

Density Matrix Description of 2-Level Atoms

More about the Density Matrix

Let $|\psi\rangle = \sum_j c_j |\mu_j\rangle$, where $\{|\mu_j\rangle\}$ is a basis and the index k labels members of the ensemble

Populations:

(real-valued)

$$\rho_{nn} = \langle c_n^{(k)} c_n^{(k)*} \rangle_k = \langle |c_n^{(k)}|^2 \rangle_k$$

Bayesianist:

Single system: ρ_{nn} is prob of observing the state $|\mu_n\rangle$

Frequentist:

Ensemble: The state $|\mu_n\rangle$ occurs with frequency ρ_{nn}

Coherences:

(complex-valued)

$$\rho_{np} = \langle c_n^{(k)} c_p^{(k)*} \rangle_k \quad \rho_{pn} = \rho_{np}^*$$

Note: Defining $c_q = |c_q| e^{i\theta_q}$ we have

$$|\langle c_n^{(k)} c_p^{(k)*} \rangle_k| = |\langle |c_n^{(k)}| |c_p^{(k)}| e^{i(\theta_n^{(k)} - \theta_p^{(k)})} \rangle_k| \leq \langle |c_n^{(k)}| |c_p^{(k)}| \rangle_k$$

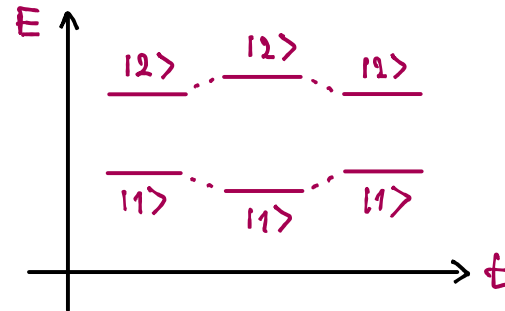
It follows that

$$\rho_{np} \rho_{pn} \leq \rho_{nn} \rho_{pp}$$

with = for pure states

$$\rho = \begin{pmatrix} \rho_{11} & \dots & \rho_{1n} \\ \vdots & \ddots & \vdots \\ \rho_{n1} & \dots & \rho_{nn} \end{pmatrix}$$

Example: 2-level atom w/random perturbations



Perturbing events cause random phase shifts $e^{i\theta}$ between states.

The ensemble average $\rho_{np} = \langle c_n c_p^* e^{i\theta} \rangle_k$ is reduced by the randomly fluctuating phase

Dipole Radiation:

$$\begin{aligned} \langle \hat{n} \rangle &= \text{Tr} [\rho \hat{n}] = \text{Tr} \left[\begin{pmatrix} \rho_{11} & \rho_{12} \\ \rho_{21} & \rho_{22} \end{pmatrix} \begin{pmatrix} 0 & \vec{n}_{12} \\ \vec{n}_{21} & 0 \end{pmatrix} \right] \\ &= \rho_{12} \vec{n}_{21} + \rho_{21} \vec{n}_{12} = 2 \text{Re} [\rho_{12} \vec{n}_{21}] \end{aligned}$$

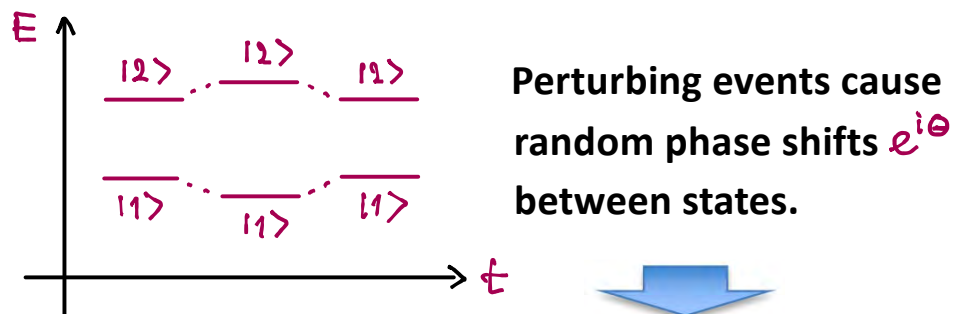
For an ensemble of pure states w/different θ

$$\langle \hat{n} \rangle = 2 \sum_k \rho_k \underbrace{\text{Re} [\rho_{12}^{(k)} \vec{n}_{21}]}_{\text{phase}} \vec{n}_{21}$$

