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Light-Matter Interaction

Hierarchy of Sophistication:

- Classical — Classical light, classical matter
- Semiclassical — Classical light, quantum matter
- Quantum — Quantum light, quantum matter

Possible attitudes:

- Purist — Most complete description possible
- Minimalist — Quantum only when necessary
- Pragmatic — Quantum or classical, based on what is simplest and still works

OPTI 544: All of the above in turn

Classical Theory of Light-Matter Interaction

Self-consistent, fully classical description

Electromagnetic Field $\rightarrow$ Atom/Molecule/Solid

Motivation: We will

- Develop Concepts $\alpha(\omega), n, \chi$
- Develop Intuition
- Classical is often adequate, sometimes accurate
- A Quantum Theory has classical limits $\Rightarrow$ Identify/understand regime of validity
- The Classical description is a useful starting point for Nonlinear and Quantum Optics

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The Electromagnetic Field: Basic Eqs. in SI Units

Maxwell's eqs.

(no free charges, currents ⇨ dielectrics)

(i) $\nabla \cdot \vec{D} = \rho = 0$ — $\vec{D}$: Dielectric displacement

(ii) $\nabla \cdot \vec{B} = 0$ — $\vec{B}$: Magnetic induction

(iii) $\nabla \times \vec{E} = -\dfrac{\partial \vec{B}}{\partial t}$ — $\vec{E}$: Electric field

(iv) $\nabla \times \vec{H} = \dfrac{\partial \vec{D}}{\partial t} + \vec{J}$ — $\vec{H}$: Magnetic field

Material Response:

(v) $\vec{B} = \mu_0 \vec{H} + \vec{M}$ ⟵ Non-magnetic ⇨ $\vec{M} = 0$

(vi) $\vec{D} = \varepsilon_0 \vec{E} + \vec{P}$ ⟵ Info about response in dipole moment density (polarization density)

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We need equations that describe:

- the behavior of $\vec{E}$ for given $\vec{P}$
- the medium response $\vec{P}$ for given $\vec{E}$

Wave Equation:

Take curl of (iii), then use (iv)

$$\nabla\times(\nabla\times\vec{E}) = -\nabla\times\frac{\partial\vec{B}}{\partial t} = -\frac{\partial}{\partial t}(\nabla\times\vec{B}) = -\mu_0\frac{\partial^2\vec{D}}{\partial t^2}$$

Next, use the identity

$$\nabla\times(\nabla\times\vec{E}) = \nabla(\nabla\cdot\vec{E}) - \nabla^2\vec{E}$$

to obtain $\vec{D} = \nabla(\nabla\cdot\vec{E}) = -\mu_0\dfrac{\partial^2\vec{D}}{\partial t^2}$

Finally, let $\vec{D} = \varepsilon_0\vec{E} + \vec{P}$ and use $\varepsilon_0\mu_0 = \dfrac{1}{c^2}$

to obtain

$$\boxed{-\nabla(\nabla\cdot\vec{E}) + \nabla^2\vec{E} = \frac{1}{c^2}\frac{\partial^2\vec{E}}{\partial t^2} + \frac{1}{\varepsilon_0 c^2}\frac{\partial^2\vec{P}}{\partial t^2}}$$

This is the Wave Equation, still exact in this form

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Transverse Fields

Definition: a field for which $\nabla\cdot\vec{E} = 0$ is Transverse

Example: a plane wave, $\vec{E}(\vec{r},t) = \vec{E}(t)\,e^{i\vec{k}\cdot\vec{r}}$,

where $\vec{E}(t)\perp\vec{k}$, is transverse.

The physical field is $Re[\vec{E}(\vec{r},t)]$

For transverse fields the wave equation simplifies to

$$\boxed{\nabla^2\vec{E} - \frac{1}{c^2}\frac{\partial^2}{\partial t^2}\vec{E} = \frac{1}{\varepsilon_0 c^2}\frac{\partial^2}{\partial t^2}\vec{P}}$$

This version of the wave equation can be a poor approximation in non-isotropic media!

