

Wigner-Weisskopf Theory of Spontaneous Decay

Today: Decay of atomic excited state due to interaction with the quantum electromagnetic field

Setup

Hamiltonian: (Schrödinger Picture)

$$\hat{H} = \hat{H}_F + \hat{H}_A + \hat{H}_{AF} =$$

$$\sum_{\vec{k}} \hbar \omega_{\vec{k}} \hat{a}_{\vec{k}}^+ \hat{a}_{\vec{k}} + \frac{1}{2} \hbar \omega_{21} \hat{\sigma}_z + \sum_{\vec{k}} \hbar (g_{\vec{k}} \hat{\sigma}_+ \hat{a}_{\vec{k}} + g_{\vec{k}}^* \hat{\sigma}_- \hat{a}_{\vec{k}}^+)$$

Hamiltonian: (Interaction Pict., Res. Approx.)

$$\hat{H}_I(t) = \sum_{\vec{k}, \lambda} \hbar g_{\vec{k}, \lambda} \hat{\sigma}_+ \hat{a}_{\vec{k}, \lambda} e^{i(\omega_{21} - \omega_{\vec{k}})t} + \text{H.c.}$$

S. E.: $i\hbar \frac{d}{dt} |\psi(t)\rangle = \hat{H}_I(t) |\psi(t)\rangle$

Expand

$$|\psi(t)\rangle = c_{2,0}(t) |2,0\rangle + \sum_{\vec{k}, \lambda} c_{1,1_{\vec{k}, \lambda}}(t) |1, 1_{\vec{k}, \lambda}\rangle$$



$$\dot{c}_{2,0}(t) = -i \sum_{\vec{k}, \lambda} g_{\vec{k}, \lambda} e^{i(\omega_{21} - \omega_{\vec{k}})t} c_{1,1_{\vec{k}, \lambda}}(t)$$

$$\dot{c}_{1,1_{\vec{k}, \lambda}}(t) = -i g_{\vec{k}, \lambda}^* e^{-i(\omega_{21} - \omega_{\vec{k}})t} c_{2,0}(t)$$

infinite # of these

Formal Solution:

$$c_{1,1_{\vec{k}, \lambda}}(t) = -i g_{\vec{k}, \lambda}^* \int_0^t e^{-i(\omega_{21} - \omega_{\vec{k}})t'} c_{2,0}(t') dt'$$

Wigner-Weisskopf Theory of Spontaneous Decay

Expand

$$|2(t)\rangle = c_{2,0}(t)|2,0\rangle + \sum_{\vec{k},\lambda} c_{1,1_{\vec{k},\lambda}}(t)|1,1_{\vec{k},\lambda}\rangle$$



$$\dot{c}_{2,0}(t) = -i \sum_{\vec{k},\lambda} g_{\vec{k},\lambda} e^{i(\omega_{eg} - \omega_{\vec{k}})t} c_{1,1_{\vec{k},\lambda}}(t)$$

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$$\dot{c}_{2,0}(t) = -\sum_{\vec{k},\lambda} |g_{\vec{k},\lambda}|^2 \int_0^t e^{i(\omega_{eg} - \omega_{\vec{k}})(t-t')} c_{2,0}(t') dt'$$

Time Dep. Perturbation Theory:

Study short-time limit, $c_{2,0}(t) \sim 1$

➡ **Fermi's Golden Rule**

No good for this problem!

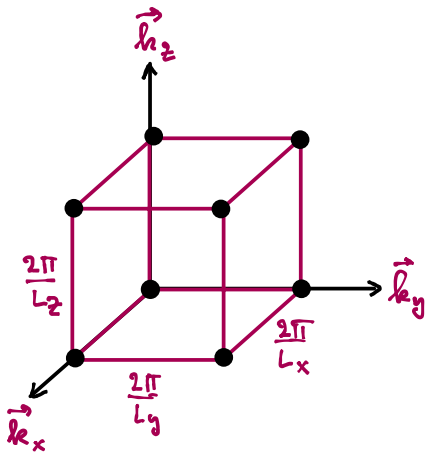
Wigner-Weisskopf Theory of Spontaneous Decay