Oscillating dipole w/phase that varies between atoms with different perturbation history

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$$\begin{aligned} \langle \hat{\vec{p}} \rangle &= \text{Tr} [\hat{\rho} \hat{\vec{p}}] = \text{Tr} \left[\begin{pmatrix} \rho_{11} & \rho_{12} \\ \rho_{21} & \rho_{22} \end{pmatrix} \begin{pmatrix} 0 & \vec{p}_{12} \\ \vec{p}_{21} & 0 \end{pmatrix} \right] \\ &= \rho_{12} \vec{p}_{21} + \rho_{21} \vec{p}_{12} = 2 \text{Re} [\rho_{12} \vec{p}_{21}] \end{aligned}$$

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Time Evolution of the Density Matrix

Challenge: We need "equations of motion" that combine the Schrödinger Equation with the effect of processes that can change $\text{Tr} \rho^2$ (measure of purity)

Approach: We do not have time for a rigorous derivation, so will rely on plausible arguments to justify the equations

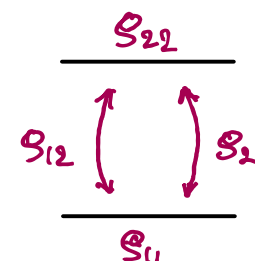
Schrödinger Evolution: In general, we have

$$\dot{\hat{\rho}} = -\frac{i}{\hbar} [\hat{H}, \hat{\rho}] = -\frac{i}{\hbar} (\hat{H}\hat{\rho} - \hat{\rho}\hat{H})$$

matrix elements

$$\dot{\rho}_{nm} = -\frac{i}{\hbar} \sum_{k=1,2} (H_{nk} \rho_{km} - \rho_{nk} H_{km})$$

2-Level Atom \Rightarrow $\begin{cases} 2 \text{ populations} \\ 2 \text{ coherences} \end{cases}$



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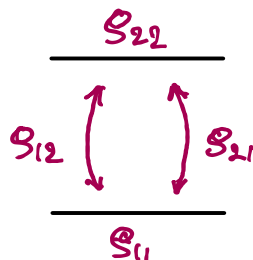
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$$\dot{\rho}_{nm} = -\frac{i}{\hbar} \sum_{k=1,2} (H_{nk} \rho_{km} - \rho_{nk} H_{km}) \quad (*)$$

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Consider the 2-Level Rabi problem with

$$H = H_0 + V \quad \& \quad V_{12} = \frac{1}{2} \hbar X_{12} e^{-i\omega t} + \text{c.c.}$$

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Set $X_{12} = X$, $X_{21} = X^*$, substitute $\rho_{12} = \tilde{\rho}_{12} e^{i\omega t}$
 (Pure state $\Rightarrow \rho_{12} = a_1 a_2^* = c_1 (c_2 e^{-i\omega t})$)
slow variable

Substitute in (*), make RWA, and drop \sim

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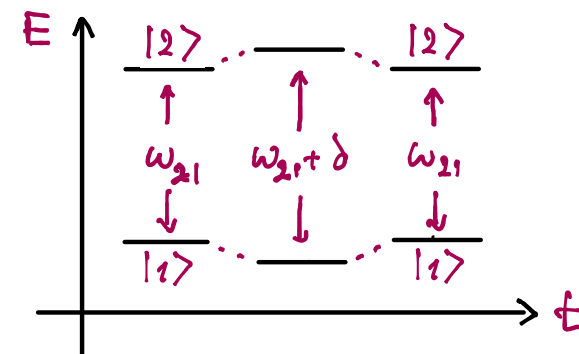
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Two atoms near each other \Rightarrow energy levels shift, free evol. of S_{12} changed



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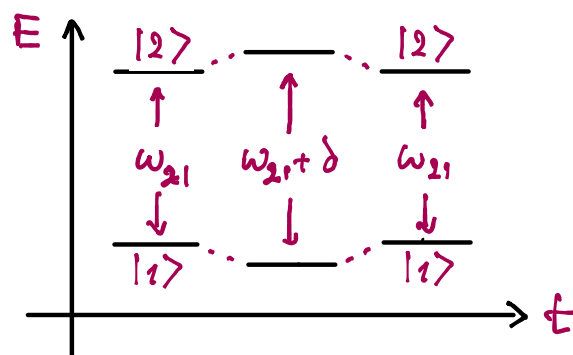
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$$\dot{\rho}_{12} = -i[\omega_{21} + \delta\omega(t)]\rho_{12}$$

collisional history \downarrow

$$\Rightarrow \rho_{12}(t) = \rho_{12}(0) e^{-i\omega_{21}t} e^{-i\int_0^t dt' \delta\omega(t')}$$

