Deutsch: Toy problem that shows Quantum Advantage



Classical Box: Need 2 queries $f(0) \otimes f(1)$

Quantum Advantage

David Deutsch: Toy problem that shows Quantum Advantage



Quantum Box: In 3 steps can show that

(1)
$$U_f:|x\rangle|y\rangle \Rightarrow |x\rangle|y\oplus f(x)\rangle$$

(2)
$$U_f: |X\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) \rightarrow |X\rangle \frac{1}{\sqrt{2}} (|f(x)\rangle - |1\otimes f(x)\rangle)$$

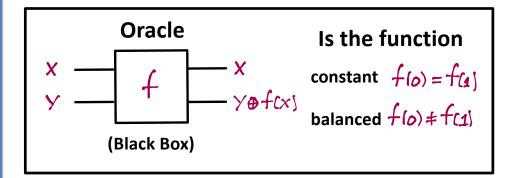
$$=|x>\frac{1}{\sqrt{2}}(-1)^{f(x)}(10>-14>)$$

(3)
$$U_{f}: \sqrt{\frac{1}{2}}(10) + (1) \frac{1}{\sqrt{2}}(10) - (1)$$

$$\rightarrow \frac{1}{2}((-1)^{f(0)}) + (-1)^{f(1)}(1) (10) - (12)$$

Quantum Advantage

David Deutsch: Toy problem that shows Quantum Advantage



Quantum Computation:

Input
$$|x\rangle|\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)\frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$

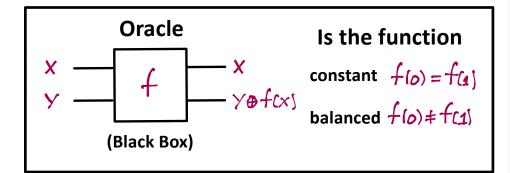
$$U_{f}: \frac{1}{2}([-1)^{f(0)}]0>+(-1)^{f(1)}]1>)([0>-11>)$$

Measure 1st qubit in basis
$$|\pm\rangle = \frac{1}{\sqrt{2}} (|0\rangle \pm |4\rangle)$$

Quantum Speedup: can solve w/1 query

Quantum Advantage

David Deutsch: Toy problem that shows Quantum Advantage



Quantum Computation:

Input
$$|x\rangle|y\rangle = \frac{1}{\sqrt{2}}(|0\rangle+|1\rangle)\frac{1}{\sqrt{2}}(|0\rangle-|1\rangle)$$

$$U_{f}: \frac{1}{2}((-1)^{f(0)})0>+(-1)^{f(1)})(0>-11>)$$

Measure 1st qubit in basis $|\pm\rangle = \frac{1}{\sqrt{2}} (|0\rangle \pm |1\rangle)$

→ |+> if constant, |-> if balanced

Quantum Speedup: can solve w/1 query

Key aspect of Deutsch's algorithm:
We are looking for a global property
of the function f

Generally:
$$U_{+}: [x>[0> \rightarrow] \times > |f(x)>$$

Input $|\Psi_{in}> = \left[\frac{1}{\sqrt{2}}(|0>+|1>)\right]^{\otimes N}|0>$
 $= \frac{1}{2^{N/2}} \sum_{x=0}^{2^{N}-1} |x>|0>$

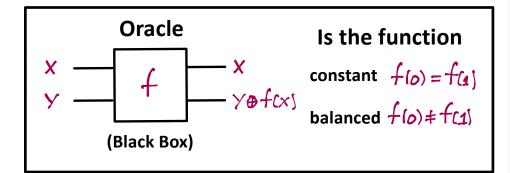
compute once

Output $|\Psi_{in}> = \frac{1}{2^{N/2}} \sum_{x=0}^{2^{N}-1} |x>|f(x)>$

Global properties encoded in state, trick is to extract desired information

Quantum Advantage

David Deutsch: Toy problem that shows Quantum Advantage



Quantum Computation:

Input
$$|x\rangle|y\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)\frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$

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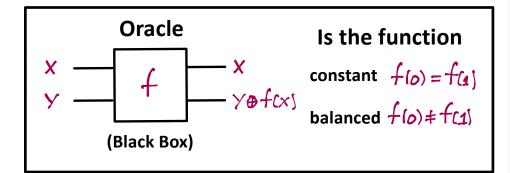
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 $|x| > |x| > |x|$

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Quantum Advantage

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Quantum Computation:

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Key aspect of Deutsch's algorithm:
We are looking for a global property
of the function *f*

Generally:
$$U_{f}: [x>[0> \rightarrow 1\times > 1+cx]>$$

Input $[\Psi_{in}> = \left(\frac{1}{\sqrt{2}}(10> + 11>)\right)^{\otimes N} [0>$
 $=\frac{1}{2^{N-1}}\sum_{x=0}^{2^{N}-1}|x>10>$

compute once

Output $[\Psi_{in}> = \frac{1}{2^{N/2}}\sum_{x=0}^{2^{N}-1}|x>10>$

Peter Shor: Period finding, QFT, Factoring

Key aspect of Deutsch's algorithm:
We are looking for a global property
of the function *f*

Generally:
$$U_{f}: \{x > \{o > \rightarrow \} | x > | f(x) > \}$$

Input $| \mathcal{A}_{in} > = \left[\frac{1}{\sqrt{2}}(1o > + |f(x))\right]^{\bigotimes N} | o > \}$

$$= \frac{1}{2^{N/2}} \sum_{x=0}^{2^{N}-1} |x > 1o > \}$$

compute once

$$| \mathcal{A}_{in} > = \frac{1}{2^{N/2}} \sum_{x=0}^{2^{N}-1} |x > 1f(x) > \}$$

Output $| \mathcal{A}_{oot} > = \frac{1}{2^{N/2}} \sum_{x=0}^{2^{N}-1} |x > |f(x) > \}$

Peter Shor: Period finding, QFT, Factoring

Next: Will this work with real-world Quantum Hardware?

Faulty gates, decoherence!

Quantum Error Correction

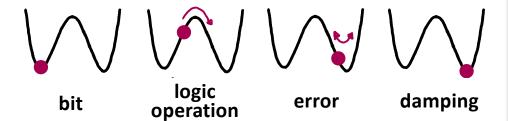
Fundamental Problem



Quantum States are fragile, especially when entangled

Classical Computation

Dissipation helps



No dissipation



Errors build up

Quantum Computation |

- * Cannot tolerate dissipation
- * Destroys superposition and entanglement

What to do? Error Correction!

Classical Error Correction:

Simple example: Redundancy protects against bit flips

Encode:
$$0 \rightarrow (000)$$
 $1 \rightarrow (111)$

Errors:
$$(000) \rightarrow (100)$$
 correct by majority vote

Quantum Computation

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- * Destroys superposition and entanglement

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Classical Error Correction:

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Encode:
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 correct by majority vote

Von Neumann:

- * A classical computer w/faulty components can work, given enough redundancy
- * Classical error correction is well developed and highly sophisticated...

* Quantum Errors

1) Bit Flip
$$\frac{10> \rightarrow 11>}{11> \rightarrow 10>}$$
, phase flip $\frac{10> \rightarrow 10>}{11> \rightarrow -11>}$

- 2) Small errors (alo)+bl1> a,b can change by & errors accumulate
- 4) No cloning Cannot protect by making copies

Von Neumann:

- * A classical computer w/faulty components can work, given enough redundancy
- * Classical error correction is well developed and highly sophisticated...

* Quantum Errors

1) Bit Flip
$$\frac{10> \rightarrow 11>}{11> \rightarrow 10>}$$
, phase flip $\frac{10> \rightarrow 10>}{11> \rightarrow -11>}$

- 2) Small errors (\(\alpha\lorentum\right) + bli\)
 errors accumulate
- 3) Measurement disturbs of quantum states

Example: Peter Shor's code for bit flip error when P(error) << 1Encode: $|0\rangle \rightarrow |0\rangle \geq |000\rangle$ (3 bit code) $|1\rangle \rightarrow |1\rangle \geq |1\rangle \rightarrow |1\rangle$ Single-qubit measurement collapse of state, destroys info, no majority voting!