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Quantization in a box w/periodic B. C.

$$\Rightarrow k_i = n_i \frac{2\pi}{L}, \quad n \text{ integer}$$

In \vec{k} space the modes form a grid:



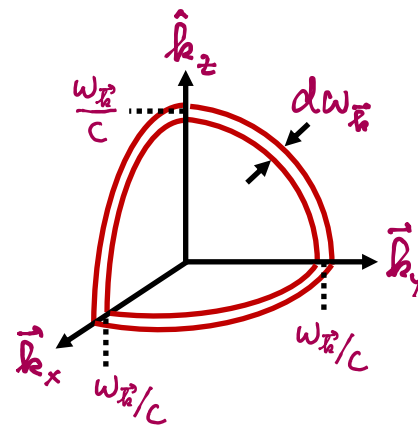
$$\frac{1 \text{ mode}}{\left(\frac{2\pi}{L_x}\right)\left(\frac{2\pi}{L_y}\right)\left(\frac{2\pi}{L_z}\right)} = \frac{V}{(2\pi)^3} = \mathcal{D}(\vec{k})$$

↑
Density of Modes

Convert sum to integral over modes:

$$\begin{aligned} \sum_{\vec{k}} &\rightarrow \int d^3 \hat{k} \mathcal{D}(\vec{k}) = \int k^2 d(\hat{k}) dk \mathcal{D}(\vec{k}) \\ &= \int d(\hat{k}) d\omega_{\vec{k}} \frac{\omega_{\vec{k}}^2}{c^3} \mathcal{D}(\vec{k}) = \int d(\hat{k}) d\omega_{\vec{k}} \mathcal{D}(\omega_{\vec{k}}) \end{aligned}$$

where $\hat{k} = \vec{k}/k$ is a unit vector along \vec{k} and



$$\mathcal{D}(\omega_{\vec{k}}) = \frac{V}{(2\pi)^3} \frac{\omega_{\vec{k}}^2}{c^3}$$

mode density in shell of \vec{k} - space of radius $\omega_{\vec{k}}/c$

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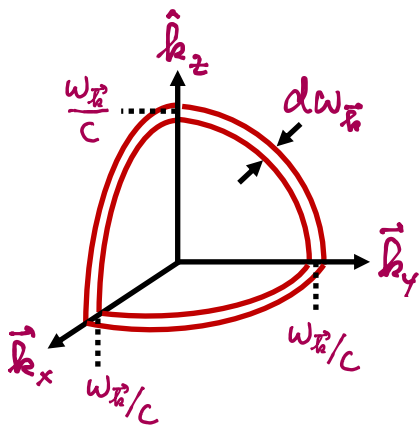
Thus, in the Continuum Limit

$$\sum_{\vec{k}, \lambda} |g_{\vec{k}, \lambda}|^2 \rightarrow \int_0^\infty d\omega_{\vec{k}} \mathcal{D}(\omega_{\vec{k}}) \sum_{\lambda} |d(\hat{k})|^2 |g_{\vec{k}, \lambda}|^2$$

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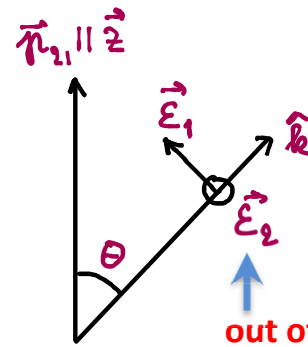