We need the ensemble average of $\rho_{12}(t)$

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It follows that: $\rho_{12}(t) = \rho_{12}(0) e^{-i\omega_{21}t} e^{-t/\tau}$

Add this decay to the equation of motion to get

$$\dot{\rho}_{12} = (\dot{\rho}_{12})_{S.E.} + (\dot{\rho}_{12})_{E.C.} = -(i\omega_{21} - 1/\tau)\rho_{12}$$

Simple Model of Inelastic Collisions

As modeled by, e. g., Milloni & Eberly,
this is a steady loss of atoms



$$\begin{aligned} \dot{\rho}_{11} &= (\dot{\rho}_{11})_{S.E.} - \Gamma_1 \rho_{11} \\ \dot{\rho}_{22} &= (\dot{\rho}_{22})_{S.E.} - \Gamma_2 \rho_{22} \end{aligned}$$

_____ |2>

 ↘ Γ₂

_____ |1>

 ↘ Γ₁

This is strange because $\text{Tr} \rho(t)$ is not preserved

Convenient when working with quantities

$$N \langle \vec{n} \rangle \propto N (\vec{n}_{12} \rho_{11} + \vec{n}_{21} \rho_{12})$$

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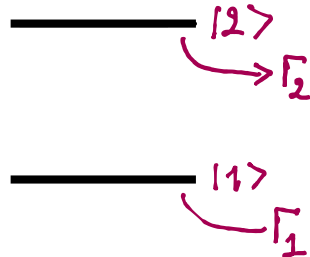
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$$N \langle \vec{p} \rangle \propto N (\vec{p}_{12} \rho_{11} + \vec{p}_{21} \rho_{12})$$

Effect on probability amplitudes

Populations are ensemble averages of the type

$$\rho_{11}(t) = \langle |a_1(t)|^2 \rangle = \langle |a_1(0)|^2 \rangle e^{-\Gamma_1 t}$$

$$\rho_{22}(t) = \langle |a_2(t)|^2 \rangle = \langle |a_2(0)|^2 \rangle e^{-\Gamma_2 t}$$

When the populations decay, the averages of the probability amplitudes must decay accordingly,


$$\langle |a_1(t)| \rangle = \langle |a_1(0)| \rangle e^{-\Gamma_1/2 t}$$

$$\langle |a_2(t)| \rangle = \langle |a_2(0)| \rangle e^{-\Gamma_2/2 t}$$

Thus, for the coherences

$$\rho_{12}(t) = \langle a_1(t) a_2(t)^* \rangle = \langle a_1(0) a_2(0)^* \rangle e^{-\Gamma_1/2 t} e^{-\Gamma_2/2 t}$$

This gives us



$$\dot{\rho}_{12} = (\dot{\rho}_{12})_{S.E.} - 1/\tau \rho_{12} - \frac{\Gamma_1 + \Gamma_2}{2} \rho_{12}$$

Density Matrix Description of 2-Level Atoms

Effect on probability amplitudes

Populations are ensemble averages of the type

$$S_{11}(t) = \langle |a_1(t)|^2 \rangle = \langle |a_1(0)|^2 \rangle e^{-\Gamma_1 t}$$

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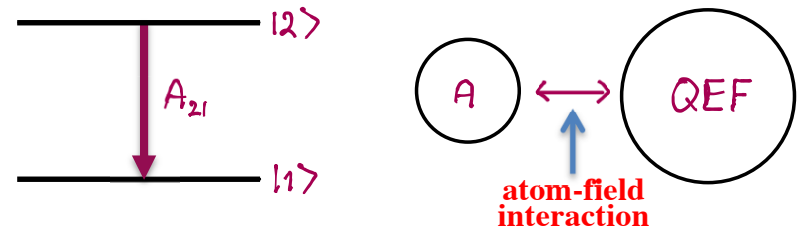
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$$\dot{S}_{12} = (\dot{S}_{12})_{\text{S.E.}} - \frac{1}{\tau} S_{12} - \frac{\Gamma_1 + \Gamma_2}{2} S_{12}$$

elastic
inelastic

Spontaneous Decay

This process occurs due to interaction between the Atom and the Quantized Electromagnetic Field



Warm-up: A Bayesian recipe for Mixed States

Alice has two 2-level atoms in the ground state.