Collective 2-qubit measurement:

- for $|x,y,\frac{1}{2}\rangle$ measure $\frac{\cancel{200}}{\cancel{200}}$ (never measure individual bits)
- if 1000>, 1111> these observables = 0
- if one bit-flip, at least one observable = 1
- easy to check that (y + 2) = binary address of qubit flip

$$|000\rangle \Rightarrow |010\rangle \qquad (1.0) = 2nd bit$$

Example: Peter Shor's code for bit flip

error when P (error) << 1

Encode: $\frac{10>\rightarrow 10>\approx 1000>}{11>\rightarrow 17>\equiv 1111>}$ (3 bit code)

alo>+b11> - alooo>+b1111>

Single-qubit measurement

collapse of state, destroys info, no majority voting!

Collective 2-qubit measurement:

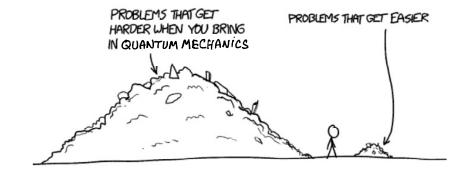
- for $|x,y,2\rangle$ measure $\frac{\cancel{200}}{\cancel{200}}$ (never measure individual bits)
- if 1000 >, 1111 > these observables = 0
- if one bit-flip, at least one observable = 1
- easy to check that (y €2,×⊕2) = binary address of qubit flip

$$|000\rangle \Rightarrow |010\rangle$$
 $(1,0) = 2nd bit$

Small errors: $|000\rangle \rightarrow |000\rangle + \mathcal{E}|001\rangle$ $|111\rangle \rightarrow |111\rangle + \mathcal{E}|110\rangle$

Quantum mechanics to the rescue!

- mostly no error detected
 - collapse into 1000 > resp. 1111>
- sometime error detected
 - ollapse into foot resp. [110]
 - full bit flip, correct as such



Source: xkcd.com

Example: Peter Shor's code for bit flip

error when P (error) << 1

Encode: $|0\rangle \rightarrow |\overline{0}\rangle \ge |000\rangle$ (3 bit code)

alo>+b11> - alooo>+b[111>

Single-qubit measurement

collapse of state, destroys info, no majority voting!

Collective 2-qubit measurement:

- for $|x,y,2\rangle$ measure $\begin{cases} x \oplus 2 \\ x \oplus 2 \end{cases}$ (never measure individual bits)
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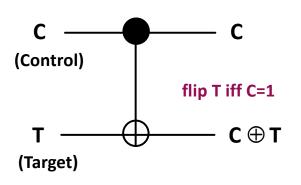
How to implement?

Quantum circuit + single qubit measurement

Quantum Gates – work on superpositions, and entangled states

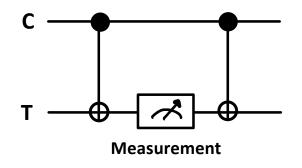
Controlled-NOT (CNOT)

Truth Table



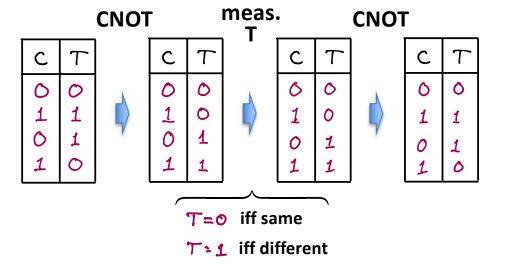
С	T	C ⊕ T
G	0	0
0	1	1
1	0	1
1	1	0

Quantum Circuit for joint measurement

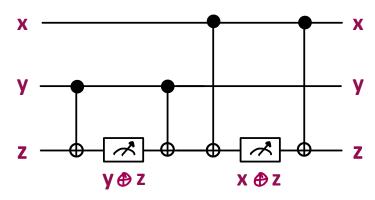


Measurement in {10>, 11>} basis
yields C⊕T

Circuit maps logical basis states as



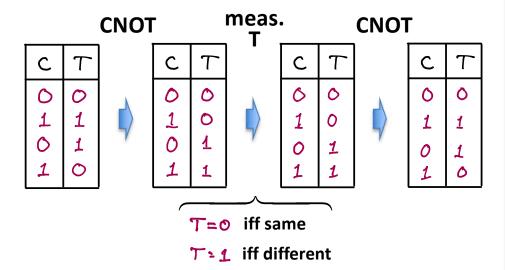
Full circuit to obtain Error Syndrome



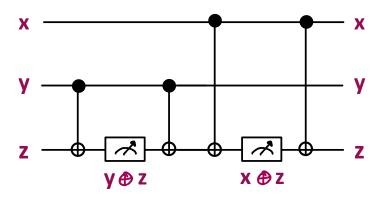
* iff qubit flip, binary address = (y €2,×⊕2)

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Circuit maps logical basis states as



Full circuit to obtain Error Syndrome



* iif qubit flip, binary address = (ヶ段, ×やと)

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Quantum Phase Error

Encoding
$$|0\rangle \rightarrow |\overline{0}\rangle = \frac{1}{2^{1/2}} (|0\rangle + |1\rangle) (|0\rangle + |1\rangle) (|0\rangle + |1\rangle)$$

$$|1\rangle \rightarrow |\overline{1}\rangle = \frac{1}{2^{3/2}} (|0\rangle - |1\rangle) (|0\rangle - |1\rangle) (|0\rangle - |1\rangle)$$

Relabel
$$\frac{\frac{1}{\sqrt{2}}(10)+11)}{\frac{1}{\sqrt{2}}(10)-11)} = 10'$$

Measure
$$\{|0'\rangle, |t'\rangle\} \rightarrow \gamma' \oplus 2', x' \oplus 2'$$
 in basis

Error Syndrome

- * Iff phase error, binary address = $(\gamma' \oplus 2')$
- Analogous to bit-flip code, just in different basis

Quantum Phase Error

Encoding

$$|0\rangle \rightarrow |\overline{0}\rangle = \frac{1}{2^{\frac{1}{2}}} (|0\rangle + |1\rangle) (|0\rangle + |1\rangle) (|0\rangle + |1\rangle)$$

$$|1\rangle \rightarrow |\overline{1}\rangle = \frac{1}{2^{\frac{1}{2}}} (|0\rangle - |1\rangle) (|0\rangle - |1\rangle) (|0\rangle - |1\rangle)$$

Relabel

$$\frac{1}{\sqrt{2}}(10) + 11) = 11'$$

Measure in basis

Error Syndrome

- * Iff phase error, binary address = $(\gamma' \oplus 2')$
- * Analogous to bit-flip code, just in different basis

Shor's 9-bit code

- * Combines flip/phase error correction
- * Corrects one flip or phase error

General principle of error correction

- * Encode p logical qubits in n physical qubits.
- * Valid Logical States form 2^p -dimensional subspace \mathcal{E}_p (code space) in n-qubit $(2^n$ -dimensional) Hilbert space \mathcal{E}_N
- * Errors displace system into orthogonal (distinguishable) subspaces.

