

# Quantized Light – Matter Interactions

With this notation

(4)

$$\hat{H}_{AF} = \hbar \sum_{\vec{k}} (g_{\vec{k}} \hat{\sigma}_+ \hat{a}_{\vec{k}} + \cancel{g_{\vec{k}} \hat{\sigma}_+ \hat{a}_{\vec{k}}^+} + \cancel{g_{\vec{k}}^* \hat{\sigma}_- \hat{a}_{\vec{k}}} + g_{\vec{k}}^* \hat{\sigma}_- \hat{a}_{\vec{k}}^+)$$

Energy conservation?



$$\hat{H}_{AF} = \hbar \sum_{\vec{k}} (g_{\vec{k}} \hat{\sigma}_+ \hat{a}_{\vec{k}} + g_{\vec{k}}^* \hat{\sigma}_- \hat{a}_{\vec{k}}^+)$$

Putting it all together

(5)

$$\hat{H} = \hat{H}_F + \hat{H}_A + \hat{H}_{AF} =$$

$$\sum_{\vec{k}} \hbar \omega_{\vec{k}} \hat{a}_{\vec{k}}^+ \hat{a}_{\vec{k}} + \frac{1}{2} \hbar \omega_{21} \hat{\sigma}_z + \sum_{\vec{k}} \hbar (g_{\vec{k}} \hat{\sigma}_+ \hat{a}_{\vec{k}} + g_{\vec{k}}^* \hat{\sigma}_- \hat{a}_{\vec{k}}^+)$$



Foundational result for remainder of course

We changed the zero point for energy by subtracting

$$\sum_{\vec{k}} \frac{1}{2} \hbar \omega_{\vec{k}} \quad \text{and} \quad \frac{1}{2} (E_2 - E_1)$$

field                      atom

# Quantized Light – Matter Interactions

With this notation

$$\hat{H}_{AF} = \hbar \sum_{\vec{k}} (g_{\vec{k}} \hat{\sigma}_+ \hat{a}_{\vec{k}} + \cancel{g_{\vec{k}} \hat{\sigma}_+ \hat{a}_{\vec{k}}^{\dagger}} + \cancel{g_{\vec{k}}^* \hat{\sigma}_- \hat{a}_{\vec{k}}} + g_{\vec{k}}^* \hat{\sigma}_- \hat{a}_{\vec{k}}^{\dagger}) \quad (4)$$

Energy conservation?



$$\hat{H}_{AF} = \hbar \sum_{\vec{k}} (g_{\vec{k}} \hat{\sigma}_+ \hat{a}_{\vec{k}} + g_{\vec{k}}^* \hat{\sigma}_- \hat{a}_{\vec{k}}^{\dagger})$$

Putting it all together

$$\hat{H} = \hat{H}_F + \hat{H}_A + \hat{H}_{AF} = \quad (5)$$

$$\sum_{\vec{k}} \hbar \omega_{\vec{k}} \hat{a}_{\vec{k}}^{\dagger} \hat{a}_{\vec{k}} + \frac{1}{2} \hbar \omega_{21} \hat{\sigma}_z + \sum_{\vec{k}} \hbar (g_{\vec{k}} \hat{\sigma}_+ \hat{a}_{\vec{k}} + g_{\vec{k}}^* \hat{\sigma}_- \hat{a}_{\vec{k}}^{\dagger})$$

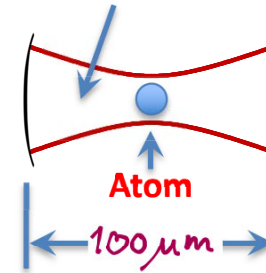
We changed the zero point for energy by subtracting

$$\sum_{\vec{k}} \frac{1}{2} \hbar \omega_{\vec{k}} \quad \text{field} \quad \text{and} \quad \frac{1}{2} (\epsilon_2 - \epsilon_1) \quad \text{atom}$$

## Interaction with Single-mode Fields

Good approx. in small, high-Q Cavity

Gaussian beam mode



$$c/2L \gg A_{21}$$

$$|g_{\vec{k}}| \gg A_{21}, \delta$$

Single-mode (Jaynes-Cummings) Hamiltonian

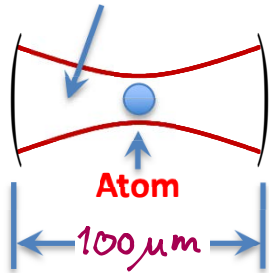
$$\hat{H} = \underbrace{\hbar \omega \hat{a}^{\dagger} \hat{a}}_{H_0} + \underbrace{\frac{1}{2} \hbar \omega_{21} \hat{\sigma}_z + \hbar (g_{\vec{k}} \hat{\sigma}_+ + g_{\vec{k}}^* \hat{\sigma}_-) (\hat{a}_{\vec{k}} + \hat{a}_{\vec{k}}^{\dagger})}_{H_{AF}}$$

# Quantized Light – Matter Interactions

## Interaction with Single-mode Fields

Good approx. in small, high-Q Cavity

Gaussian beam mode



$$c/2L \gg A_{21}$$

$$|g_{\vec{k}}| \gg A_{21}, \delta$$

Single-mode (Jaynes-Cummings) Hamiltonian

$$\hat{H} = \underbrace{\hbar\omega \hat{a}^\dagger \hat{a} + \frac{1}{2} \hbar\omega_{21} \hat{\sigma}_z}_{H_0} + \underbrace{\hbar(g_{\vec{k}} \hat{\sigma}_+ + g_{\vec{k}}^* \hat{\sigma}_-)(\hat{a}_{\vec{k}} + \hat{a}_{\vec{k}}^\dagger)}_{H_{AF}}$$