Step (1) She applies a Hamiltonian that drives the evolution

$$|1\rangle_A |1\rangle_B \rightarrow a_1 |1\rangle_A |1\rangle_B + a_2 |2\rangle_A |2\rangle_B$$

Step (2) She gives atom B to Bob and asks him to measure if it is in $|1\rangle_B$ or $|2\rangle_B$ and keep the result secret forever.

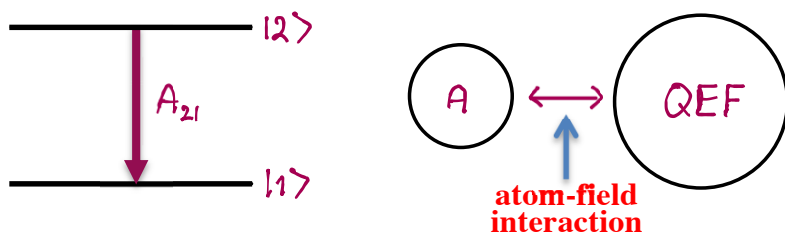
Result: Alice now has a 2-level atom in the state

$$S = |a_1|^2 |1\rangle_{BB} \langle 1| + |a_2|^2 |2\rangle_{BB} \langle 2|$$

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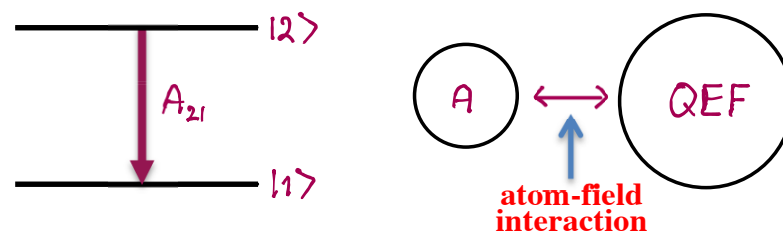
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$$\rho = |a_1|^2 |1\rangle_B \langle 1| + |a_2|^2 |2\rangle_B \langle 2|$$

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Final OPTI 544 Lectures:

$$|\psi(0)\rangle = |2\rangle_A |vac\rangle_{QEF} \quad \Rightarrow \quad \text{evolution over time } t$$

$$|\psi(t)\rangle = C_{2,0}(t) |2\rangle_A |vac\rangle_{QEF} + \sum_k C_{1,1k}(t) |1\rangle_A |n_k=1\rangle_{QEF}$$

↑
photon "in the atom"
↑
photon in field mode k

Cannot recover info in continuum of field modes

Probability $|C_{2,0}(t)|^2$ of having **no decay**

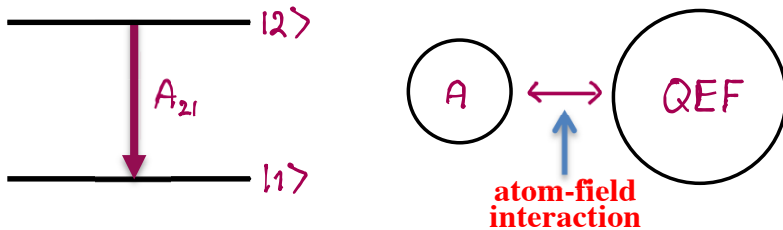
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Conclusion: Decay moves population $|2\rangle \rightarrow |1\rangle$
at rate A_{21} , damps coherence at rate $A_{21}/2$

$$\dot{\rho}_{11} = A_{21} \rho_{22}, \quad \dot{\rho}_{22} = -A_{21} \rho_{11}$$

$$\dot{\rho}_{12} = -\frac{A_{21}}{2} \rho_{12} = \dot{\rho}_{21}^*$$

Putting it all together:

$$\dot{\rho}_{11} = -\Gamma_1 \rho_{11} + A_{21} \rho_{22} - \frac{i}{2} (X \rho_{12} - X^* \rho_{21})$$