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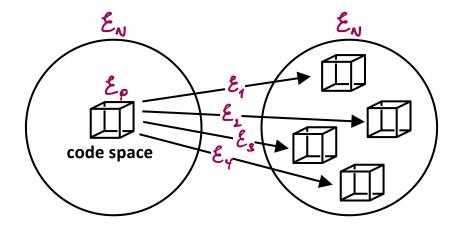
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Geometric illustration



What about non-Unitary errors?

e. g., decay
$$\begin{array}{c} 107 \rightarrow 107 \\ \hline 117 \rightarrow 107 \end{array}$$

Problem: Errors not displaced into

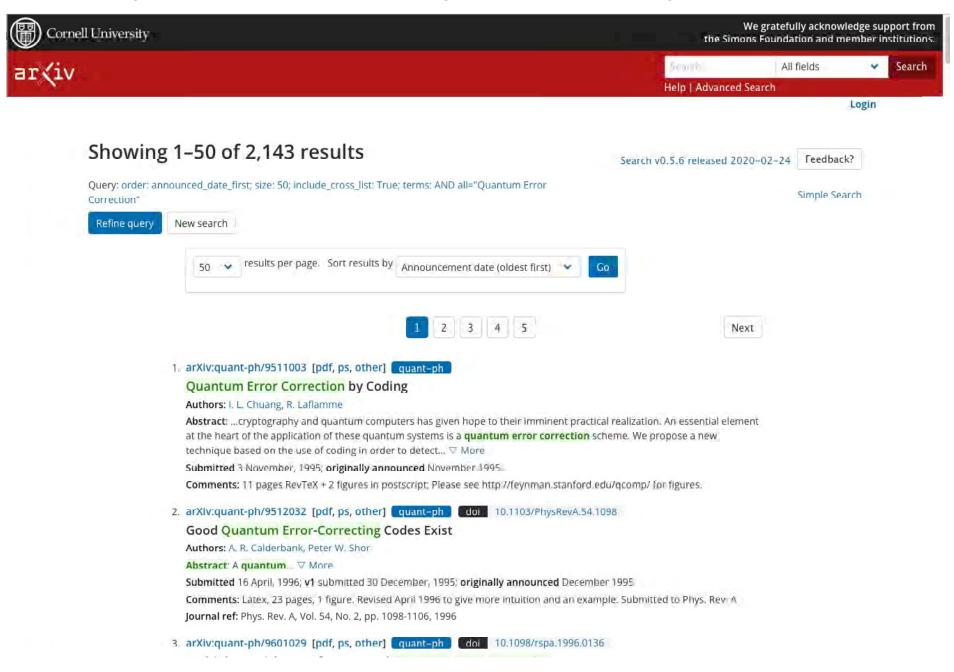
orthogonal subspaces

Solution: "Quantum jump codes",

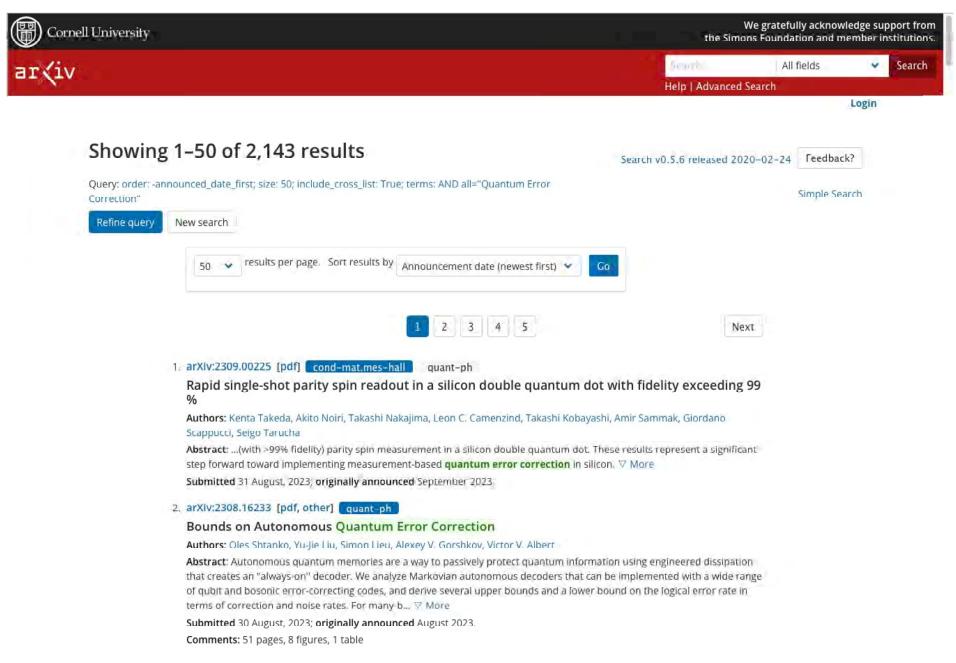
monitors the environment

Other kinds of errors?

Catnip for Theoretical Physicists & Computer Scientists



Catnip for Theoretical Physicists & Computer Scientists



Quantum Hardware

Physical Implementation is Extremely demanding!

Requirements

- 1. Storage: Quantum memory.
- 2. Gates: We put computation U_f together from 1 and 2-qubit operations.
- 3. Readout: Method to measure qubits.
- 4. Isolation: No coupling to environment to avoid decoherence & errors
- 5. Precision: Gates, readouts must be highly accurate

Inherent Contradictions

2. Gates vs 4. Isolation

coupling between qubits no coupling to environment

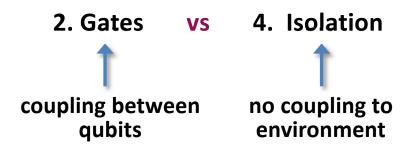


To build a Quantum Computer



Choose, find or invent a system with acceptable tradeoffs

Inherent Contradictions





To build a Quantum Computer

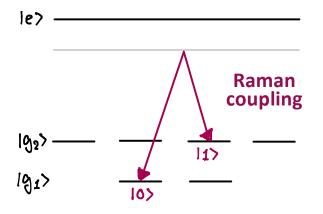


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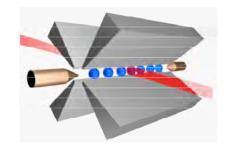
Ion Trap Quantum Computing

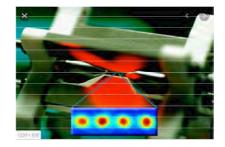
First to demonstrate a Quantum Gate

* Qubit is encoded in the electronic ground state of an atomic ion



* Early design with a few ions in large trap

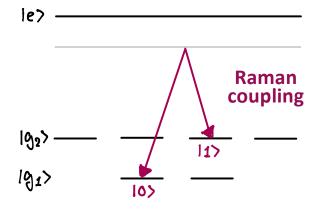




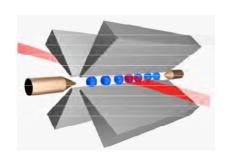
Ion Trap Quantum Computing

First to demonstrate a Quantum Gate

* Qubit is encoded in the electronic ground state of an atomic ion



* Early design with a few ions in large trap





Requirements

1. Storage: 10s-100s coherence time

2. Gates: Use collective vibrations as

"quantum bus"

3. Readout: Fluorescence



Cirac & Zoller: 5 laser pulses 🔷



CNOT gate between any **2** ions in linear array

Wineland: 3 laser pulses enough for CNOT

Use this example serves as conceptual template

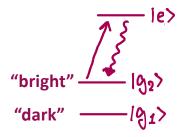
Requirements

1. Storage: 10s-100s coherence time

2. Gates: Use collective vibrations as

"quantum bus"