Paradigm for spin-1/2 coupled to QHO

- Atom in high-Q cavity \*)
- Quantum dot in high-Q
- Rydberg atom in superconducting  $\mu\text{W}$  Cavity
- Superconducting qubit in superconducting  $\mu\text{W}$  cavity
- Superconducting qubit in superconducting  $\mu\text{W}$  stripline cavity (circuit QED)
- Trapped ion with quantized COM motion \*)

\*) Nobel Prize in Physics 2012

# Quantized Light – Matter Interactions

## Paradigm for spin-1/2 coupled to QHO

- Atom in high-Q cavity \*)
- Quantum dot in high-Q
- Rydberg atom in superconducting  $\mu\text{w}$  Cavity
- Superconducting qubit in superconducting  $\mu\text{w}$  cavity
- Superconducting qubit in superconducting  $\mu\text{w}$  stripline cavity (circuit QED)
- Trapped ion with quantized COM motion \*)

\*) Nobel Prize in Physics 2012

## More about Single-mode Cavity QED

$$\hat{H} = \underbrace{\hbar\omega\hat{a}^\dagger\hat{a} + \frac{1}{2}\hbar\omega_2\hat{\sigma}_z}_{H_0} + \underbrace{\hbar(g_{\vec{R}}\hat{\sigma}_+ + g_{\vec{R}}^*\hat{\sigma}_-)(\hat{a}_{\vec{R}} + \hat{a}_{\vec{R}}^\dagger)}_{H_{AF}}$$

For simplicity  $\vec{n}_{21} = \vec{n}_{12} \Rightarrow g_{\vec{R}} = g_{\vec{R}}^* = g_{\vec{R}}$

Note:  $\hat{H}_{AF}$  conserves excitation number, couples  $|2, n\rangle \leftrightarrow |1, n+1\rangle$



Series of 2-level systems, one for each  $n$

All 2-level systems are alike  
Rabi problem!

Switch to Interaction Picture, 6-7 in Notes:

$$\left. \begin{aligned} \hat{H}_S &\rightarrow \hat{H}_I = e^{i\frac{\hat{H}_0 t}{\hbar}} \hat{H}_{AF} e^{-i\frac{\hat{H}_0 t}{\hbar}} \\ |\psi_S(t)\rangle &\rightarrow |\psi_I(t)\rangle = e^{i\frac{\hat{H}_0 t}{\hbar}} |\psi_S(t)\rangle \end{aligned} \right\} \rightarrow$$

# Quantized Light – Matter Interactions

## More about Single-mode Cavity QED

$$\hat{H} = \underbrace{\hbar\omega\hat{a}^\dagger\hat{a} + \frac{1}{2}\hbar\omega_2\hat{\sigma}_2}_{H_0} + \underbrace{\hbar(g_{\vec{R}}\hat{\sigma}_+ + g_{\vec{R}}^*\hat{\sigma}_-)(\hat{a}_{\vec{R}} + \hat{a}_{\vec{R}}^\dagger)}_{H_{AF}}$$

For simplicity  $\vec{n}_{21} = \vec{n}_{12} \Rightarrow g_{\vec{R}} = g_{\vec{R}}^* = g_{\vec{R}}$

Note:  $\hat{H}_{AF}$  conserves excitation number, couples  $|2, n\rangle \leftrightarrow |1, n+1\rangle$

Series of 2-level systems, one for each  $n$

All 2-level systems are alike  
Rabi problem!

Switch to Interaction Picture, 6-7 in Notes:

$$\left. \begin{aligned} \hat{H}_S \rightarrow \hat{H}_I &= e^{i\frac{\hat{H}_0 t}{\hbar}} \hat{H}_{AF} e^{-i\frac{\hat{H}_0 t}{\hbar}} \\ |\psi_S(t)\rangle \rightarrow |\psi_I(t)\rangle &= e^{i\frac{\hat{H}_0 t}{\hbar}} |\psi_S(t)\rangle \end{aligned} \right\} \Rightarrow$$

$$\hat{H}_{AF} = \hbar g (\hat{\sigma}_+ \hat{a} e^{i(\omega_2 - \omega)t} + \hat{\sigma}_+ \hat{a}^\dagger e^{i(\omega_2 + \omega)t} + \hat{\sigma}_- \hat{a} e^{-i(\omega_2 + \omega)t} + \hat{\sigma}_- \hat{a}^\dagger e^{-i(\omega_2 - \omega)t})$$

RWA and resonant approximation



Jaynes-Cummings Hamiltonian

$$\hat{H}_I = \hat{H}_{AF} = \hbar g (\hat{\sigma}_+ \hat{a} e^{i\Delta t} + \hat{\sigma}_- \hat{a}^\dagger e^{-i\Delta t})$$

$$\Delta = \omega_2 - \omega$$

Can show

$$\begin{aligned} e^{i\omega\hat{a}^\dagger\hat{a}t} \hat{a} e^{-i\omega\hat{a}^\dagger\hat{a}t} &= \hat{a} e^{-i\omega t} \\ e^{i\frac{\omega_2}{2}\hat{\sigma}_2 t} \hat{\sigma}_+ e^{-i\frac{\omega_2}{2}\hat{\sigma}_2 t} &= \hat{\sigma}_+ e^{i\omega_2 t} \end{aligned}$$