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We define the "polarization average"

$$\begin{aligned} \overline{|g(\omega_{\vec{k}})|^2} &= \sum_{\lambda} |d(\hat{k})|^2 |g_{\vec{k}, \lambda}|^2 = \\ &= \frac{1}{k^2} \left(\frac{\hbar \omega_{\vec{k}}}{2 \epsilon_0 V} \right) |d(\hat{k})|^2 \sum_{\lambda} |\vec{p}_{2,1} \cdot \vec{\epsilon}_{\vec{k}, \lambda}|^2 \end{aligned}$$



in polar coordinates

$$\sum_{\lambda} |\vec{p}_{2,1} \cdot \vec{\epsilon}_{\vec{k}, \lambda}|^2 = \sin^2 \theta |\vec{p}_{2,1}|^2$$

no ϕ dependence

$$\begin{aligned} \int d(\hat{k}) \sum_{\lambda} |\vec{p}_{2,1} \cdot \vec{\epsilon}_{\vec{k}, \lambda}|^2 &= \int_0^{2\pi} d\phi \int_0^\pi d(\cos \theta) \sin^2 \theta |\vec{p}_{2,1}|^2 \\ &= 2\pi |\vec{p}_{2,1}|^2 \int_{-1}^1 du (1-u^2) = \frac{8\pi}{3} |\vec{p}_{2,1}|^2 \end{aligned}$$

Wigner-Weisskopf Theory of Spontaneous Decay

$$\dot{c}_{2,0}(t) = - \sum_{\vec{k}, \lambda} |g_{\vec{k}, \lambda}|^2 \int_0^t e^{i(\omega_{21} - \omega_{\vec{k}})(t-t')} c_{2,0}(t') dt'$$

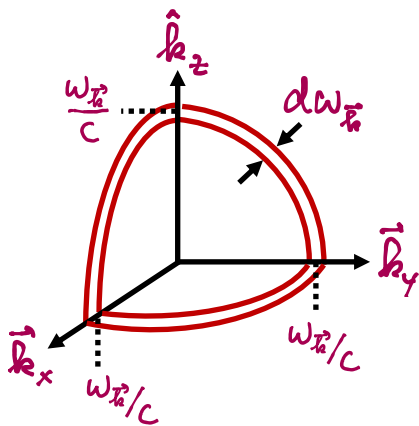
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Putting it together:

$$\begin{aligned} \dot{c}_{2,0}(t) &= - \sum_{\vec{k}, \lambda} |g_{\vec{k}, \lambda}|^2 \int_0^t e^{i(\omega_{21} - \omega_{\vec{k}})(t-t')} c_{2,0}(t') dt' \\ &= - \int_0^\infty d\omega_{\vec{k}} \overline{|g(\omega_{\vec{k}})|^2} \mathcal{D}(\omega_{\vec{k}}) \int_0^t dt' e^{i(\omega_{21} - \omega_{\vec{k}})(t-t')} c_{2,0}(t') \end{aligned}$$

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Thus, in the Continuum Limit

$$\sum_{\vec{k}, \lambda} |g_{\vec{k}, \lambda}|^2 \rightarrow \int_0^\infty d\omega_{\vec{k}} \mathcal{D}(\omega_{\vec{k}}) \sum_{\lambda} \int d(\hat{n}) |g_{\vec{k}, \lambda}|^2$$

We define the "polarization average"

$$\begin{aligned} \overline{|g(\omega_{\vec{k}})|^2} &= \sum_{\lambda} \int d(\hat{n}) |g_{\vec{k}, \lambda}|^2 = \\ &= \frac{1}{4\pi} \left(\frac{\hbar \omega_{\vec{k}}}{2 \epsilon_0 V} \right) \int d(\hat{n}) \sum_{\lambda} |\vec{e}_{\vec{k}, \lambda} \cdot \vec{e}_{\vec{n}}|^2 \\ &= \frac{4\pi \omega_{\vec{k}}^2}{2\hbar \epsilon_0 V} |\vec{e}_{\vec{k}, \lambda}|^2 \end{aligned}$$

Putting it together:

$$\begin{aligned} \dot{c}_{2,0}(t) &= - \sum_{\vec{k}, \lambda} |g_{\vec{k}, \lambda}|^2 \int_0^t e^{i(\omega_{2,1} - \omega_{\vec{k}})(t-t')} c_{2,0}(t') dt' \\ &= - \int_0^\infty d\omega_{\vec{k}} \overline{|g(\omega_{\vec{k}})|^2} \mathcal{D}(\omega_{\vec{k}}) \int_0^t dt' e^{i(\omega_{2,1} - \omega_{\vec{k}})(t-t')} c_{2,0}(t') \end{aligned}$$


Regular Perturbation Theory

Validity is limited to short times, such that the probability amplitude of the initial state does not change much. That allows us to approximate $c_{2,0}(t') \sim c_{2,0}(0)$. That does not work here.