3. Readout: Fluorescence

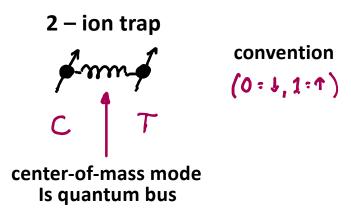


Cirac & Zoller: 5 laser pulses 🔷

CNOT gate between any **2** ions in linear array

Wineland: 3 laser pulses enough for CNOT

Use this example serves as conceptual template



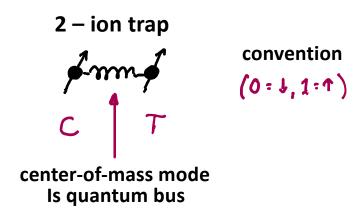
CNOT input-output map:

$$| \mathcal{L}_{c} \mathcal{L}_{\tau} \rangle \rightarrow | \mathcal{L}_{c} \mathcal{L}_{\tau} \rangle \qquad | \mathcal{L}_{c} \mathcal{L}_{\tau} \rangle \rightarrow | \mathcal{L}_{c} \mathcal{L}_{\tau} \rangle \\ | \mathcal{L}_{c} \mathcal{L}_{\tau} \rangle \rightarrow | \mathcal{L}_{c} \mathcal{L}_{\tau} \rangle \qquad | \mathcal{L}_{c} \mathcal{L}_{\tau} \rangle \rightarrow | \mathcal{L}_{c} \mathcal{L}_{\tau} \rangle \\ | \text{does nothing} \qquad \text{swaps} \quad | \mathcal{L}_{c} \mathcal{L}_{\tau} \rangle, \quad | \mathcal{L}_{c} \mathcal{L}_{\tau} \rangle$$

Notation for quantum states: $|n, \sigma_c^{(2)}, \sigma_T^{(2)}\rangle$

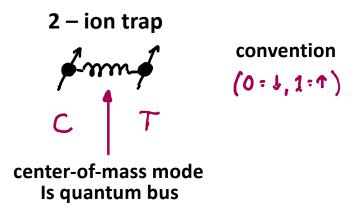
n: vibrational quantum number

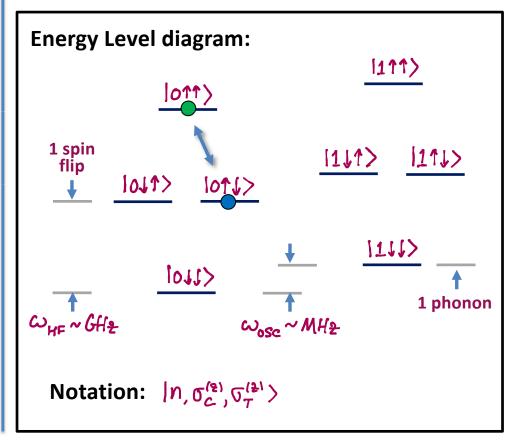
 $\sigma_{C}^{(2)}, \sigma_{T}^{(2)}$: spin states of C and T ions

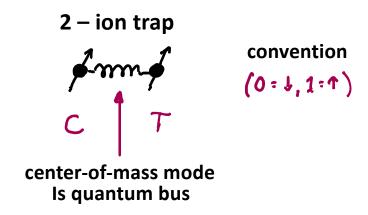


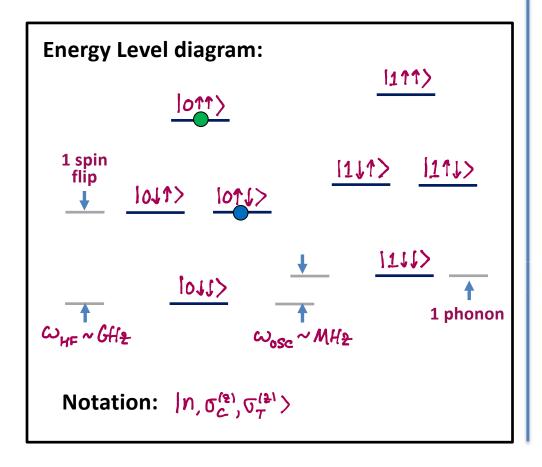
CNOT input-output map:

Notation for quantum states: $|n, \sigma_c^{(2)}, \sigma_T^{(2)}\rangle$









CNOT input-output map:

$$|\mathcal{L}_{c}|_{T} > -|\mathcal{L}_{c}|_{T} > \qquad |\mathcal{L}_{c}|_{T} > -|\mathcal{L}_{c}|_{T} > \qquad |\mathcal{L}_{c}|_{T} > -|\mathcal{L}_{c}|_{T} > \qquad |\mathcal{L}_{c}|_{T} > \qquad |\mathcal{L}_{c}|_{T}$$

Laser pulse sequence:

- (1) π pulse on C swaps $\{o\uparrow_{c}x_{\tau}\} \leftrightarrow \{1\downarrow_{c}x_{\tau}\}$
- (2) π pulse on T swaps $|1\downarrow_c\downarrow_{\tau}\rangle \Leftrightarrow |1\downarrow_c\uparrow_{\tau}\rangle$
- (3) π pulse on C swaps $|0 \uparrow_C X_T \rangle \approx |1 \downarrow_C X_T \rangle$

