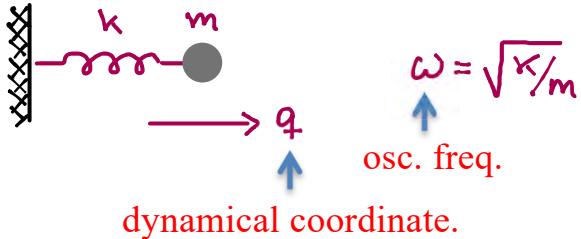


# Quantum Electrodynamics – Intro to Field Theory

## Classical Simple Harmonic Oscillator (SHO)

Particle on a spring



Kinetic Energy:

$$T = \frac{1}{2} m \dot{q}^2$$

Potential Energy:

$$V = \frac{1}{2} k q^2 = \frac{1}{2} m \omega^2 q^2$$

Lagrangian:  $\mathcal{L} = T - V = \frac{1}{2} m \dot{q}^2 - \frac{1}{2} m \omega^2 q^2$

$$\frac{\partial}{\partial t} \frac{\partial \mathcal{L}}{\partial \dot{q}} - \frac{\partial \mathcal{L}}{\partial q} = 0 \quad \Rightarrow \quad \ddot{q} + \omega^2 q = 0$$

usual eq. of motion

Conjugate momentum

$$p = \frac{\partial \mathcal{L}}{\partial \dot{q}} = m \dot{q}$$

Hamiltonian

$$\mathcal{H} = T(\dot{q} = p/m) + V(q) = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 q^2$$

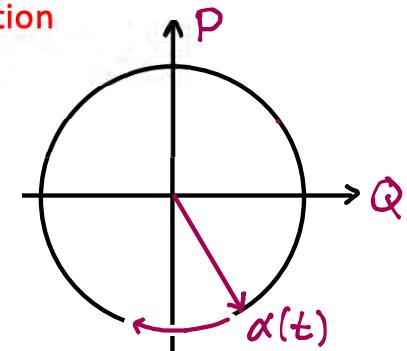
$$\left. \begin{aligned} \dot{q} &= \frac{\partial \mathcal{L}}{\partial p} = p/m \\ \dot{p} &= -\frac{\partial \mathcal{L}}{\partial q} = -m \omega^2 q \end{aligned} \right\} \quad \ddot{q} + \omega^2 q = 0$$

Phase plane

Scaled variables

$$Q \equiv q/q_0, \quad P \equiv p/p_0$$

$$\alpha = Q + iP \quad \left\{ \begin{array}{l} Q = \text{Re}[\alpha] \\ P = \text{Im}[\alpha] \\ \mathcal{H} = E_0 \alpha^* \alpha \end{array} \right.$$



# Quantum Electrodynamics – Intro to Field Theory

Conjugate momentum

$$p = \frac{\partial \mathcal{L}}{\partial \dot{q}} = m\dot{q}$$

Hamiltonian

$$\mathcal{H} = T(\dot{q} = p/m) + V(q) = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 q^2$$

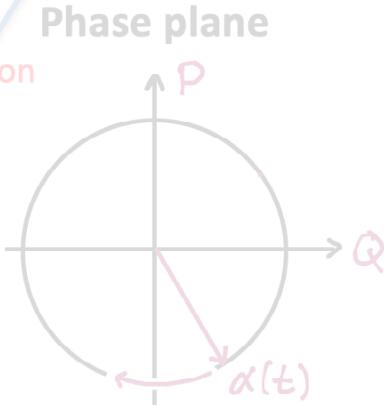
$$\begin{aligned}\dot{q} &= \frac{\partial \mathcal{L}}{\partial p} = p/m \\ \dot{p} &= -\frac{\partial \mathcal{L}}{\partial q} = -m\omega^2 q\end{aligned}$$

$$\left. \begin{aligned}\dot{q} &= p/m \\ \dot{p} &= -m\omega^2 q\end{aligned} \right\} \Rightarrow \ddot{q} + \omega^2 q = 0$$