$$\dot{\rho}_{22} = -\Gamma_2 \rho_{22} - A_{21} \rho_{11} + \frac{i}{2} (X \rho_{12} - X^* \rho_{21})$$

$$\dot{\rho}_{12} = (i\Delta - \beta) \rho_{12} + \frac{iX^*}{2} (\rho_{22} - \rho_{11}) = \dot{\rho}_{21}^*$$

where $\beta = \frac{\Gamma_1}{2} + \frac{A_{21}}{2} + \frac{\Gamma_2}{2}$

These are our desired

**Density Matrix
Equations of Motion**

Density Matrix Description of 2-Level Atoms

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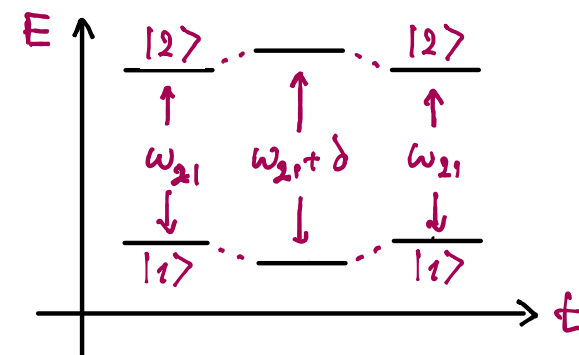
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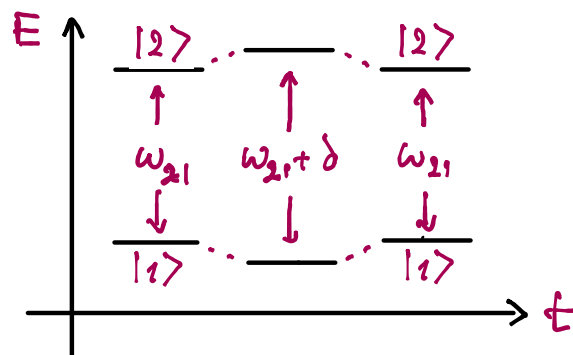
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collisional history \downarrow

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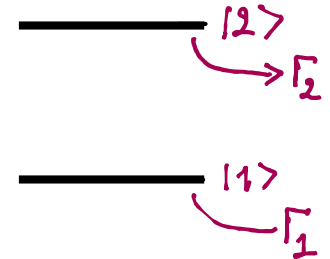
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$$N \langle \vec{p} \rangle \propto N (\vec{p}_{12} \rho_{21} + \vec{p}_{21} \rho_{12})$$

Density Matrix Description of 2-Level Atoms


It follows that: $\rho_{12}(t) = \rho_{12}(0) e^{-i\omega_{21}t} e^{-t/\tau}$

Add this decay to the equation of motion to get

$$\dot{\rho}_{12} = (\dot{\rho}_{12})_{S.E.} + (\dot{\rho}_{12})_{E.C.} = -(i\omega_{21} - 1/\tau)\rho_{12}$$

Simple Model of Inelastic Collisions

As modeled by, e. g., Milloni & Eberly, this is a steady loss of atoms



$$\begin{aligned}\dot{\rho}_{11} &= (\dot{\rho}_{11})_{S.E.} - \Gamma_1 \rho_{11} \\ \dot{\rho}_{22} &= (\dot{\rho}_{22})_{S.E.} - \Gamma_2 \rho_{22}\end{aligned}$$

This is strange because $\text{Tr} \rho(t)$ is not preserved

Convenient when working with quantities

$$N \langle \vec{p} \rangle \propto N (\vec{p}_{12} \rho_{11} + \vec{p}_{21} \rho_{12})$$

Effect on probability amplitudes

Populations are ensemble averages of the type

$$\rho_{11}(t) = \langle |a_1(t)|^2 \rangle = \langle |a_1(0)|^2 \rangle e^{-\Gamma_1 t}$$

$$\rho_{22}(t) = \langle |a_2(t)|^2 \rangle = \langle |a_2(0)|^2 \rangle e^{-\Gamma_2 t}$$

When the populations decay, the averages of the probability amplitudes must decay accordingly,

$$\langle |a_1(t)| \rangle = \langle |a_1(0)| \rangle e^{-\Gamma_1/2 t}$$

$$\langle |a_2(t)| \rangle = \langle |a_2(0)| \rangle e^{-\Gamma_2/2 t}$$

Thus, for the coherences

$$\rho_{12}(t) = \langle a_1(t) a_2(t)^* \rangle = \langle a_1(0) a_2(0)^* \rangle e^{-\Gamma_1/2 t} e^{-\Gamma_2/2 t}$$