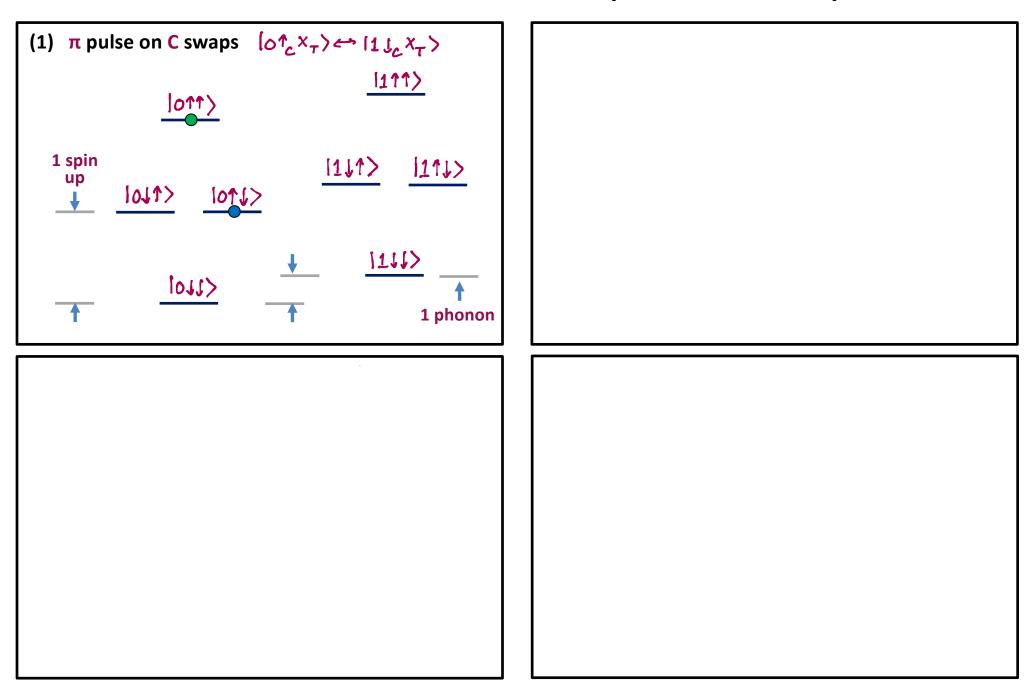


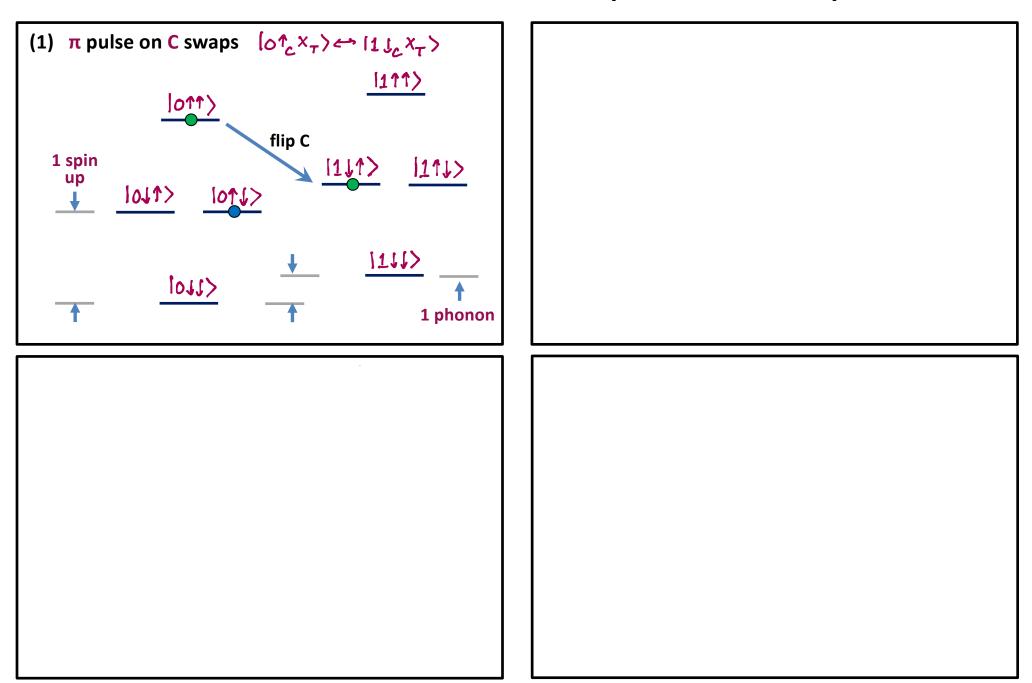
Performs CNOT Gate

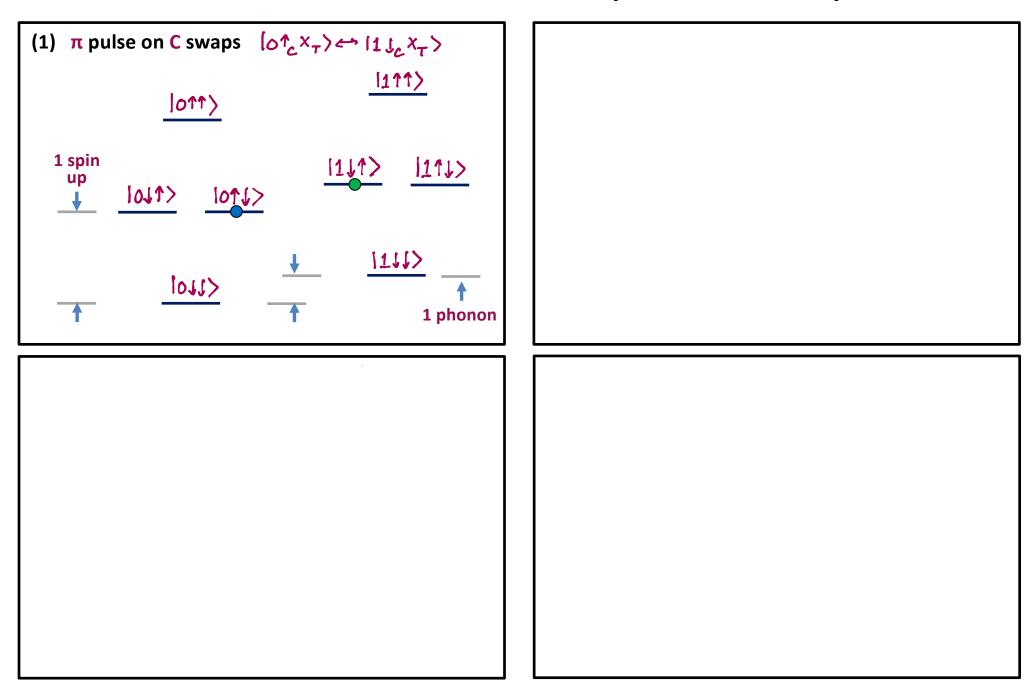
Let us go through the process in detail!

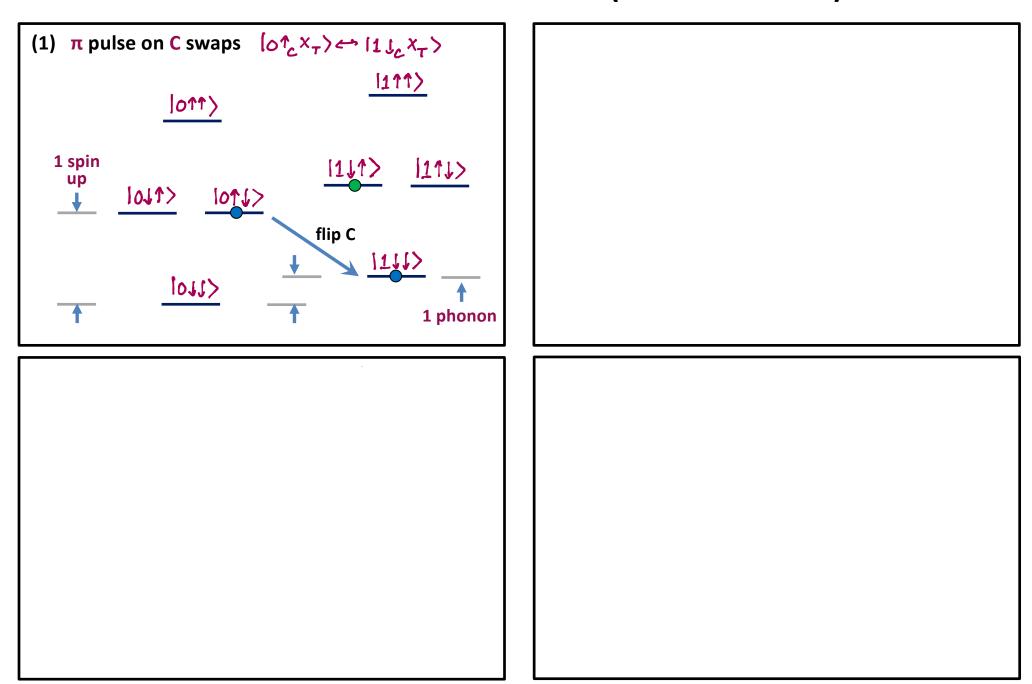
(Note: Clickthrough Panels)

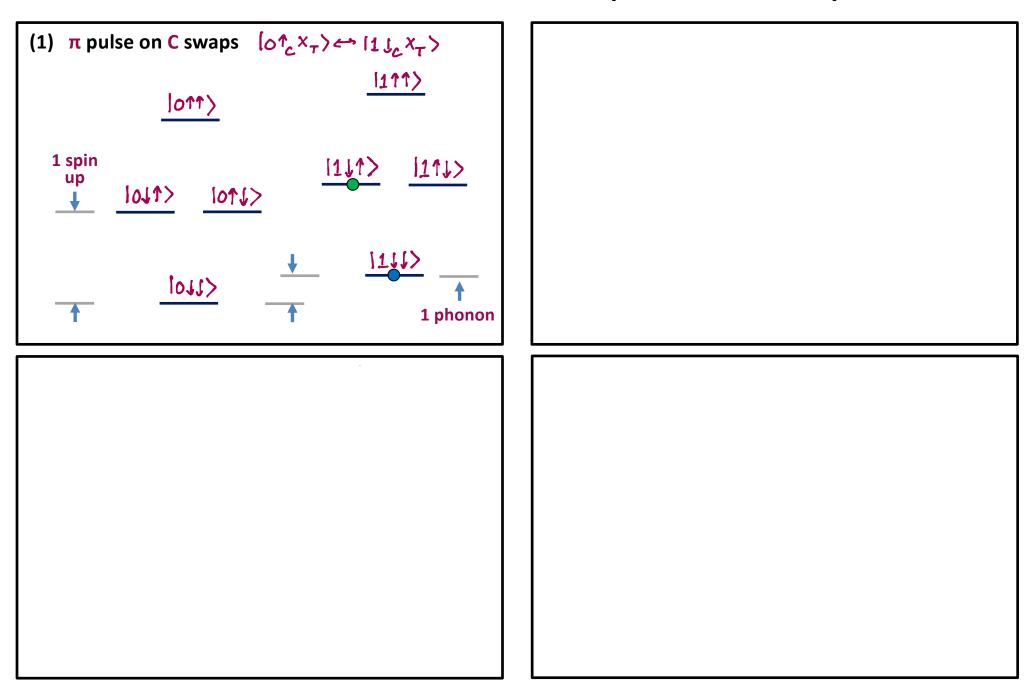
(1) π pulse on C swaps $\{ \Diamond \uparrow_c \times_{\tau} \rangle \leftrightarrow \{ 1 \}_c \times_{\tau} \}$	
-	