Scaled variables

$$Q \equiv q/q_0, \quad P \equiv p/p_0$$

$$\alpha = Q + iP \quad \left\{ \begin{aligned} Q &= \text{Re}[\alpha] \\ P &= \text{Im}[\alpha] \\ \mathcal{L} &= E_0 \alpha^* \alpha\end{aligned} \right.$$



solution

## Quantum Harmonic Oscillator

Formal Quantization Procedure:

$$q \rightarrow \hat{q}, \quad p \rightarrow \hat{p}, \quad [\hat{q}, \hat{p}] = i\hbar$$

Choose  $E_0 = \hbar\omega \quad \Rightarrow \quad q_0 = \sqrt{\frac{2\hbar}{m\omega}}, \quad p_0 = \sqrt{2m\hbar\omega}$

natural scale

$$\alpha \rightarrow \hat{\alpha} = \hat{Q} + i\hat{P} = \sqrt{\frac{m\omega}{2\hbar}} \left( \hat{q} + i\frac{\hat{p}}{m\omega} \right)$$

$$[\hat{\alpha}, \hat{\alpha}^\dagger] = 1$$

Rewrite:

$$\hat{H} = \hbar\omega (\hat{Q}^2 + \hat{P}^2) = \hbar\omega (\hat{a}^\dagger \hat{a} + \frac{1}{2})$$

$$\hat{N} = \hat{a}^\dagger \hat{a} \quad (\text{number operator})$$

# Quantum Electrodynamics – Intro to Field Theory

## Quantum Harmonic Oscillator

Formal Quantization Procedure:

$$q \rightarrow \hat{q}, \quad p \rightarrow \hat{p}, \quad [\hat{q}, \hat{p}] = i\hbar$$

Choose  $E_0 = \hbar\omega$    $\Rightarrow q_0 = \sqrt{\frac{2\hbar}{m\omega}}, \quad p_0 = \sqrt{2m\hbar\omega}$

$$\alpha \rightarrow \hat{a} = \hat{Q} + i\hat{P} = \sqrt{\frac{m\omega}{2\hbar}} \left( \hat{q} + i\frac{\hat{p}}{m\omega} \right)$$
$$[\hat{a}, \hat{a}^\dagger] = 1$$

Rewrite:

$$\hat{H} = \hbar\omega(\hat{Q}^2 + \hat{P}^2) = \hbar\omega(\hat{a}^\dagger \hat{a} + \frac{1}{2})$$
$$\hat{N} = \hat{a}^\dagger \hat{a} \quad (\text{number operator})$$

Commutator  $[\hat{H}, \hat{N}] = 0$

 joint energy/number states  $|n\rangle$

$$\hat{H}|n\rangle = \hbar\omega(n + \frac{1}{2})|n\rangle$$
$$\hat{N}|n\rangle = n|n\rangle$$

Commutators

$$\begin{aligned} [\hat{N}, \hat{a}^\dagger] &= \hat{a}^\dagger \\ [\hat{N}, \hat{a}] &= -\hat{a} \end{aligned} \quad \left. \right\} \quad \Rightarrow$$

$$\begin{aligned} \hat{a}|n\rangle &= \sqrt{n}|n-1\rangle \\ \hat{a}^\dagger|n\rangle &= \sqrt{n+1}|n+1\rangle \\ \hat{a}|0\rangle &= 0 \end{aligned}$$

Generating excited states

$$|n\rangle = \frac{1}{\sqrt{n!}} (\hat{a}^\dagger)^n |0\rangle$$

# Quantum Electrodynamics – Intro to Field Theory

Commutator  $[\hat{H}, \hat{N}] = 0$

→ joint energy/number states  $|n\rangle$

$$\hat{H}|n\rangle = \hbar\omega(n + 1/2)|n\rangle$$

$$\hat{N}|n\rangle = n|n\rangle$$

Commutators

$$\begin{aligned} [\hat{N}, \hat{a}^+] &= \hat{a}^+ \\ [\hat{N}, \hat{a}] &= -\hat{a} \end{aligned} \quad \left. \right\}$$

$$\begin{aligned} \hat{a}|n\rangle &= \sqrt{n}|n-1\rangle \\ \hat{a}^+|n\rangle &= \sqrt{n+1}|n+1\rangle \\ \hat{a}|0\rangle &= 0 \end{aligned}$$