Wigner-Weisskopf approximation:

The integral over time reduces to

$$\int_0^t dt' e^{i(\omega_{2,1} - \omega_{\vec{k}})(t-t')} \propto \delta(t-t')$$

on time scales $t-t' \gg |\omega_{2,1} - \omega_{\vec{k}}|^{-1}$ 

$$\begin{aligned} \dot{c}_{2,0}(t) &= - \sum_{\vec{k}, \lambda} |g_{\vec{k}, \lambda}|^2 \int_0^t e^{i(\omega_{2,1} - \omega_{\vec{k}})(t-t')} c_{2,0}(t') dt' \\ &= - \int_0^\infty d\omega_{\vec{k}} \overline{|g(\omega_{\vec{k}})|^2} \mathcal{D}(\omega_{\vec{k}}) \int_0^t dt' e^{i(\omega_{2,1} - \omega_{\vec{k}})(t-t')} \boxed{c_{2,0}(t)} \end{aligned}$$

Wigner-Weisskopf Theory of Spontaneous Decay


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Wigner-Weisskopf approximation:

The integral over time reduces to


$$\int_0^t dt' e^{i(\omega_{21} - \omega_{\mathbb{R}})(t-t')} \propto \delta(t-t')$$

on time scales $t-t' \gg |\omega_{21} - \omega_{\mathbb{R}}|^{-1}$ 

$$\begin{aligned} \dot{c}_{2,0}(t) &= -\sum_{\mathbb{R}, \lambda} |g_{\mathbb{R}, \lambda}|^2 \int_0^t e^{i(\omega_{21} - \omega_{\mathbb{R}})(t-t')} c_{2,0}(t') dt' \\ &= -\int_0^\infty d\omega_{\mathbb{R}} \overline{|g(\omega_{\mathbb{R}})|^2} \mathcal{D}(\omega_{\mathbb{R}}) \int_0^t e^{i(\omega_{21} - \omega_{\mathbb{R}})(t-t')} c_{2,0}(t') dt' \end{aligned}$$

Rest of Lecture: Work with this Equation

First: This eq. is of the form $\dot{c}_{2,0}(t) = \beta c_{2,0}(t)$

 { solutions oscillate at freq. $\text{Im}[\beta]$
and grow or decay at rate $\text{Re}[\beta]$

Next: Atoms couple weakly to the vacuum

$c_{2,0}(t)$ changes slowly on timescale ω_{21}^{-1} ,
evolving at a rate Γ to be determined.

That means we can let $t \rightarrow \infty$ on timescale ω_{21}^{-1}
while still keeping $t \ll \Gamma^{-1}$

Defining $-i\mathfrak{Z}(\omega_{\mathbb{R}} - \omega_{21}) = \int_0^\infty dt' e^{i(\omega_{21} - \omega_{\mathbb{R}})(t-t')}$

We can then rewrite

$$\dot{c}_{2,0}(t) = -\int_0^\infty d\omega_{\mathbb{R}} |g(\omega_{\mathbb{R}})|^2 \mathcal{D}(\omega_{\mathbb{R}}) [-i\mathfrak{Z}(\omega_{\mathbb{R}} - \omega_{21})] c_{2,0}(t)$$