(1)
$$\pi$$
 pulse on C swaps $[0\uparrow_c \times_{\tau}\rangle \leftrightarrow [1\downarrow_c \times_{\tau}\rangle$

$$\frac{|1\uparrow\uparrow\rangle}{|0\uparrow\uparrow\rangle}$$

$$\frac{|1\uparrow\uparrow\rangle}{|1\uparrow\downarrow\rangle}$$

$$\frac{|0\downarrow\uparrow\rangle}{|1\downarrow\downarrow\rangle}$$

$$\frac{|0\downarrow\downarrow\rangle}{|1\downarrow\downarrow\rangle}$$

$$\frac{|1\downarrow\downarrow\rangle}{|1\downarrow\downarrow\rangle}$$

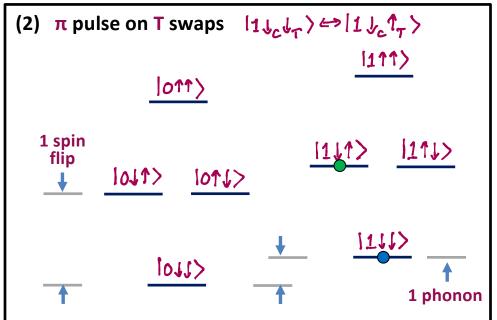
$$\frac{|1\downarrow\downarrow\rangle}{|1\downarrow\downarrow\rangle}$$

$$\frac{|1\downarrow\downarrow\rangle}{|1\downarrow\downarrow\rangle}$$

$$\frac{|1\downarrow\downarrow\rangle}{|1\downarrow\downarrow\rangle}$$

$$\frac{|1\downarrow\downarrow\rangle}{|1\downarrow\downarrow\rangle}$$





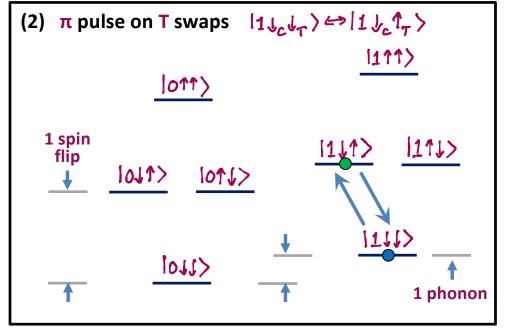


(1)
$$\pi$$
 pulse on C swaps $[0\uparrow_c \times_{\tau}\rangle \leftrightarrow [1 \downarrow_c \times_{\tau}\rangle$

$$\frac{|1\uparrow\uparrow\rangle}{|1\uparrow\downarrow\rangle}$$
1 spin $\frac{|1\downarrow\uparrow\rangle}{|1\uparrow\downarrow\rangle}$

$$\frac{|0\downarrow\uparrow\rangle}{|1\uparrow\downarrow\rangle}$$
1 phonon





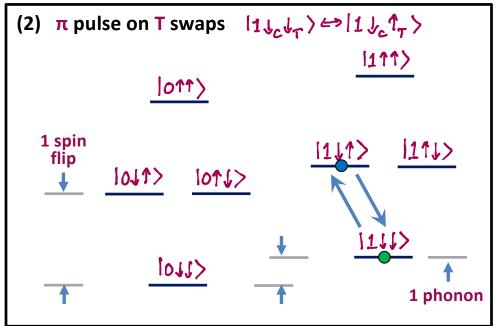


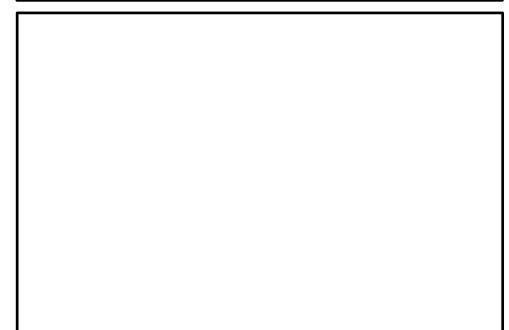
(1)
$$\pi$$
 pulse on C swaps $[0\uparrow_c \times_{\tau}\rangle \leftrightarrow [1 \downarrow_c \times_{\tau}\rangle$

$$\frac{|0\uparrow\uparrow\rangle}{|0\uparrow\uparrow\rangle}$$
1 spin $\frac{|1\downarrow\uparrow\rangle}{|1\uparrow\downarrow\rangle}$

$$\frac{|0\downarrow\uparrow\rangle}{|1\uparrow\downarrow\rangle}$$
1 phonon







(1)
$$\pi$$
 pulse on C swaps $[0\uparrow_c \times_{\tau}\rangle \leftrightarrow [1\downarrow_c \times_{\tau}\rangle$

$$\frac{|0\uparrow\uparrow\rangle}{|0\uparrow\uparrow\rangle}$$

$$\frac{|1\downarrow\uparrow\rangle}{|1\downarrow\uparrow\rangle}$$

$$\frac{|0\downarrow\uparrow\rangle}{|1\downarrow\downarrow\rangle}$$

$$\frac{|0\downarrow\downarrow\rangle}{|1\downarrow\downarrow\rangle}$$

$$\frac{|0\downarrow\downarrow\rangle}{|1\downarrow\downarrow\rangle}$$

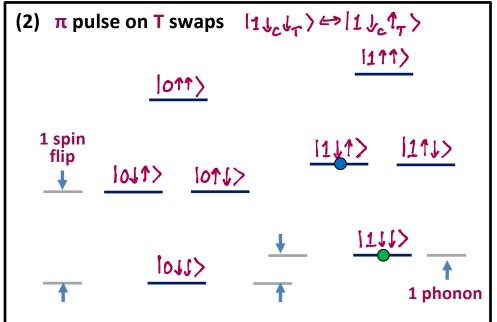
$$\frac{|1\downarrow\downarrow\rangle}{|1\downarrow\downarrow\rangle}$$

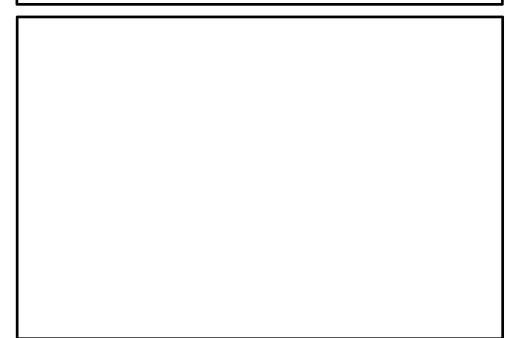
$$\frac{|1\downarrow\downarrow\rangle}{|1\downarrow\downarrow\rangle}$$