Generating excited states

$$|n\rangle = \frac{1}{\sqrt{n!}} (\hat{a}^+)^n |0\rangle$$

Expectation values for  $\hat{q}$  and  $\hat{p}$  in number states

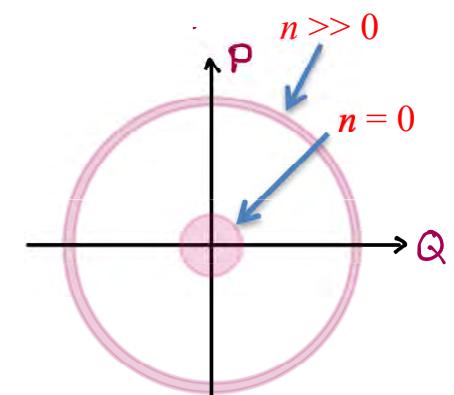
$$\langle n|\hat{q}|n\rangle = \langle n|\hat{p}|n\rangle = 0$$

$$\langle n|\hat{q}^2|n\rangle = \frac{q_0^2}{2}(n + 1/2) \neq 0$$

$$\langle n|\hat{p}^2|n\rangle = \frac{p_0^2}{2}(n + 1/2) \neq 0$$

$$\Delta q \Delta p = \frac{q_0 p_0}{2}(n + 1/2) = \hbar(n + 1/2)$$

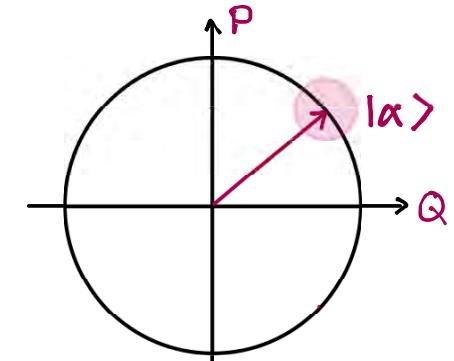
Phase space visualization  
of number states



Quasi-classical  
(coherent) state

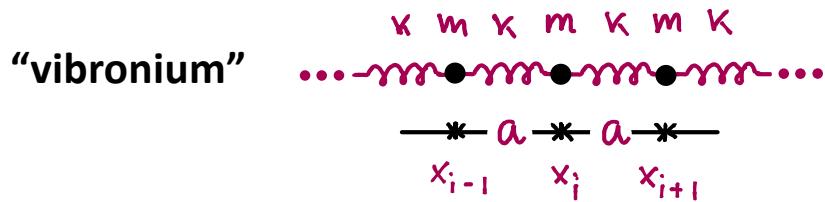
$$|\alpha\rangle = e^{-|\alpha|^2/2} \sum_i \frac{\alpha^n}{\sqrt{n!}} |n\rangle$$

$$\Delta q \Delta p = \hbar/2, \quad \Delta Q = \Delta P$$



# Quantum Electrodynamics – Intro to Field Theory

## Lagrange formulation of 1D Scalar Field



Configuration space =  $\{x_i\}$  (set of  $N$  osc. positions)

$$T = \sum_{i=1}^N \frac{1}{2} m \dot{x}_i^2, \quad V = \sum_{i=1}^N \frac{1}{2} K (x_{i+1} - x_i)^2$$

Lagrangian, equations of motion

Continuum limit  $\rightarrow$  Elastic rod

$$N \rightarrow \infty \quad m/a \rightarrow \mu \quad \leftarrow \text{linear mass density}$$

$$a \rightarrow dx \quad K a \rightarrow Y \quad \leftarrow \text{Youngs modulus}$$

$$\{x_i\} \rightarrow \eta(x) \quad \leftarrow \text{displacement field (sound)}$$