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Isotropic Media

Absent a preferred direction, the induced $\vec{P}$ must be parallel to the driving field $\vec{E}$

Linear response, most general case:

$$\vec{D}(t) = \varepsilon_0 \vec{E}(t) + \vec{P}(t)$$
$$= \varepsilon_0 \vec{E}(t) + \varepsilon_0 \int_{-\infty}^{t} dt' R(t-t')\, \vec{E}(t')$$

where the *response function* $R(t-t')$ is a scalar and we have $R(\tau) = 0$ for $\tau < 0$

Take divergence on both sides and use M.E. (i)

$$\nabla\cdot\vec{D}(t) = \varepsilon_0 \nabla\cdot\vec{E}(t) + \varepsilon_0 \int_{-\infty}^{t} dt' R(t-t')\, \nabla\cdot\vec{E}(t') = 0$$

$\Rightarrow$ $\nabla\cdot\vec{E}(t) = -\int_{-\infty}^{t} dt' R(t-t')\, \nabla\cdot\vec{E}(t')$ for all t

It follows that $\nabla\cdot\vec{E}(t) = 0$ for all t,

OR $R(\tau) = -2\delta(\tau)$

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<u>OR</u> $R(\tau) = -2\delta(\tau)$

Note: if $R(\tau) \propto \delta(\tau)$ (instantaneous response) then

$$\varepsilon_0 \int_{-\infty}^{t} dt' R(t-t') \vec{E}(t') = \varepsilon_0 \chi \vec{E}(t)$$

susceptibility

The case $R(\tau) = -2\delta(\tau)$ is an example of negative susceptibility, $\chi < 0$, which only occurs in certain engineered metamaterials.

Electric fields are transverse in linear, isotropic dielectric media

(including the vacuum)

Wave Equation in free space

$$\boxed{\nabla^2 \vec{E} - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \vec{E} = 0}$$

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Light-Matter Interaction

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Monochromatic trial solution $\vec{E}(\vec{r},t) = \vec{E}_0(\vec{r}) e^{-i\omega t}$

$$\nabla^2 \vec{E}_0(\vec{r}) e^{-i\omega t} + \frac{\omega^2}{c^2} \vec{E}_0(\vec{r}) e^{-i\omega t} = 0$$

Equation for the spatial component alone:

$$\nabla^2 \vec{E}_0(\vec{r}) + |\vec{k}|^2 \vec{E}_0(\vec{r}) = 0 \, , \; |\vec{k}| = \omega/c$$

Plane wave solutions

$$\boxed{\vec{E}_0(\vec{r}) = \vec{\varepsilon} E_0 e^{i\vec{k}\cdot\vec{r}} \, , \; |\vec{k}| = \omega/c}$$

<u>Optical Cavities</u>: Here we need to solve the wave equation subject to boundary conditions. See, e. g., Millony & Eberly for examples such as rectangular cavities, Fabryt-Perot etalons, and spherical mirror resonators.

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Wave Equation in Fourier Space:

In Configuration Space:

$$\nabla^2 \vec{E}(\vec{r},t) - \frac{1}{c^2}\frac{\partial^2}{\partial t^2}\vec{E}(\vec{r},t) = \frac{1}{\varepsilon_0 c^2}\frac{\partial^2}{\partial t^2}\vec{P}(\vec{r},t)$$

In Fourier Space:

$$\boxed{k^2 \vec{E}(\vec{k},\omega) - \frac{\omega^2}{c^2}\vec{E}(\vec{k},\omega) = \frac{\omega^2}{\varepsilon_0 c^2}\vec{P}(\vec{k},\omega)}$$

Note: In the Fourier domain the wave equation is purely algebraic – there are no derivatives or integrals. This becomes important later in the course when we quantize the electromagnetic field.

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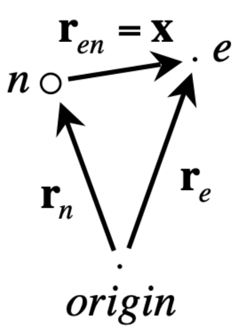
Light-Matter Interaction

Theory of Atomic Response

So far, we have a model for the field. Next, we need a model of how the constituents of the medium responds to the field.

This will allow us to find the polarization density $\vec{P}$ as function of the field $\vec{E}$

Classical "atom"



Simple model:
nucleus + electron

Lorentz Force

$$\boxed{\vec{F} = q(\vec{E} + \vec{v}\times\vec{B})}$$

~ 0 if non-relativistic

Newton:

$$\text{(i)} \quad m_n \frac{d^2}{dt^2}\vec{r}_n(t) = -e\vec{E}(\vec{r}_n,t) - \vec{F}_{en}(\vec{r}_{en},t)$$

$$\text{(ii)} \quad m_e \frac{d^2}{dt^2}\vec{r}_e(t) = e\vec{E}(\vec{r}_e,t) + \vec{F}_{en}(\vec{r}_{en},t)$$

This is a standard 2-body problem which we can re-cast as in terms of relative and COM motion.