$$\frac{|1\downarrow\downarrow\rangle}{|1\downarrow\downarrow\rangle}$$

$$\frac{|1\downarrow\downarrow\rangle}{|1\downarrow\downarrow\rangle}$$



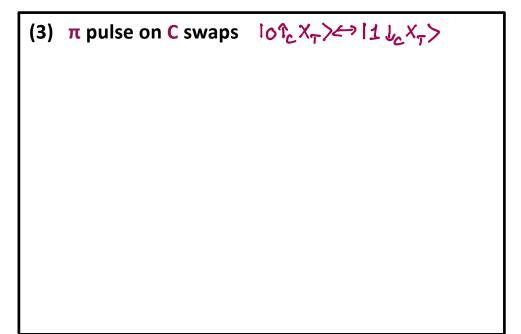




(1)
$$\pi$$
 pulse on C swaps $[0\uparrow_c \times_{\tau}\rangle \leftrightarrow [1 \downarrow_c \times_{\tau}\rangle$

$$\frac{|0\uparrow\uparrow\rangle}{|0\uparrow\uparrow\rangle}$$
1 spin $\frac{|1\downarrow\uparrow\rangle}{|1\uparrow\downarrow\rangle}$

$$\frac{|0\downarrow\uparrow\rangle}{|1\uparrow\downarrow\rangle}$$
1 phonon



```
(2) \pi pulse on T swaps |1\downarrow_c\downarrow_{\uparrow}\rangle \Leftrightarrow |1\downarrow_c\uparrow_{\uparrow}\rangle
|0\uparrow\uparrow\rangle
|1\uparrow\uparrow\rangle
|1\uparrow\uparrow\rangle
|1\uparrow\uparrow\rangle
|1\uparrow\uparrow\rangle
|1\uparrow\uparrow\rangle
|1\uparrow\downarrow\rangle
|1\downarrow\downarrow\rangle
|1\downarrow\rangle
|1\rangle
|1\rangle
```



(1)
$$\pi$$
 pulse on C swaps $[o\uparrow_c \times_{\tau}\rangle \leftrightarrow [1\downarrow_c \times_{\tau}\rangle$

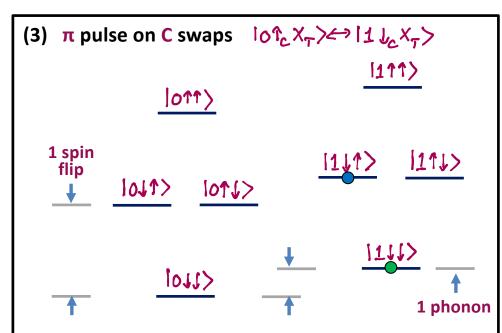
$$|o\uparrow\uparrow\rangle$$

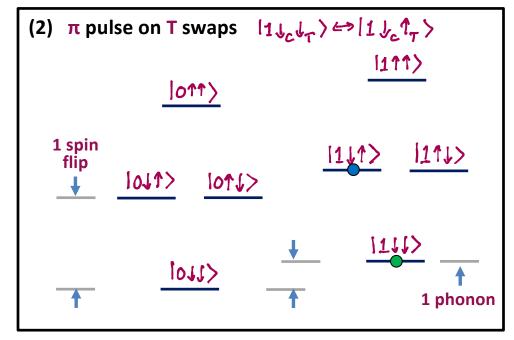
$$|spin \text{ up} \text{ } |o\downarrow\uparrow\rangle$$

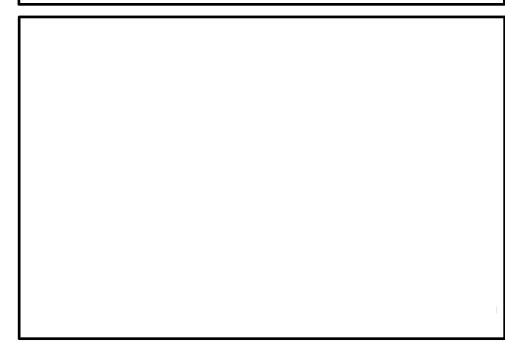
$$|o\uparrow\downarrow\rangle$$

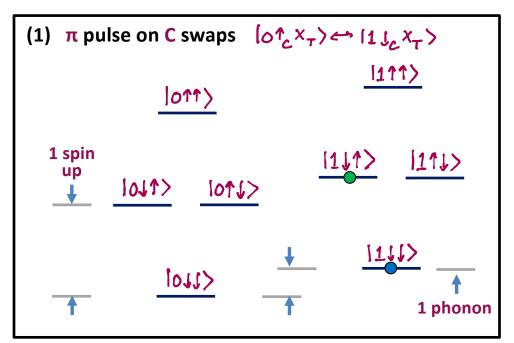
$$|i\downarrow\uparrow\rangle$$

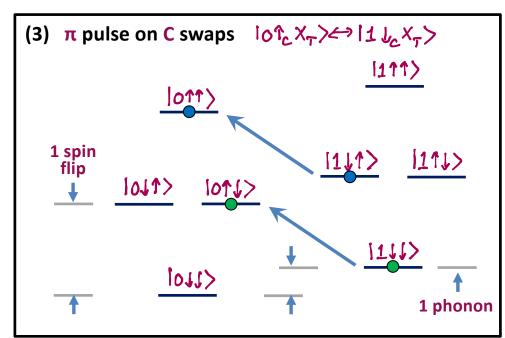
$$|i\downarrow\downarrow\rangle$$

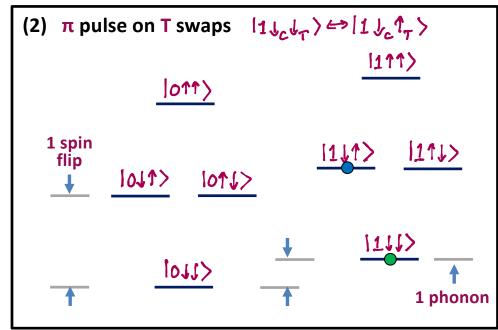


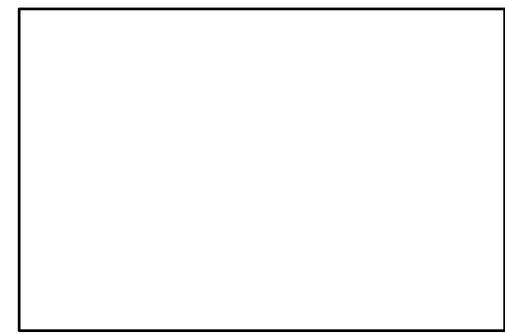






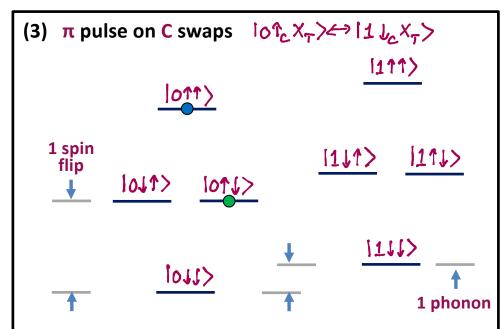


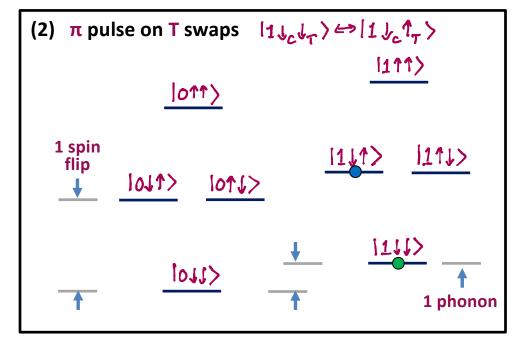


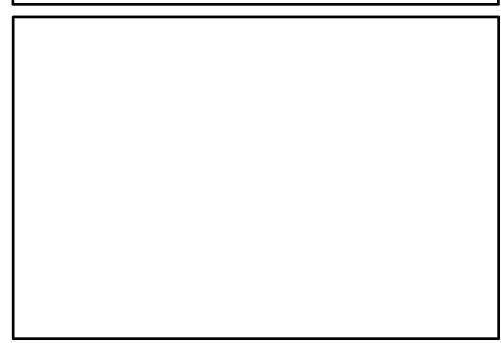


(1)
$$\pi$$
 pulse on C swaps $[0\uparrow_c \times_{\tau}\rangle \leftrightarrow [1\downarrow_c \times_{\tau}\rangle$

$$\frac{|1\uparrow\uparrow\rangle}{|1\uparrow\uparrow\rangle}$$
1 spin up |10 $\downarrow\uparrow\rangle$ |11 $\downarrow\uparrow\rangle$ |11 $\downarrow\downarrow\rangle$ |11 $\downarrow\rangle$ |11 $\downarrow\downarrow\rangle$ |11 $\downarrow\downarrow\rangle$ |11 $\downarrow\downarrow\rangle$ |11 $\downarrow\downarrow\rangle$ |11 $\downarrow\downarrow\rangle$ |11 $\downarrow\downarrow\rangle$ |11 $\downarrow\downarrow\rangle$ |11 $\downarrow\downarrow\rangle$ |11 $\downarrow\downarrow\rangle$ |1