## Rewrite

$$T = \lim_{N \rightarrow \infty} \sum_{i=1}^N a \frac{1}{2} \left( \frac{m}{a} \right) \dot{x}_i^2 = \int dx \frac{1}{2} \mu \left( \frac{\partial \eta}{\partial t} \right)^2$$

$$V = \lim_{N \rightarrow \infty} \sum_{i=1}^N a \frac{1}{2} K a \left( \frac{x_{i+1} - x_i}{a} \right)^2 = \int dx \frac{1}{2} Y \left( \frac{\partial \eta}{\partial x} \right)^2$$

Lagrangian:

$$\mathcal{L} = T - V = \int dx \frac{1}{2} \mu \left( \frac{\partial \eta}{\partial t} \right)^2 - \int dx \frac{1}{2} Y \left( \frac{\partial \eta}{\partial x} \right)^2$$

Notes, Homework  $\rightarrow$  Scalar wave equation

$$\frac{\partial^2 \eta}{\partial t^2} - \frac{Y}{\mu} \frac{\partial^2 \eta}{\partial x^2} = 0$$

– Not yet ready for Quantization –

# Quantum Electrodynamics – Intro to Field Theory

Rewrite

$$T = \lim_{N \rightarrow \infty} \sum_{i=1}^N a \frac{1}{2} \left( \frac{m}{a} \dot{x}_i \right)^2 = \int dx \frac{1}{2} \mu \left( \frac{\partial \eta}{\partial t} \right)^2$$

$$V = \lim_{N \rightarrow \infty} \sum_{i=1}^N a \frac{1}{2} \times a \left( \frac{x_{i+1} - x_i}{a} \right)^2 = \int dx \frac{1}{2} \gamma \left( \frac{\partial \eta}{\partial x} \right)^2$$

Lagrangian:

$$\mathcal{L} = T - V = \int dx \frac{1}{2} \mu \left( \frac{\partial \eta}{\partial t} \right)^2 - \int dx \frac{1}{2} \gamma \left( \frac{\partial \eta}{\partial x} \right)^2$$

Notes, Homework → Scalar wave equation

$$\frac{\partial^2 \eta}{\partial t^2} - \frac{\gamma}{\mu} \frac{\partial^2 \eta}{\partial x^2} = 0$$

– Not yet ready for Quantization –

Normal Mode Decomposition

Field in cavity:



Solutions to wave eq.

Let  $\eta(x, t) = g(t)u(x) = g_0 e^{i\omega t} u(x)$

$$\ddot{u} - \nu^2 u'' = -\omega^2 g(t) u(x) - \nu^2 g(t) u''(x) = 0$$



$$u''(x) = -k^2 u(x), \quad k = \omega/\nu$$

Solutions in cavity:

$$u_{kx}(x) = \sqrt{\frac{2}{L}} \sin(kx), \quad k = \frac{n\pi}{L}$$

These standing waves are a set of Normal Modes

# Quantum Electrodynamics – Intro to Field Theory

## Normal Mode Decomposition

Field in cavity:



Solutions to wave eq.

Let  $y(x,t) = g(t)u(x) = g_0 e^{i\omega t} u(x)$

$$i\ddot{y} - \nu^2 y'' = -\omega^2 g(t) u(x) - \nu^2 g(t) u''(x) = 0$$



$$u''(x) = -k^2 u(x), \quad k = \omega/\nu$$

Solutions in cavity:

$$u_{nk}(x) = \sqrt{\frac{2}{L}} \sin(kx), \quad k = \frac{n\pi}{L}$$

These standing waves are a set of Normal Modes

These modes are orthonormal and complete



$$y(x,t) = \sqrt{L} \sum_{nk} q_{nk}(t) u_{nk}(x)$$

Normal mode expansion of  $y(x,t)$  in basis  $u_{nk}(x)$

Lagrangian for the acoustic field:

$$\begin{aligned} T &= \int dx \frac{1}{2} \mu \left( \frac{\partial y}{\partial t} \right)^2 = \sum_{nk,nk'} \underbrace{\frac{1}{2} \mu L q_{nk} \dot{q}_{nk}}_M \underbrace{\int dx u_{nk}(x) u_{nk'}(x)}_{\delta_{nk,nk'}} \\ &= \sum_{nk} \frac{1}{2} M \dot{q}_{nk}^2 \end{aligned}$$