# Quantized Light – Matter Interactions

## More about Single-mode Cavity QED

$$\hat{H} = \underbrace{\hbar\omega\hat{a}^\dagger\hat{a} + \frac{1}{2}\hbar\omega_{21}\hat{\sigma}_z}_{H_0} + \underbrace{\hbar(g_{\vec{R}}\hat{\sigma}_+ + g_{\vec{R}}^*\hat{\sigma}_-)(\hat{a}_{\vec{R}} + \hat{a}_{\vec{R}}^\dagger)}_{H_{AF}}$$

For simplicity  $\vec{n}_{21} = \vec{n}_{12} \Rightarrow g_{\vec{R}} = g_{\vec{R}}^* = g_{\vec{R}}$

Note:  $\hat{H}_{AF}$  conserves excitation number, couples  $|2, n\rangle \leftrightarrow |1, n+1\rangle$

Series of 2-level systems, one for each  $n$

**All 2-level systems are alike**  
**Rabi problem!**

Switch to Interaction Picture, 6-7 in Notes:

$$\left. \begin{aligned} \hat{H}_S \rightarrow \hat{H}_I &= e^{i\frac{\hat{H}_0 t}{\hbar}} \hat{H}_{AF} e^{-i\frac{\hat{H}_0 t}{\hbar}} \\ |\psi_S(t)\rangle \rightarrow |\psi_I(t)\rangle &= e^{i\frac{\hat{H}_0 t}{\hbar}} |\psi_S(t)\rangle \end{aligned} \right\} \Rightarrow$$

$$\hat{H}_{AF} = \hbar g (\hat{\sigma}_+ \hat{a} e^{i(\omega_{21} - \omega)t} + \hat{\sigma}_+ \hat{a}^\dagger e^{i(\omega_{21} + \omega)t} + \hat{\sigma}_- \hat{a} e^{-i(\omega_{21} + \omega)t} + \hat{\sigma}_- \hat{a}^\dagger e^{-i(\omega_{21} - \omega)t})$$

RWA and resonant approximation

Jaynes-Cummings Hamiltonian

$$\hat{H}_I = \hat{H}_{AF} = \hbar g (\hat{\sigma}_+ \hat{a} e^{i\Delta t} + \hat{\sigma}_- \hat{a}^\dagger e^{-i\Delta t})$$

$$\Delta = \omega_{21} - \omega$$

Begin 04-26-2021

Eigenstates of  $\hat{H}_0 = \hat{H}_F + \hat{H}_A$

State	Energy
$ 2, n\rangle$	$\hbar\omega n + \hbar\omega_{21}$
$ 1, n+1\rangle$	$\hbar\omega(n+1) - \hbar\omega_{21}$

# Quantized Light – Matter Interactions

$$\hat{H}_{AF} = \hbar g (\hat{\sigma}_+ \hat{a} e^{i(\omega_{21} - \omega)t} + \hat{\sigma}_+ \hat{a}^\dagger e^{i(\omega_{21} + \omega)t} + \hat{\sigma}_- \hat{a} e^{-i(\omega_{21} + \omega)t} + \hat{\sigma}_- \hat{a}^\dagger e^{-i(\omega_{21} - \omega)t})$$

RWA and resonant approximation



Jaynes-Cummings Hamiltonian

$$\hat{H}_I = \hat{H}_{AF} = \hbar g (\hat{\sigma}_+ \hat{a} e^{i\Delta t} + \hat{\sigma}_- \hat{a}^\dagger e^{-i\Delta t})$$

$$\Delta = \omega_{21} - \omega$$

Eigenstates of  $\hat{H}_0 = \hat{H}_F + \hat{H}_A$

State

Energy

$|2, n\rangle$

$\hbar\omega n + \hbar\omega_{21}$

$|1, n+1\rangle$

$\hbar\omega(n+1) - \hbar\omega_{21}$

Cavity QED version of the Rabi Problem

$$|\psi(0)\rangle = |2, n\rangle$$

$$|\psi(t)\rangle = C_{1, n+1} |1, n+1\rangle + C_{2, n} |2, n\rangle$$

Matrix elements

$$\langle 2, n | \hat{H}_{AF} | 1, n+1 \rangle = \hbar g \sqrt{n+1} e^{i\Delta t}$$

$$\langle 1, n+1 | \hat{H}_{AF} | 2, n \rangle = \hbar g \sqrt{n+1} e^{-i\Delta t}$$



Schrödinger Equation

$$i\hbar \frac{d}{dt} \begin{pmatrix} C_{1, n+1} \\ C_{2, n} \end{pmatrix} =$$

$$\hbar g \sqrt{n+1} \begin{pmatrix} 0 & e^{-i\Delta t} \\ e^{i\Delta t} & 0 \end{pmatrix} \begin{pmatrix} C_{1, n+1} \\ C_{2, n} \end{pmatrix}$$

# Quantized Light – Matter Interactions

## Cavity QED version of the Rabi Problem

$$|\psi(0)\rangle = |2, n\rangle$$

$$|\psi(t)\rangle = c_{1, n+1} |1, n+1\rangle + c_{2, n} |2, n\rangle$$

## Matrix elements

$$\langle 2, n | \hat{H}_{AF} | 1, n+1 \rangle = \hbar g \sqrt{n+1} e^{i\Delta t}$$

$$\langle 1, n+1 | \hat{H}_{AF} | 2, n \rangle = \hbar g \sqrt{n+1} e^{-i\Delta t}$$



## Schrödinger Equation

$$i\hbar \frac{d}{dt} \begin{pmatrix} c_{1, n+1} \\ c_{2, n} \end{pmatrix} =$$

$$\hbar g \sqrt{n+1} \begin{pmatrix} 0 & e^{-i\Delta t} \\ e^{i\Delta t} & 0 \end{pmatrix} \begin{pmatrix} c_{1, n+1} \\ c_{2, n} \end{pmatrix}$$



$$\dot{c}_{1, n+1} = -ig\sqrt{n+1} e^{-i\Delta t} c_{2, n}$$
$$\dot{c}_{2, n} = -ig\sqrt{n+1} e^{i\Delta t} c_{1, n+1}$$

Substitute  $c_{1, n+1} \rightarrow c_1$ ,  $c_{2, n} \rightarrow c_2 e^{i\Delta t}$

Looks **exactly** like Semiclassical Rabi problem

Solve for  $c_1(0) = 0$ ,  $c_2(0) = 1$



$$c_{2, n}(t) = \left[ \cos\left(\frac{\Omega_n t}{2}\right) - i \frac{\Delta}{\Omega_n} \sin\left(\frac{\Omega_n t}{2}\right) \right] e^{i\Delta t/2}$$

$$c_{1, n+1} = -i \frac{2g\sqrt{n+1}}{\Omega_n} \sin\left(\frac{\Omega_n t}{2}\right) e^{-i\Delta t/2}$$

$$\Omega_n = (4g^2(n+1) + \Delta^2)^{1/2}$$



# Quantized Light – Matter Interactions



$$\dot{c}_{1,n+1} = -ig\sqrt{n+1} e^{-i\Delta t} c_{2,n}$$

$$\dot{c}_{2,n} = -ig\sqrt{n+1} e^{i\Delta t} c_{1,n+1}$$

Substitute  $c_{1,n+1} \rightarrow c_1$ ,  $c_{2,n} \rightarrow c_2 e^{i\Delta t}$

Looks **exactly** like Semiclassical Rabi problem

Solve for  $c_1(0) = 0$ ,  $c_2(0) = 1$



$$c_{2,n}(t) = \left[ \cos\left(\frac{\Omega_n t}{2}\right) - i \frac{\Delta}{\Omega_n} \sin\left(\frac{\Omega_n t}{2}\right) \right] e^{i\Delta t/2}$$

$$c_{1,n+1} = -i \frac{2g\sqrt{n+1}}{\Omega_n} \sin\left(\frac{\Omega_n t}{2}\right) e^{-i\Delta t/2}$$

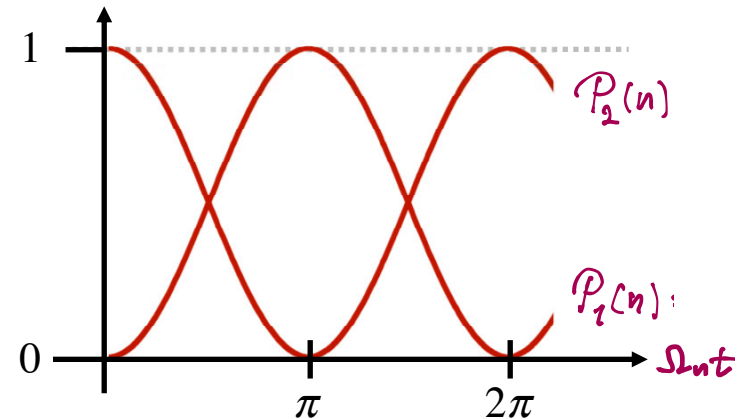
$$\Omega_n = (4g^2(n+1) + \Delta^2)^{1/2}$$

## Rabi Oscillations

$$P_2(n) = \cos^2\left(\frac{\Omega_n t}{2}\right) + \left(\frac{\Delta}{\Omega_n}\right)^2 \sin^2\left(\frac{\Omega_n t}{2}\right)$$

$$P_1(n) = \frac{4g^2(n+1)}{\Omega_n^2} \sin^2\left(\frac{\Omega_n t}{2}\right)$$

Example:  $\Delta = 0$



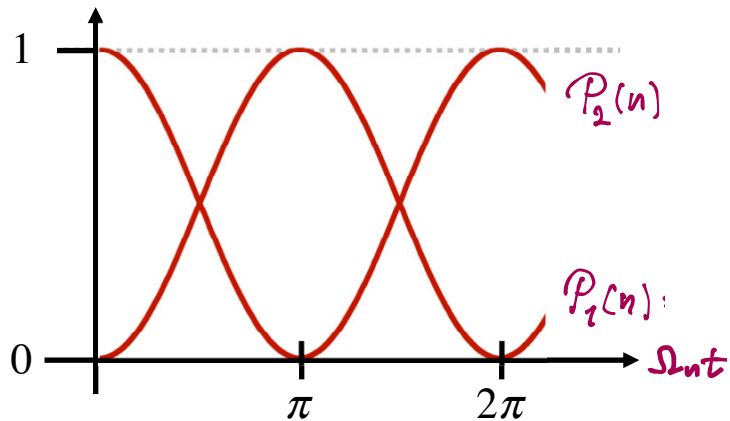
# Quantized Light – Matter Interactions

## Rabi Oscillations

$$P_2(n) = \cos^2\left(\frac{\Omega_n t}{2}\right) + \left(\frac{\Delta}{\Omega_n}\right)^2 \sin^2\left(\frac{\Omega_n t}{2}\right)$$

$$P_1(n) = \frac{4g^2(n+1)}{\Omega_n^2} \sin^2\left(\frac{\Omega_n t}{2}\right)$$

Example:  $\Delta = 0$



## Vacuum Rabi Oscillations

If  $|2(0)\rangle = |2,0\rangle \rightarrow$  no photons in field

yet  $|2,0\rangle$  evolves into  $|1,1\rangle$

**Uniquely QED phenomenon!**

$$\text{Asymmetry} \begin{cases} |2, n=0\rangle \rightarrow |1, n=1\rangle \\ |1, n=0\rangle \rightarrow |1, n=0\rangle \end{cases}$$

holds germ of **Spontaneous Decay**

# Quantized Light – Matter Interactions

## Vacuum Rabi Oscillations

If  $|\psi(0)\rangle = |2,0\rangle \Rightarrow$  no photons in field

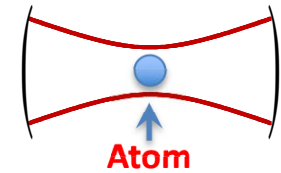
yet  $|2,0\rangle$  evolves into  $|1,1\rangle$

**Uniquely QED phenomenon!**

Asymmetry  $\left\{ \begin{array}{l} |2, n=0\rangle \rightarrow |1, n=1\rangle \\ |1, n=0\rangle \rightarrow |1, n=1\rangle \end{array} \right.$

holds germ of **Spontaneous Decay**

## Next: More Cavity QED



2-level atom

Single cavity mode

What happens with a Coherent State in the Cavity mode?

**(Quantum-Classical correspondence)**

Initial atom-field state:

$$|\psi(0)\rangle = |1\rangle \otimes |\alpha\rangle = \sum_n C_n |1, n\rangle, \quad C_n = e^{-\frac{1}{2}|\alpha|^2} \frac{\alpha^n}{\sqrt{n!}}$$

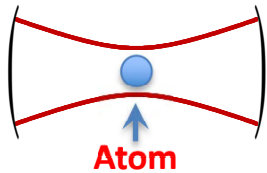
↑
↑  
atom
field

From Rabi solutions:  $\Delta = 0 \Rightarrow$

$$C_{1,n} = \cos\left(\frac{\Omega n t}{2}\right), \quad C_{2,n-1} = -i \sin\left(\frac{\Omega n t}{2}\right)$$

# Quantized Light – Matter Interactions

Next: More Cavity QED



2-level atom

Single cavity mode

What happens with a Coherent State  
in the Cavity mode?

(Quantum-Classical correspondence)

Initial atom-field state:

$$|\psi(0)\rangle = |1\rangle \otimes |\alpha\rangle = \sum_n C_n |1, n\rangle, \quad C_n = e^{-\frac{1}{2}|\alpha|^2} \frac{\alpha^n}{\sqrt{n!}}$$

↑
↑  
atom
field

From Rabi solutions:  $\Delta = 0 \Rightarrow$

$$C_{1,n} = \cos\left(\frac{\Omega_n t}{2}\right), \quad C_{2,n-1} = -i \sin\left(\frac{\Omega_n t}{2}\right)$$

Therefore

uncoupled

$$|\psi(t)\rangle = C_0 |1, 0\rangle + \sum_{n=1}^{\infty} C_n \left[ \cos\left(\frac{\Omega_n t}{2}\right) |1, n\rangle - i \sin\left(\frac{\Omega_n t}{2}\right) |2, n-1\rangle \right]$$

Consider the Atomic Excited State Population

$$\begin{aligned} P_2(t) &= \sum_{n=0}^{\infty} P_{2,n} = \sum_{n=0}^{\infty} |\langle 2, n | \psi(t) \rangle|^2 \\ &= \sum_{n=0}^{\infty} |C_n|^2 \sin^2\left(\frac{\Omega_n t}{2}\right) \\ &= \sum_{n=0}^{\infty} \frac{|\alpha|^{2n}}{n!} e^{-|\alpha|^2} \sin^2\left(\frac{\Omega_n t}{2}\right) \end{aligned}$$

Use  $|\alpha|^2 = \bar{n}$  and  $\Omega_n = 2g\sqrt{n}$

$$P_2(t) = \sum_{n=0}^{\infty} \frac{(\bar{n})^n e^{-\bar{n}}}{n!} \sin^2(g\sqrt{n}t)$$

# Quantized Light – Matter Interactions

Therefore

uncoupled

$$| \psi(t) \rangle = c_0 | 1, 0 \rangle + \sum_{n=1}^{\infty} c_n \left[ \cos\left(\frac{\Omega_n t}{2}\right) | 1, n \rangle - i \sin\left(\frac{\Omega_n t}{2}\right) | 2, n-1 \rangle \right]$$

Consider the Atomic Excited State Population

$$\begin{aligned} P_2(t) &= \sum_{n=0}^{\infty} P_{2,n} = \sum_{n=0}^{\infty} | \langle 2, n | \psi(t) \rangle |^2 \\ &= \sum_{n=0}^{\infty} |c_n|^2 \sin^2\left(\frac{\Omega_n t}{2}\right) \\ &= \sum_{n=0}^{\infty} \frac{|\alpha|^{2n}}{n!} e^{-|\alpha|^2} \sin^2\left(\frac{\Omega_n t}{2}\right) \end{aligned}$$

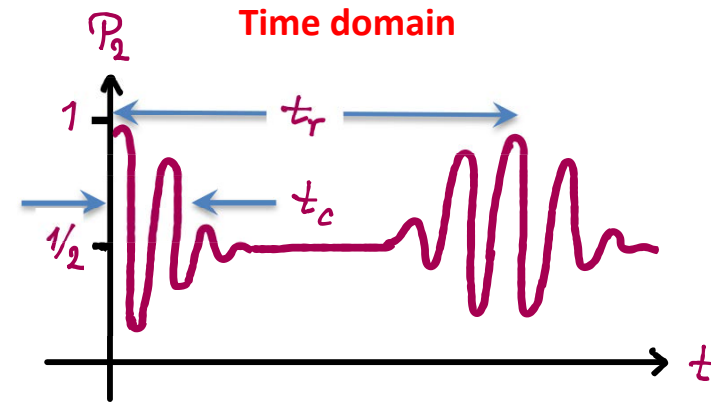
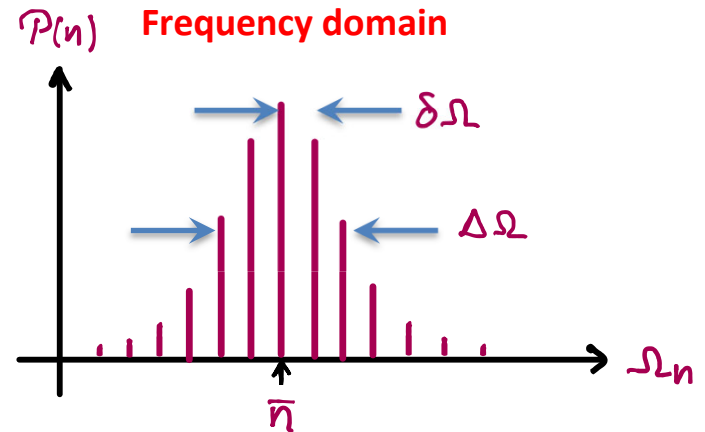
Use  $|\alpha|^2 = \bar{n}$  and  $\Omega_n = 2g\sqrt{n}$

$$P_2(t) = \sum_{n=0}^{\infty} \frac{(\bar{n})^n e^{-\bar{n}}}{n!} \sin^2(g\sqrt{n}t)$$

- Poisson weighted average of sinusoids
- Sinusoids gradually dephase over time



Collapse of oscillation amplitude

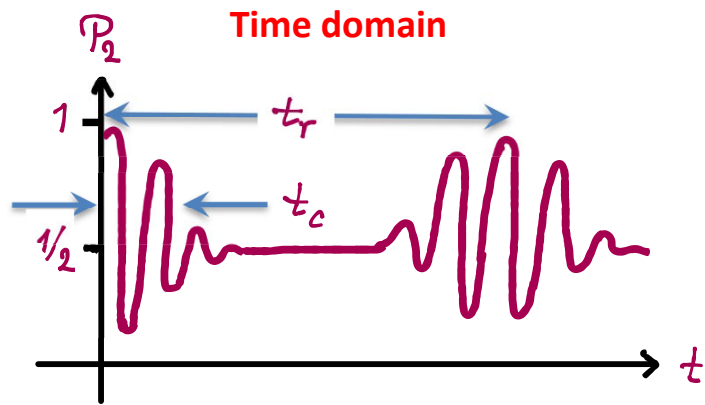
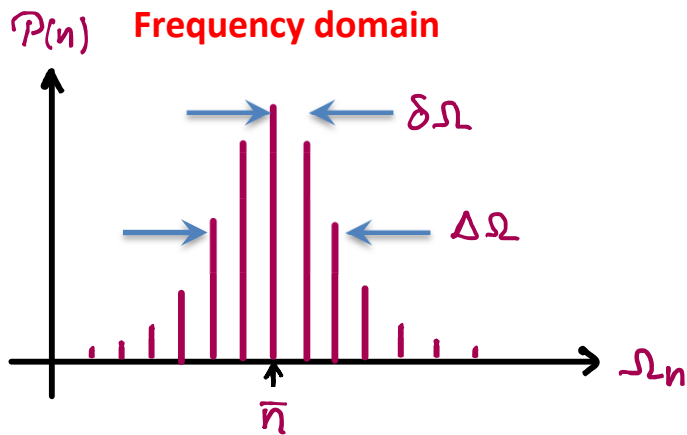


# Quantized Light – Matter Interactions

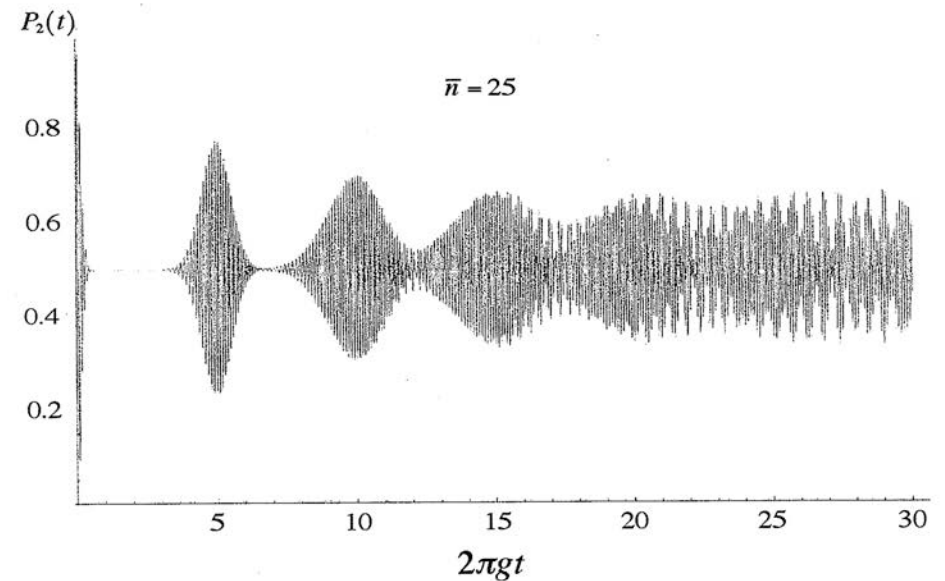
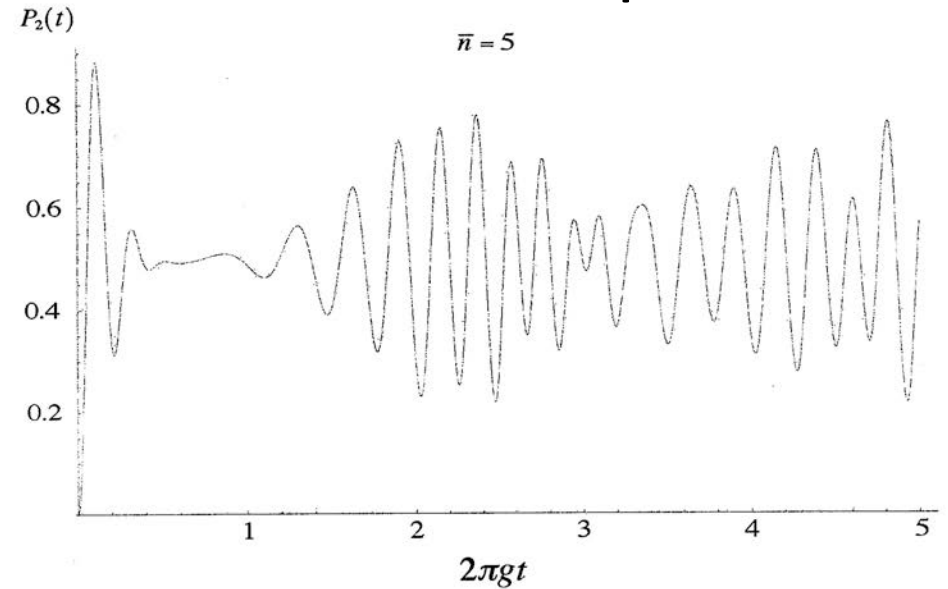
- Poisson weighted average of sinusoids
- Sinusoids gradually dephase over time



**Collapse of oscillation amplitude**



## Numerical examples

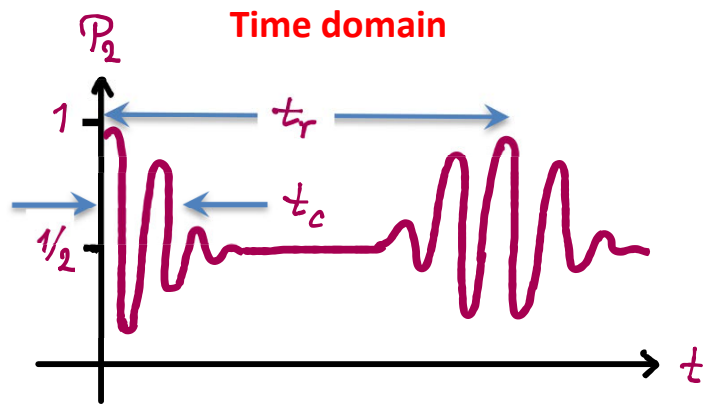
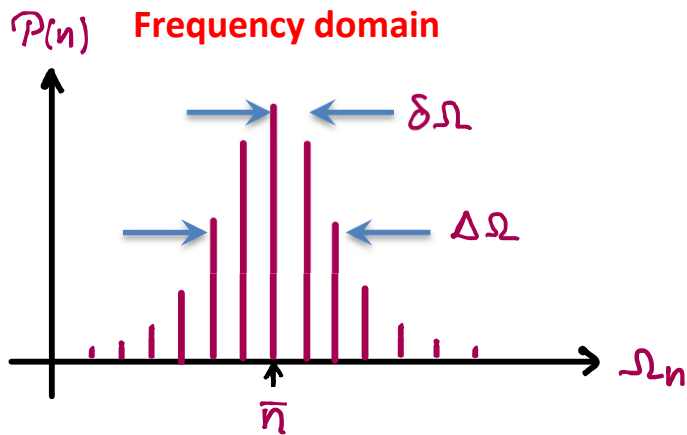


# Quantized Light – Matter Interactions

- Poisson weighted average of sinusoids
- Sinusoids gradually dephase over time



Collapse of oscillation amplitude



Use  $\Delta n = \sqrt{\bar{n}} \rightarrow \Delta\Omega \sim \Delta\Omega_{\bar{n}+\sqrt{\bar{n}}} - \Delta\Omega_{\bar{n}-\sqrt{\bar{n}}}$

$$t_c = \frac{1}{\Delta\Omega} \sim \frac{1}{2g\sqrt{\bar{n}+\sqrt{\bar{n}}} - 2g\sqrt{\bar{n}-\sqrt{\bar{n}}}} \sim \frac{1}{2g}$$

for  $\bar{n} \gg \sqrt{\bar{n}}$

Rephasing: when  $(\Omega_{\bar{n}} - \Omega_{\bar{n}-1})t_r \approx 2\pi m$

Similar arguments -> **Revival time**

$$t_r \sim \frac{2\pi}{\delta\Omega} \sim \frac{2\pi\sqrt{\bar{n}}}{g}$$

# Quantized Light – Matter Interactions

Use  $\Delta n = \sqrt{\bar{n}} \Rightarrow \Delta \Omega \sim \Delta \Omega_{\bar{n}+\sqrt{\bar{n}}} - \Delta \Omega_{\bar{n}-\sqrt{\bar{n}}}$

$$t_c = \frac{1}{\Delta \Omega} \sim \frac{1}{2g\sqrt{\bar{n}+\sqrt{\bar{n}}} - 2g\sqrt{\bar{n}-\sqrt{\bar{n}}}} \sim \frac{1}{2g}$$

for  $\bar{n} \gg \sqrt{\bar{n}}$

Rephasing: when  $(\Omega_{\bar{n}} - \Omega_{\bar{n}-1})t_r \approx 2\pi m$

Similar arguments  $\rightarrow$  **Revival time**

$$t_r \sim \frac{2\pi}{\delta \Omega} \sim \frac{2\pi\sqrt{\bar{n}}}{g}$$

## Collapse & Revival Dynamics

### Pure Quantum Phenomenon

(“graininess” of photons)

Classical limit  $\left\{ \begin{array}{l} V \rightarrow \infty \Rightarrow \mathcal{E}_h \rightarrow 0 \Rightarrow g \rightarrow 0 \\ \bar{n} \rightarrow \infty \Rightarrow \frac{\Delta \Omega}{\Omega_{\bar{n}}} \rightarrow 0 \Rightarrow \Omega_{\bar{n}} > 0 \end{array} \right. \Rightarrow t_c \rightarrow \infty$   
 well defined

$$\Omega_{\bar{n}} = 2g\sqrt{\bar{n}} = \frac{\vec{n}_{\alpha 1} \cdot 2\hat{\Sigma}_R \mathcal{E}_R \sqrt{\bar{n}}}{\hbar} = \frac{\vec{n}_{\alpha 1} \cdot \hat{\vec{E}}}{\hbar}$$

Classical Rabi frequency

mean field  $\langle \alpha(t) | \hat{\vec{E}} | \alpha(t) \rangle$



# Quantized Light – Matter Interactions

Use  $\Delta n = \sqrt{\bar{n}}$   $\rightarrow$   $\Delta \Omega \sim \Delta \Omega_{\bar{n}+\sqrt{\bar{n}}} - \Delta \Omega_{\bar{n}-\sqrt{\bar{n}}}$

$$t_c = \frac{1}{\Delta \Omega} \sim \frac{1}{2g\sqrt{\bar{n}+\sqrt{\bar{n}}} - 2g\sqrt{\bar{n}-\sqrt{\bar{n}}}} \sim \frac{1}{2g}$$

for  $\bar{n} \gg \sqrt{\bar{n}}$

Rephasing: when  $(\Omega_{\bar{n}} - \Omega_{\bar{n}-1})t_r \approx 2\pi m$

Similar arguments  $\rightarrow$  **Revival time**

$$t_r \sim \frac{2\pi}{\delta \Omega} \sim \frac{2\pi\sqrt{\bar{n}}}{g}$$

## Collapse & Revival Dynamics

### Pure Quantum Phenomenon

(“graininess” of photons)

Classical limit  $\left\{ \begin{array}{l} V \rightarrow \infty \rightarrow \xi_h \rightarrow 0 \rightarrow g \rightarrow 0 \rightarrow t_c \rightarrow \infty \\ \bar{n} \rightarrow \infty \rightarrow \frac{\Delta \Omega}{\Omega_{\bar{n}}} \rightarrow 0 \rightarrow \Omega_{\bar{n}} > 0 \end{array} \right.$  well defined

$$\Omega_{\bar{n}} = 2g\sqrt{\bar{n}} = \frac{\vec{n}_{21} \cdot 2\hat{\Sigma}_2 \xi_h \sqrt{\bar{n}}}{\hbar} = \frac{\vec{n}_{21} \cdot \vec{E}}{\hbar}$$

Classical Rabi frequency

mean field  $\langle \alpha(t) | \hat{E} | \alpha(t) \rangle$