Wigner-Weisskopf Theory of Spontaneous Decay

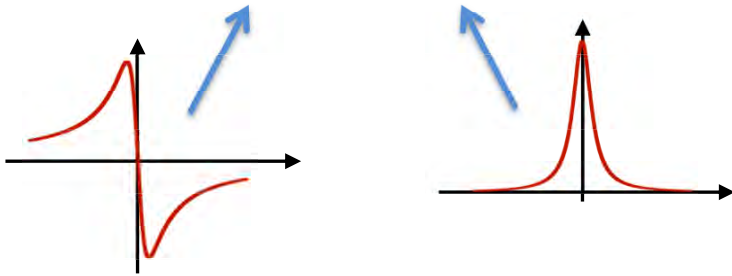
Math Details: Consider the function

$$\xi(\Omega) = i \int_0^{\infty} d\tau e^{-i\Omega\tau}, \quad \Omega = \omega_{\vec{k}} - \omega_{21}$$

Approximate the integral as

$$\xi(\Omega) = \lim_{\epsilon \rightarrow 0^+} i \int_0^{\infty} d\tau e^{-i\Omega\tau - \epsilon\tau} = \lim_{\epsilon \rightarrow 0^+} \left(\frac{i}{-i\Omega - \epsilon} \right)$$

$$= \lim_{\epsilon \rightarrow 0^+} \left(\frac{-\Omega^2}{\Omega^2 + \epsilon^2} + i \frac{\epsilon}{\Omega^2 + \epsilon^2} \right)$$



$$= \mathcal{P}\left(\frac{i}{\Omega}\right) + i\pi\delta(\Omega)$$

Cauchy's Principal Part

Thus,

$$\dot{c}_{2,0}(t) \approx - \left[\int_0^{\infty} d\omega_{\vec{k}} |g(\omega_{\vec{k}})|^2 \mathcal{D}(\omega_{\vec{k}}) \left(\pi \delta(\omega_{21} - \omega_{\vec{k}}) - i \mathcal{P}\left(\frac{1}{\omega_{21} - \omega_{\vec{k}}}\right) \right) \right] c_{2,0}(t)$$

which is of the form

$$\begin{aligned} \dot{c}_{2,0} &= - \left(\frac{A_{21}}{2} - i\delta \right) c_{2,0}(t) \rightarrow \\ c_{2,0}(t) &= e^{-A_{21}/2 t} e^{i\delta(t)} c_{2,0}(0) \end{aligned}$$

Decay Rate:

$$\begin{aligned} A_{21} &= - \int_0^{\infty} d\omega_{\vec{k}} 2\pi |g(\omega_{\vec{k}})|^2 \mathcal{D}(\omega_{\vec{k}}) \delta(\omega_{21} - \omega_{\vec{k}}) \\ &= 2\pi |g(\omega_{21})|^2 \mathcal{D}(\omega_{21}) \\ &= 2\pi \times \frac{4\pi\omega_{21}}{3\hbar\epsilon_0 V} |\vec{p}_{21}|^2 \times \frac{V}{(2\pi)^3} \frac{\omega_{21}^2}{c^3} \end{aligned}$$

Wigner-Weisskopf Theory of Spontaneous Decay

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$$A_{21} = \frac{\omega_{21}^3 |\vec{p}_{21}|^2}{3\pi \epsilon_0 \hbar c^3} = \frac{8\pi^2 |\vec{p}_{21}|^2}{3\pi \hbar \lambda^3}$$

This is the value used earlier!

Frequency shift:

$$\delta = \lim_{\epsilon \rightarrow 0} \left(i \int_0^\infty d\omega_{\vec{k}} |g(\omega_{\vec{k}})|^2 \mathcal{D}(\omega_{\vec{k}}) \frac{\omega_{21} - \omega_{\vec{k}}}{(\omega_{21} - \omega_{\vec{k}})^2 + \epsilon^2} \right)$$

Integral is not well behaved!

Further development by Willis Lamb



Lamb shift

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Lamb shift

Hydrogen 2P lifetime:

- accurate calculations for

$$\lambda = 121.57 \text{ nm}$$

$$|\vec{p}_{21}| = 0.745 a_0 e = \frac{4\pi\epsilon_0 \hbar^2}{m_e e^2} \times 0.745 e$$



$$(A_{21})^{-1} = \tau_{2p} = 1.595 \text{ ns}$$

$$\text{Measured: } \tau_{2p} = 1.60 \text{ ns}$$