$$\begin{aligned} V &= \int dx \frac{1}{2} \gamma \left( \frac{\partial y}{\partial x} \right)^2 = \sum_{nk,nk'} \underbrace{\frac{1}{2} \gamma L q_{nk} q_{nk'}}_{M} \underbrace{\int dx \left( \frac{\partial u_{nk}}{\partial x} \right) \left( \frac{\partial u_{nk'}}{\partial x} \right)}_{\delta_{nk,nk'}} \\ &= \sum_{nk} \frac{1}{2} M \omega_{nk}^2 q_{nk}^2 \end{aligned}$$



$$\mathcal{L} = T - V = \sum_{nk} \left( \frac{1}{2} M \dot{q}_{nk}^2 - \frac{1}{2} M \omega_{nk}^2 q_{nk}^2 \right) = \sum_{nk} \mathcal{L}_{nk}$$

# Quantum Electrodynamics – Intro to Field Theory

The rest now follows from the Lagrangian

$$\mathcal{L} = T - V = \sum_k \left( \frac{1}{2} M \dot{q}_k^2 - \frac{1}{2} M \omega_k^2 q_k^2 \right) = \sum_k \mathcal{L}_k$$



Canonical Momentum

$$p_k = \frac{\partial \mathcal{L}}{\partial \dot{q}_k} = M \dot{q}_k$$

Hamiltonian

$$\mathcal{H}(\{p_k, q_k\}) = T + V = \sum_k \left( \frac{p_k^2}{2M} + \frac{1}{2} M \omega_k^2 q_k^2 \right)$$

( collection of SHO's, one for each normal mode )

Following the standard recipe...

$$E_{0,k} = \hbar \omega_k, \quad q_{0,k} = \sqrt{2\hbar/M\omega_k}, \quad p_{0,k} = \sqrt{2M\hbar\omega_k}$$

$$Q_k = q_k/q_{0,k}, \quad P_k = p_k/p_{0,k}, \quad \alpha_k = Q_k + iP_k$$

... we get solutions

$$\alpha_{k*}(t) = Q_k(t) + iP_k(t) = \alpha_k(0) e^{-i\omega_k t}$$

This finally gives us

$$\mathcal{H} = \sum_k \hbar \omega_k (Q_k^2 + P_k^2) = \sum_k \hbar \omega_k \alpha_k^* \alpha_k$$

$$\begin{aligned} g(x, L) &= \sqrt{L} \sum_k q_{k*}(t) u_k(x) \\ &= \frac{1}{2} \sum_k \sqrt{L q_{0,k}^2} (\alpha_k(t) u_k(x) + \alpha_k^*(t) u_k^*(x)) \end{aligned}$$

allows real or complex  $u_k(x)$

# Quantum Electrodynamics – Intro to Field Theory

... we get solutions

$$\alpha_{k\epsilon}(t) = Q_{k\epsilon}(t) + i P_{k\epsilon}(t) = \alpha_{k\epsilon}(0) e^{-i\omega_k t}$$

This finally gives us

$$\mathcal{H} = \sum_{k\epsilon} \hbar \omega_{k\epsilon} (Q_{k\epsilon}^2 + P_{k\epsilon}^2) = \sum_{k\epsilon} \hbar \omega_{k\epsilon} \alpha_{k\epsilon}^* \alpha_{k\epsilon}$$

$$\begin{aligned} g(x, t) &= \sqrt{L} \sum_{k\epsilon} q_{k\epsilon}(t) u_{k\epsilon}(x) \\ &= \frac{1}{2} \sum_{k\epsilon} \sqrt{L q_{0,k}^2} (\alpha_{k\epsilon}(t) u_{k\epsilon}(x) + \alpha_{k\epsilon}^*(t) u_{k\epsilon}^*(x)) \end{aligned}$$

allows real or complex  $u_k(x)$

Formal Quantization Procedure:

$$q_{k\epsilon} \rightarrow \hat{q}_{k\epsilon}, \quad p_{k\epsilon} \rightarrow \hat{p}_{k\epsilon}, \quad \alpha_{k\epsilon} \rightarrow \hat{\alpha}_{k\epsilon}$$

$$[\hat{q}_k, \hat{p}_{k'\epsilon'}] = i\hbar \delta_{kk'}, \quad [\hat{a}_k, \hat{a}_{k'}^+] = \delta_{kk'}, \quad [\hat{a}_k, \hat{a}_{k'}] = 0$$