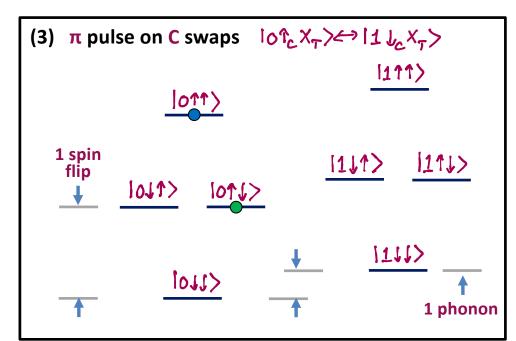


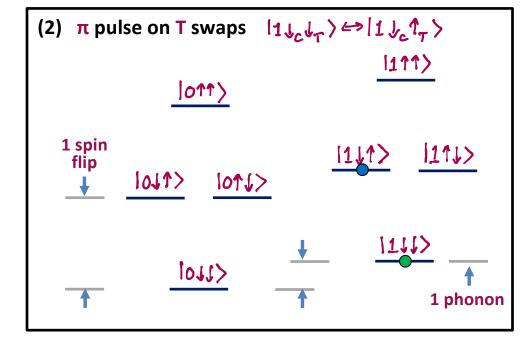


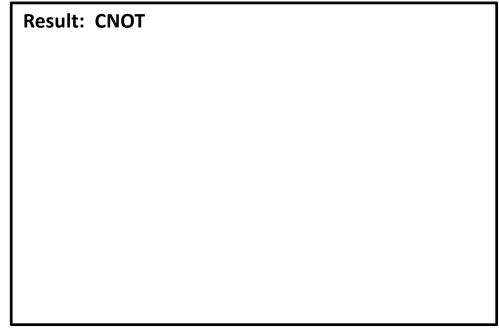
(1)
$$\pi$$
 pulse on C swaps $[0\uparrow_c \times_{\tau}\rangle \leftrightarrow [1 \downarrow_c \times_{\tau}\rangle$

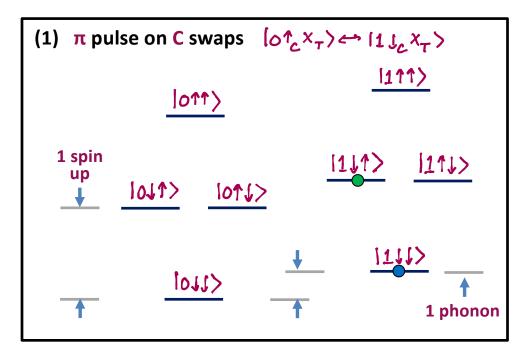
$$\frac{|0\uparrow\uparrow\rangle}{|0\uparrow\uparrow\rangle}$$
1 spin $\frac{|1\downarrow\uparrow\rangle}{|1\uparrow\downarrow\rangle}$

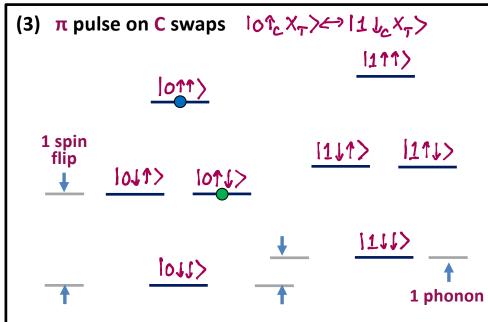
$$\frac{|0\downarrow\uparrow\rangle}{|1\uparrow\downarrow\rangle}$$
1 phonon

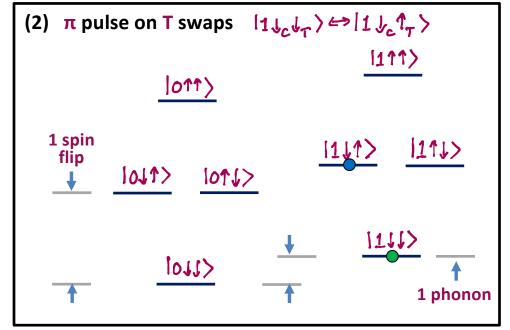


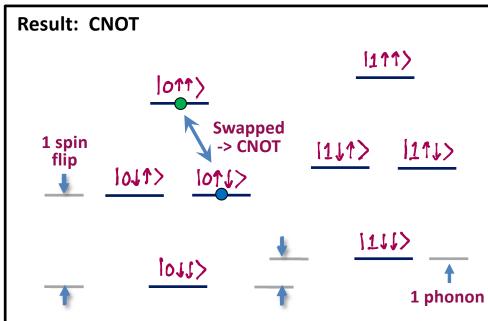


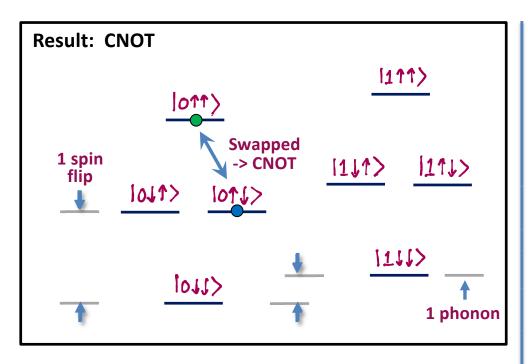


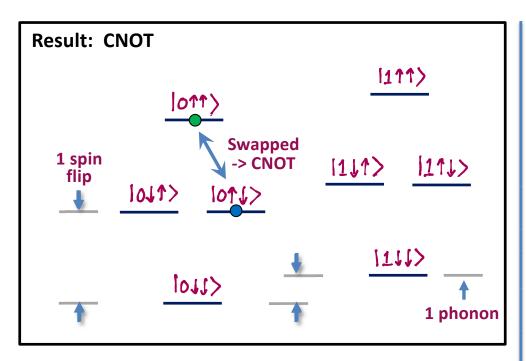












Note: The sequence (1) -> (2) -> (3) leaves the spin and vibrational degrees of freedom unentangled after the CNOT

Note: Todays Ion Trap QC experiments rely on much more sophisticated, accurate and robust gate protocols

Status: Many important milestones achieved

- * Entanglement of ≥ 20 ions (2018)
- * Highest gate & readout fidelities, longest coherence times
- * Error Correction, Fault Tolerance proof of principle demonstrations
- * Complex algorithms on few ions, quantum simulations with ≥ 50
- * Research groups in academia, National Labs, Industry

Some leading groups

NIST Innsbruck Quantinuum

Sandia NL Duke U IonQ

Many, many others

Some links to get started

Amazon Braket (IonQ, other Technologies) https://aws.amazon.com/braket/

Quantinuum (Ion Trap Quantum Computing https://www.quantinuum.com

lonQ https://ionq.com

NIST https://www.nist.gov/pml/time-and-frequency-division/ion-storage

Innsbrück

https://www.uibk.ac.at/th-physik/qic-group/