This gives us

elastic

inelastic

$$\dot{\rho}_{12} = (\dot{\rho}_{12})_{S.E.} - 1/\tau \rho_{12} - \frac{\Gamma_1 + \Gamma_2}{2} \rho_{12}$$

Density Matrix Description of 2-Level Atoms

Effect on probability amplitudes

Populations are ensemble averages of the type

$$S_{11}(t) = \langle |a_1(t)|^2 \rangle = \langle |a_1(0)|^2 \rangle e^{-\Gamma_1 t}$$

$$S_{22}(t) = \langle |a_2(t)|^2 \rangle = \langle |a_2(0)|^2 \rangle e^{-\Gamma_2 t}$$

When the populations decay, the averages of the probability amplitudes must decay accordingly,

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Thus, for the coherences

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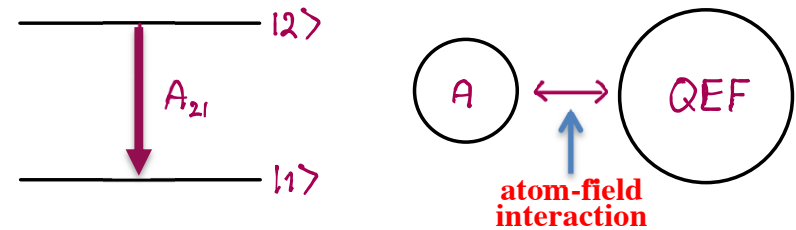
This gives us

$$\dot{S}_{12} = (\dot{S}_{12})_{S.E.} - \frac{1}{T} S_{12} - \frac{\Gamma_1 + \Gamma_2}{2} S_{12}$$

elastic
inelastic

Spontaneous Decay

This process occurs due to interaction between the Atom and the Quantized Electromagnetic Field



Warm-up: A Bayesian recipe for Mixed States

Alice has two 2-level atoms in the ground state.

Step (1) She applies a Hamiltonian that drives the evolution

$$|1\rangle_A |1\rangle_B \rightarrow a_1 |1\rangle_A |1\rangle_B + a_2 |2\rangle_A |2\rangle_B$$

Step (2) She gives atom B to Bob and asks him to measure if it is in $|1\rangle_B$ or $|2\rangle_B$ and keep the result secret forever.

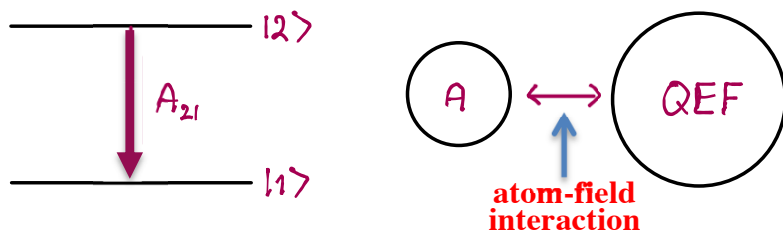
Result: Alice now has a 2-level atom in the state

$$S = |a_1|^2 |1\rangle_{BB} \langle 1| + |a_2|^2 |2\rangle_{BB} \langle 2|$$

Density Matrix Description of 2-Level Atoms

Spontaneous Decay

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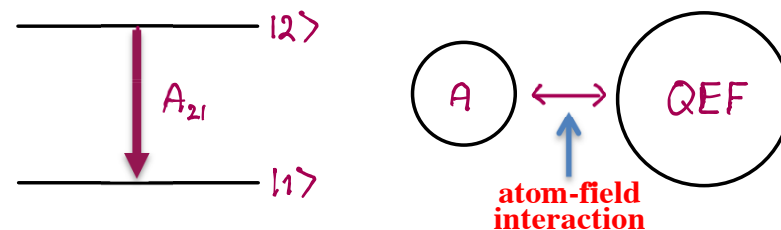
Step (2) She gives atom B to Bob and asks him to measure if it is in $|1\rangle_B$ or $|2\rangle_B$ and keep the result secret forever.