Note:  $k \neq k'$  → operators commute  
(normal modes = independent degs. of freedom)

Hamiltonian & Quantized fields

$$\hat{H} = \sum_{k\epsilon} \hbar \omega_{k\epsilon} (\hat{a}_{k\epsilon}^+ \hat{a}_{k\epsilon} + \frac{1}{2})$$

$$\hat{g}(x) = \sqrt{L} \sum_{k\epsilon} \hat{q}_{k\epsilon} u_{k\epsilon}(x) = \sum_{k\epsilon} \sqrt{L q_{0,k}^2} (\hat{a}_{k\epsilon} u_{k\epsilon}(x) + \hat{a}_{k\epsilon}^+ u_{k\epsilon}^*(x))$$

$$\hat{\pi}(x) = \frac{1}{\sqrt{L}} \sum_{k\epsilon} \hat{p}_{k\epsilon} u_{k\epsilon}(x) = -i \sum_{k\epsilon} \sqrt{\frac{n_{0,k}^2}{L}} (\hat{a}_{k\epsilon} u_{k\epsilon}(x) - \hat{a}_{k\epsilon}^+ u_{k\epsilon}^*(x))$$

field  $\hat{g}(x)$  and canonical momentum field  $\hat{\pi}(x)$

$$\rightarrow [\hat{g}(x), \hat{\pi}(x')] = i\hbar \delta(x-x')$$

# Quantum Electrodynamics – Intro to Field Theory

Formal Quantization Procedure:

$$q_k \rightarrow \hat{q}_k, \quad p_k \rightarrow \hat{p}_k, \quad a_k \rightarrow \hat{a}_k$$

$$[\hat{q}_k, \hat{p}_{k'}] = i\hbar \delta_{kk'}, \quad [\hat{a}_k, \hat{a}_{k'}^\dagger] = \delta_{kk'}, \quad [\hat{a}_k, \hat{a}_{k'}] = 0$$

Note:  $k \neq k'$  operators commute  
( normal modes = independent degs. of freedom )

Hamiltonian & Quantized fields

$$\hat{H} = \sum_k \hbar \omega_k (\hat{a}_k^\dagger \hat{a}_k + \frac{1}{2})$$

$$\hat{\eta}(x) = \sqrt{L} \sum_k \hat{q}_k u_k(x) = \sum_k \sqrt{\frac{\hbar \omega_k}{2m}} (\hat{a}_k u_k(x) + \hat{a}_k^\dagger u_k^*(x))$$

$$\hat{\pi}(x) = \frac{1}{\sqrt{L}} \sum_k \hat{p}_k u_k(x) = -i \sum_k \sqrt{\frac{\hbar \omega_k}{2m}} (\hat{a}_k u_k(x) - \hat{a}_k^\dagger u_k^*(x))$$

field  $\hat{\eta}(x)$  and canonical momentum field  $\hat{\pi}(x)$

$$\Rightarrow [\hat{\eta}(x), \hat{\pi}(x')] = i\hbar \delta(x-x')$$

Quantum States:

The quantum field is a collection of QSHO's



State space = tensor product of SHO spaces

Fock Space  $\mathcal{E} = \mathcal{E}_{k_1} \otimes \mathcal{E}_{k_2} \otimes \mathcal{E}_{k_3} \otimes \dots, \mathcal{E}_{k_j}$

Fock State  $| \{n_{k_1}, n_{k_2}, \dots \} \rangle = |n_{k_1}\rangle \otimes |n_{k_2}\rangle \otimes \dots |n_{k_j}\rangle$

$\hat{a}_{k_i}, \hat{a}_{k_i}^\dagger$  destroy/create excitations in mode  $k_i$

Vacuum State  $|0\rangle$  zero quanta in every mode

Favorite Question: What is a Phonon?

End 03-30-2022