We define:

$$\vec{x} = \vec{r}_{en} = \vec{r}_e - \vec{r}_n \qquad m = \frac{m_e m_n}{m_e + m_n} \sim m_e$$

$$\vec{R} = \frac{m_e\vec{r}_e + m_n\vec{r}_n}{M} \qquad M = m_e + m_n \sim m_n$$

$\vec{x}$ Relative coord. $\quad m$ Reduced mass

$\vec{R}$ Center-of-mass $\quad M$ Total mass

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Light-Matter Interaction

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$\vec{x}$ **Relative coord.** $\qquad m$ **Reduced mass**

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Sub into (i), (ii) and rewrite:

$$M\frac{d^2}{dt^2}\vec{R} = e\left[\vec{E}\left(\vec{R}+\frac{m_n}{M}\vec{x},t\right) - \vec{E}\left(\vec{R}-\frac{m_e}{M}\vec{x},t\right)\right]$$

$$m\frac{d^2}{dt^2}\vec{x} = \frac{e}{2}\left[\vec{E}\left(\vec{R}+\frac{m_n}{M}\vec{x},t\right) + \vec{E}\left(\vec{R}-\frac{m_e}{M}\vec{x},t\right)\right] + \vec{F}_{en}(\vec{x}) + \frac{1}{2}(m_n - m_e)\frac{d^2}{dt^2}\vec{R}$$

Basic result, no approximations !

Milloni & Eberly, main text $\Rightarrow$ **Set** $\vec{R} \approx \vec{r}_n$, $\vec{x} \approx \vec{r}_{en}$

Throw away eq. for $\vec{R}$

<u>Electric Dipole approximation</u>

Atomic dimensions		Optical Wavelength
$\lvert\vec{x}\rvert \sim 1\,\text{Å}$	$\ll$	$\lambda \sim 10^4\,\text{Å}$

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<u>EDA</u>: the field is <u>nearly constant</u> on the scale of an atom

Good approximation: 1st order expansion in $\vec{x}$

$$\vec{E}\left(\vec{R}-\frac{m_e}{M}\vec{x},t\right) \approx \vec{E}(\vec{R},t) - \frac{m_e}{M}(\vec{x}\cdot\nabla)\vec{E}(\vec{R}(t))$$

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$$M\frac{d^2}{dt^2}\vec{R} \approx e(\vec{x}\cdot\nabla)E(\vec{R},t) \qquad \textbf{COM}$$

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Physical Interpretation:

$\vec{p} = e\vec{x}$: electric dipole moment of the atom

The Eqs. Of motion can then be recast as

$$M\frac{d^2}{dt^2}\vec{R} \approx (\vec{p}\cdot\nabla)\vec{E}(\vec{R},t) = \vec{F} = -\nabla_R V(\vec{x},\vec{R},t)$$

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where $V(\vec{x},\vec{r},t) = -\vec{p}\cdot\vec{E}(\vec{r},t)$

electric-dipole interaction

Note: The COM Eq. is the foundation for a range of laser Atom Traps and Optical Tweezers. We will not explore this further in OPTI 544 lectures, but good review articles can be found in the published literature.

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Light-Matter Interaction

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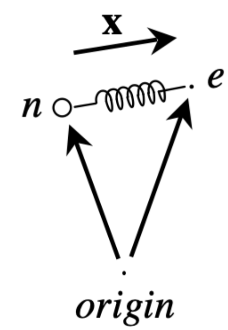
The Electron Oscillator/Lorentz Oscillator

Simple model w/a harmonically bound electron:



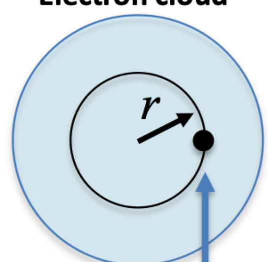
$\vec{F}_{en} = -m\omega_0^2\vec{x}$

resonance frequency

This is meant as a model of the atomic <u>response</u>, not a model of the atom itself.

<u>Nevertheless</u>: QM suggest the atom consists of a point-like nucleus and a spherical electron cloud

Electron cloud



nucleus

Force from charge inside r as if entire charge was at the center

Force from charge outside r is zero

Force $F \propto \frac{r^3}{r^2} \propto r$

harmonic restoring force

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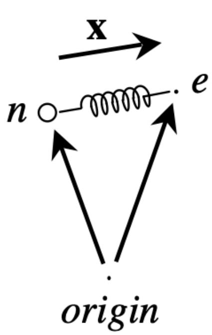

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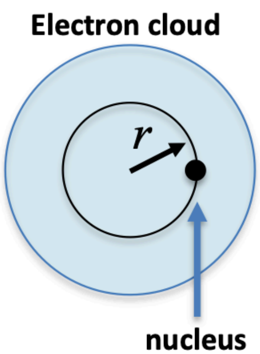

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harmonic restoring force

Now substitute $\vec{F}_{en} = -m\omega_0^2 \vec{x}$ into eq. for $\vec{x}$

$$\boxed{\frac{d^2}{dt^2}\vec{x} + \omega_0^2 \vec{x} = \frac{e}{m}\vec{E}(\vec{R},t)}$$

Combine with $\vec{P} = N\vec{p}$, $\vec{p} = e\vec{x}$ where N is the number density of atoms. This relates the macroscopic $\vec{P}$ to the microscopic $\vec{x}$

We now have

Maxwell's Equations
The Lorentz model

Maxwell-Lorentz Equations
We can seek self-consistent solutions to wave propagation

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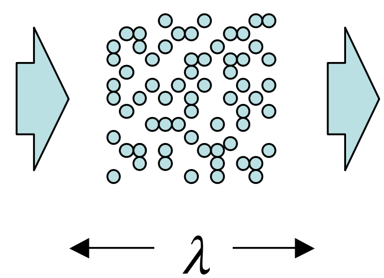
Classical Model of Absorption & Dispersion

Maxwell's Eqs: Oscillating dipole loses energy

⇨ Must include damping in Eq, of Motion

<u>Note</u>: In perfectly homogeneous media the coherently scattered light from a collection of Lorentz oscillators interferes constructively only in the forward direction ⇨



No energy loss for a propagating fields
(See note set "Classical Light-Matter")

QM to the rescue: Part of the radiation from quantum mechanical atoms is incoherent.

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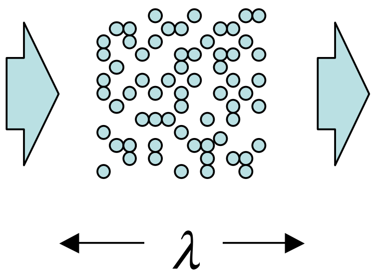

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QM to the rescue: Part of the radiation from quantum mechanical atoms is incoherent.

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The Lorentz Model with Damping

We add an ad hoc friction term w/ $\beta \ll \omega_0$

damping rate

This gives us our basic equation for the atomic response:

$$\boxed{\frac{d^2}{dt^2}\vec{x} + 2\beta\frac{d}{dt}\vec{x} + \omega_0^2\vec{x} = \frac{e}{m}\vec{E}(\vec{R},t)}$$

This type of differential equation generally has both oscillating and decaying terms. Solutions without source terms generally decay as $e^{-\beta t}$

⇨ We adopt a trial solution

Driving Field $\quad E(\vec{R},t) = \vec{\varepsilon} E_0 e^{-i(\omega t - kz)}$

Response $\quad \vec{x}(\vec{R},t) = \vec{a}\, e^{-i(\omega t - kz)}$

complex amplitude

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Solution for $\vec{a}$:

$$\boxed{\vec{a} = -\vec{\varepsilon}\,\frac{(e/m)E_0}{\omega^2 - \omega_0^2 + 2i\beta\omega}}$$

<u>Physical Quantities:</u>

Field

$$Re[\vec{E}(\vec{R},t)] = \vec{\varepsilon}E_0\cos(\omega t)$$

Dipole ($\vec{\varepsilon}$ real)

$$Re[\vec{p}(\vec{R},t)] = Re[e\vec{x}(\vec{R},t)]$$

$$= \vec{\varepsilon}E_0\frac{e^2}{m}\frac{(\omega_0^2-\omega^2)\cos(\omega t - kz) + 2\beta\omega\sin(\omega t - kz)}{(\omega_0^2-\omega^2) + 4\beta^2\omega^2}$$

<u>Note</u>: $\vec{p}$ and $\vec{E}$ generally oscillate out of phase

$\omega \ll \omega_0$ ➡ $\vec{p}$ & $\vec{E}$ in-phase

$\omega = \omega_0$ ➡ $\vec{p}$ lags $\vec{E}$ by $\pi/2$

$\omega \gg \omega_0$ ➡ $\vec{p}$ Lags $\vec{E}$ by π

Best to stick with complex notation !

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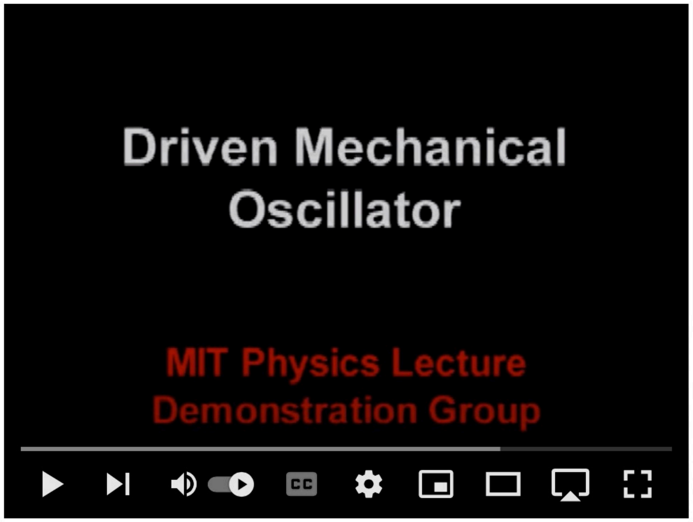
Video of driven – damped harmonic oscillator

https://www.youtube.com/watch?v=aZNnwQ8HJHU



Complex polarizability:

$$\vec{p} = e\vec{x} = e\vec{a}e^{-i(\omega t - kz)} \equiv \alpha(\omega)\vec{\varepsilon}E_0 e^{-i(\omega t - kz)}$$

$$\alpha(\omega) = \frac{e^2/m}{\omega_0^2 - \omega^2 - 2i\beta\omega} = \frac{e^2}{m}\frac{\omega_0^2 - \omega^2 + 2i\beta\omega}{(\omega_0^2 - \omega^2)^2 + 4\beta^2\omega^2}$$

Easy to show that if $\vec{E}(\vec{R},t) = \vec{\varepsilon}E_0 e^{-i(\omega t - kz)}$

and $\vec{P} = N\vec{p}$

then the wave equation reduces to

$$\left(-k^2 + \frac{\omega^2}{c^2}\right)\vec{\varepsilon}E_0 e^{-i(\omega t - kz)} = -\frac{\omega^2}{c^2}\frac{N\alpha(\omega)}{\varepsilon_0}\vec{\varepsilon}E_0 e^{-i(\omega t - kz)}$$

➡ plane wave solutions with $k = n(\omega)\omega/c$ where

$$\boxed{k^2 = \frac{\omega^2}{c^2}\left[1 + \frac{N\alpha(\omega)}{\varepsilon_0}\right] \equiv \frac{\omega^2}{c^2}n(\omega)}$$

Complex index of refraction

Light-Matter Interaction

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Complex index of refraction

Complex <u>Index of Refraction</u> – Physical discussion

Let $\boxed{n(\omega) = n_R(\omega) + i n_i(\omega)}$

Plane wave propagation $k = n(\omega)\omega/c$

$$\vec{E}(z,t) = \vec{\varepsilon} E_0 e^{-i(\omega t - kz)}$$

$$= \vec{\varepsilon} E_0 e^{-i(\omega t - [n(\omega)\omega/c]z)}$$

$$= \vec{\varepsilon} E_0 e^{-n_i(\omega)\omega z/c} e^{-i\omega(t - n_R(\omega)z/c)}$$

We can now identify

$\dfrac{c}{\omega n_i(\omega)}$	←	attenuation length
$\dfrac{c}{n_R(\omega)}$	←	phase velocity

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Absorption

The intensity of a plane wave field E is

$$I_\omega(z) = \frac{1}{2} n_R(\omega) c \varepsilon_0 |E(c,z)|^2 = I_0(0) e^{-2n_i(\omega)\omega z/c}$$

$$\equiv I_0 e^{-a(\omega)z}$$

where the absorption coefficient is

$$a(\omega) \equiv 2n_I(\omega)\omega/c = \frac{2\omega}{c} \mathrm{Im}\left[\left(1 + \frac{N\alpha(\omega)}{\varepsilon_0}\right)^{1/2}\right]$$

Possibility of gain ?

No – there is no energy source !

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Absorption and Dispersion in Gases

Approximations:

$\lvert\omega_0-\omega\rvert \ll \omega_0, \omega$	near resonance
$\lvert n(\omega)\rvert \sim 1$	weakly polarizable

Let $\omega_0^2-\omega^2=(\omega_0+\omega)(\omega_0-\omega)\approx 2\omega(\omega_0-\omega)$

$$\alpha(\omega)=\frac{e^2/m}{\omega_0^2-\omega^2-2i\beta\omega}=\frac{e^2/2m\omega}{\omega_0-\omega-i\beta}$$

$$=\frac{e^2}{2m\omega}\frac{\omega_0-\omega+i\beta}{(\omega_0-\omega)^2+\beta^2}$$

Furthermore

$$n(\omega)^2=1+\frac{N\alpha(\omega)}{\varepsilon_0}=1+\varepsilon,\ \varepsilon\ll 1$$

Expand to 1st order $(1+\varepsilon)^{1/2}\approx 1+\varepsilon/2$

Putting it together

$$n_R(\omega)=1+\frac{Ne^2}{4\varepsilon_0 m\omega}\frac{\omega_0-\omega}{(\omega_0-\omega)^2+\beta^2}$$

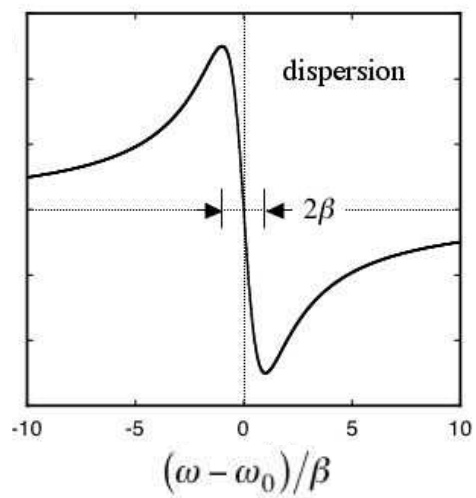
dispersive line shape

$$n_I(\omega)=\frac{Ne^2}{4\varepsilon_0 m\omega}\frac{\beta}{(\omega_0-\omega)^2+\beta^2}$$

Lorentzian line shape

General behavior:

$n_R(\omega)$



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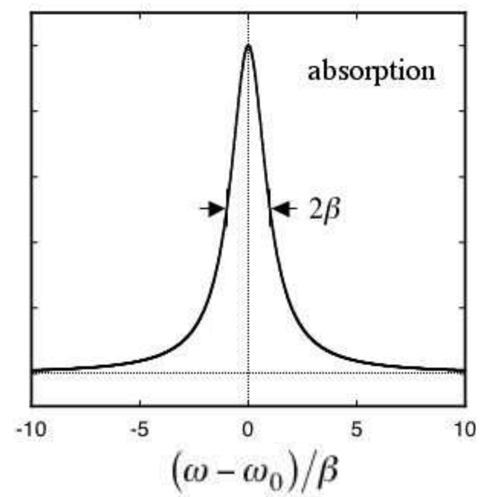

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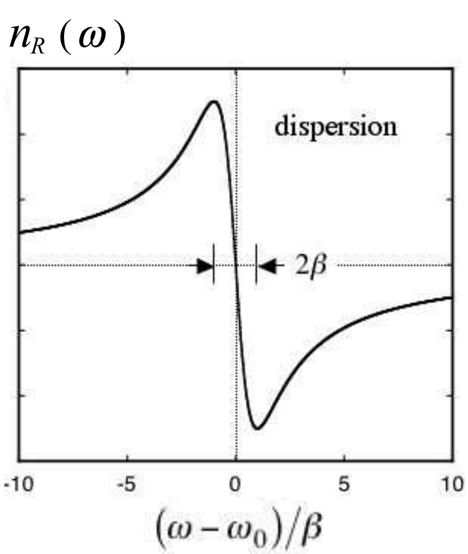

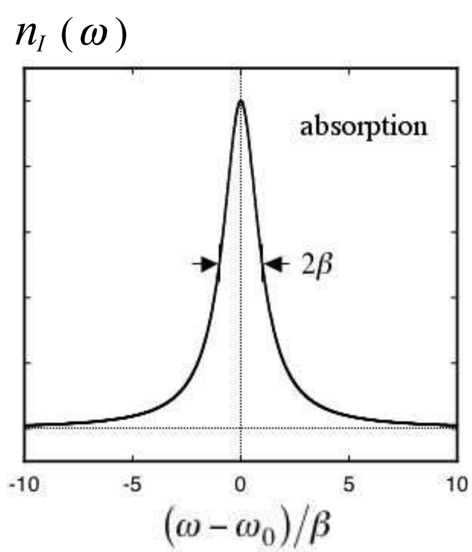

Note: for $|\omega_0-\omega| \gg \beta$

dispersion $\propto \dfrac{1}{(\omega_0-\omega)}$

absorption $\propto \dfrac{1}{(\omega_0-\omega)^2}$

➡ We can have loss-less dispersive media

Note: If we introduce the detuning $\Delta = (\omega_0-\omega)$ we can rewrite $n_R(\omega)$, $n_I(\omega)$ as

$$n_R(\Delta) = 1 + \frac{Ne^2}{4\varepsilon_0 m\omega}\,\frac{\Delta}{\Delta^2+\beta^2}$$

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From the above we see that

$n_R(\omega) < 1$ for $\omega > \omega_0$ ➡ $\dfrac{c}{n_R(\omega)} \gg c$

Superluminal propagation ?

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Superluminal propagation ?

Free Electrons

Consider the limit $\omega \gg \omega_0$

➡ effectively unbound electrons

This is a reasonable model of plasmas & metals

In this limit we have

$$\alpha(\omega) = \frac{e^2/m}{\omega_0^2-\omega^2-2i\beta\omega} \approx -\frac{e^2}{m\omega}$$ ➡

$$n(\omega) = \sqrt{1+\frac{N(\alpha)}{\varepsilon_0}} \approx \sqrt{1-\frac{Ne^2}{\varepsilon_0 m\omega^2}} \equiv \sqrt{1-\frac{\omega_p^2}{\omega^2}}$$

We introduce the <u>Plasma Frequency</u>:

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We introduce the **Plasma Frequency** :

$$\boxed{\omega_p = \sqrt{\frac{Ne^2}{\varepsilon_0 m}}}$$

Let $\left.\begin{matrix} \omega_0 \ll \omega \ll \omega_p \\ |\omega_0 - \omega| \gg \beta \end{matrix}\right\}$ ➡ $n(\omega)$ purely imaginary - but no loss!

We now have

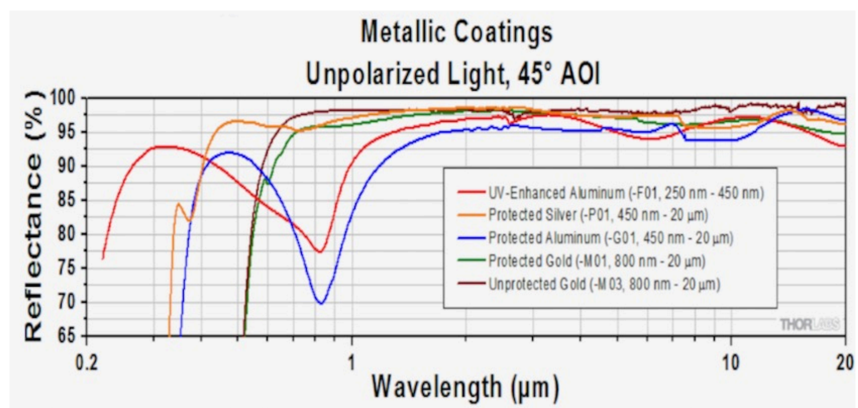
$$\vec{E}(z,t) = \vec{\varepsilon} E_0 e^{-i\omega[t - n(\omega)z/c]}$$

$$= \vec{\varepsilon} E_0 e^{-i\omega t} e^{i(z/c)\sqrt{\omega^2 - \omega_p^2}}$$

$$= \vec{\varepsilon} E_0 e^{-i\omega t} e^{-b(\omega) z}$$

where $b(\omega) = -\frac{i}{c}\sqrt{\omega^2 - \omega_p^2}$

Reflection at surface, penetration depth $\sim 1/b(\omega)$



Examples of this kind of medium includes plasmas, and metals such as aluminum, silver and gold which are known to be excellent mirrors for visible and IR radiation.