Result: Alice now has a 2-level atom in the state

$$\rho = |a_1|^2 |1\rangle_B \langle 1| + |a_2|^2 |2\rangle_B \langle 2|$$

Spontaneous Decay

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Final OPTI 544 Lectures:

$$|\psi(0)\rangle = |2\rangle_A |\text{vac}\rangle_{\text{QEF}} \quad \rightarrow \quad \text{evolution over time } t$$

$$|\psi(t)\rangle = C_{2,0}(t) |2\rangle_A |\text{vac}\rangle_{\text{QEF}} + \sum_k C_{1,1k}(t) |1\rangle_A |n_k=1\rangle_{\text{QEF}}$$

↑
↑
 photon "in the atom" photon in field mode k

Cannot recover info in continuum of field modes

Probability $|C_{2,0}(t)|^2$ of having **no decay**

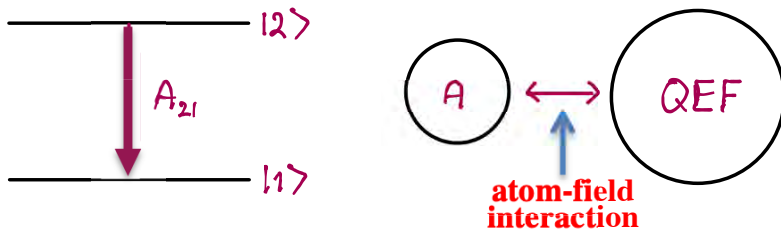
Probability $\sum_k |C_{1,1k}(t)|^2$ of having **decay**

No Coherence established between states $|1\rangle, |2\rangle$

Density Matrix Description of 2-Level Atoms

Spontaneous Decay

This process occurs due to interaction between the Atom and the Quantized Electromagnetic Field



Final OPTI 544 Lectures:

$|\psi(0)\rangle = |2\rangle_A |\text{vac}\rangle_{\text{QEF}}$ evolution over time t

$$|\psi(t)\rangle = c_{2,0}(t) |2\rangle_A |\text{vac}\rangle_{\text{QEF}} + \sum_k c_{1,1k}(t) |1\rangle_A |n_k=1\rangle_{\text{QEF}}$$

↑
excitation in the atom
↑
excitation in field mode k

Cannot recover info in continuum of field modes

Probability $|c_{2,0}(t)|^2$ of having **no decay**

Probability $\sum_k |c_{1,1k}(t)|^2$ of having **decay**

No Coherence established between states $|1\rangle, |2\rangle$

Conclusion: Decay moves population $|2\rangle \rightarrow |1\rangle$
at rate $A_{2,1}$, damps coherence at rate $A_{2,1}/2$

$$\dot{\rho}_{11} = A_{2,1} \rho_{22}, \quad \dot{\rho}_{22} = -A_{2,1} \rho_{22}$$

$$\dot{\rho}_{12} = -\frac{A_{2,1}}{2} \rho_{12} = \dot{\rho}_{21}^*$$

Putting it all together:

$$\dot{\rho}_{11} = -\Gamma_1 \rho_{11} + A_{2,1} \rho_{22} - \frac{i}{2} (X \rho_{12} - X^* \rho_{21})$$

$$\dot{\rho}_{22} = -\Gamma_2 \rho_{22} - A_{2,1} \rho_{22} + \frac{i}{2} (X \rho_{12} - X^* \rho_{21})$$

$$\dot{\rho}_{12} = (i\Delta - \beta) \rho_{12} + \frac{iX^*}{2} (\rho_{22} - \rho_{11}) = \dot{\rho}_{21}^*$$

where $\beta = \frac{\Gamma_1}{2} + \frac{A_{2,1}}{2} + \frac{\Gamma_2}{2}$

These are our desired

**Density Matrix
Equations of Motion**

Emission and Absorption – Population Rate Equations

Note: The terms remaining after elimination of S_{12}, S_{21} are commonly identified with induced or stimulated processes. They are proportional to $|X|^2, |E_0|^2$ and thus the intensity of the light field.

Def: Absorption Rate = Stimulated Emission Rate

$$R_{12} = \frac{|X|^2 \beta / 2}{\Delta^2 + \beta^2} = \frac{|\vec{\mu}_{12} \cdot \vec{E}_0 / \hbar|^2 \beta / 2}{(\omega_{21} - \omega)^2 + \beta^2}$$

Schematic:

