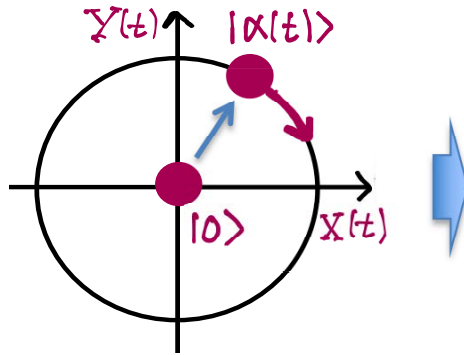


Quantum States of the Quantized Field

More about Coherent States



Coherent States
as translated
Vacuum States?

Generating Coherent States from the Vacuum

Definition: $\hat{D}(\alpha) = e^{\alpha \hat{a}^\dagger - \alpha^* \hat{a}}$

Unitary, equals translation

Glauber's formula (from BCH formula)

$$e^{\hat{A} + \hat{B}} = e^{\hat{A}} e^{\hat{B}} e^{\frac{i}{2} [\hat{A}, \hat{B}]}$$

for $[\hat{A}, [\hat{A}, \hat{B}]] = [\hat{B}, [\hat{A}, \hat{B}]] = 0$

Apply to

$$[\alpha \hat{a}^\dagger, -\alpha^* \hat{a}] = \alpha^* \alpha$$

\uparrow \uparrow \uparrow
 \hat{A} \hat{B} $[\hat{A}, \hat{B}]$

$$\hat{D}(\alpha) = e^{-|\alpha|^2/2} e^{\alpha \hat{a}^\dagger} e^{-\alpha \hat{a}}$$

Remember: $\hat{a}|0\rangle = 0$

$$e^{-\alpha \hat{a}} |0\rangle = \sum_n \frac{(-\alpha \hat{a})^n}{n!} |0\rangle = |0\rangle$$

$$\begin{aligned} \hat{D}(\alpha)|0\rangle &= e^{-|\alpha|^2/2} e^{\alpha \hat{a}^\dagger} |0\rangle \\ &= e^{-|\alpha|^2} \sum_n \frac{(\alpha \hat{a}^\dagger)^n}{n!} |0\rangle \\ &= e^{-|\alpha|^2/2} \sum_n \frac{\alpha^n}{\sqrt{n!}} |n\rangle = |\alpha\rangle \end{aligned}$$

$$\hat{D}(\alpha)|0\rangle = |\alpha\rangle$$

Quantum States of the Quantized Field

Apply to

$$[\alpha \hat{a}^\dagger, -\alpha^* \hat{a}] = \alpha^* \alpha$$

\uparrow \uparrow \uparrow
 \hat{A} \hat{B} $[\hat{A}, \hat{B}]$



$$\hat{D}(\alpha) = e^{-|\alpha|^2/2} e^{\alpha \hat{a}^\dagger} e^{-\alpha \hat{a}}$$

Remember:

$$\hat{a}|0\rangle = 0 \Rightarrow$$

$$e^{-\alpha^* \hat{a}} |0\rangle = \sum_n \frac{(-\alpha^* \hat{a})^n}{n!} |0\rangle = |0\rangle$$



$$\begin{aligned} \hat{D}(\alpha)|0\rangle &= e^{-|\alpha|^2/2} e^{\alpha \hat{a}^\dagger} |0\rangle \\ &= e^{-|\alpha|^2} \sum_n \frac{(\alpha \hat{a}^\dagger)^n}{n!} |0\rangle \\ &= e^{-|\alpha|^2/2} \sum_n \frac{\alpha^n}{\sqrt{n!}} |n\rangle = |\alpha\rangle \end{aligned}$$



$$\hat{D}(\alpha)|0\rangle = |\alpha\rangle$$

OK – $\hat{D}(\alpha)$ generates $|\alpha\rangle$ from the vacuum!

Rewrite:

$$\begin{aligned} \alpha \hat{a}^\dagger - \alpha^* \hat{a} &= (\alpha - \alpha^*) \hat{X} + i(\alpha + \alpha^*) \hat{Y} \\ &= i2\gamma \hat{X} + i2\chi \hat{Y} \end{aligned}$$

where $\chi = \langle \alpha | \hat{X} | \alpha \rangle$, $\gamma = \langle \alpha | \hat{Y} | \alpha \rangle$

Glauber's formula again:

$$\hat{D}(\alpha) = e^{i2\gamma \hat{X} + i2\chi \hat{Y}} = e^{-\chi\gamma/4} e^{i2\gamma \hat{X}} e^{i2\chi \hat{Y}}$$

Recall: $\hat{S}(q) = e^{-iq\hat{P}/\hbar} \Rightarrow$ translation by q

$\hat{S}(p) = e^{-ip\hat{Q}/\hbar} \Rightarrow$ translation by p

where

$$q = q_0 \chi, \quad p = p_0 \gamma$$

$$\hat{q} = q_0 \hat{X}, \quad \hat{p} = p_0 \hat{Y}$$

$$\& \quad \chi_0 p_0 = 2\hbar$$

Quantum States of the Quantized Field

OK – $\hat{D}(\alpha)$ generates $|\alpha\rangle$ from the vacuum!

Rewrite:

$$\alpha \hat{a}^\dagger - \alpha^* \hat{a} = (\alpha - \alpha^*) \hat{X} + i(\alpha + \alpha^*) \hat{Y} \\ = i2Y \hat{X} + i2X \hat{Y}$$

where $X = \langle \alpha | \hat{X} | \alpha \rangle$, $Y = \langle \alpha | \hat{Y} | \alpha \rangle$

Glauber's formula again:

$$\hat{D}(\alpha) = e^{i2Y \hat{X} + i2X \hat{Y}} = e^{-XY/4} e^{i2Y \hat{X}} e^{i2X \hat{Y}}$$

Recall: $\hat{S}(q) = e^{-iq\hat{p}/\hbar}$ translation by q

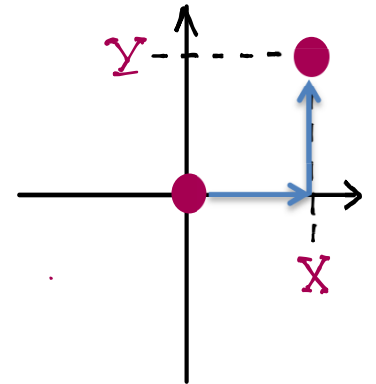
Recall: $\hat{S}(p) = e^{-ip\hat{q}/\hbar}$ translation by p

where $q = q_0 X$, $p = p_0 Y$ & $X_0 p_0 = 2\hbar$
 $\hat{q} = q_0 \hat{X}$, $\hat{p} = p_0 \hat{Y}$

This gives us

$$\hat{S}(q) = \hat{S}(X) = e^{i2X \hat{Y}}, \quad \hat{S}(p) = \hat{S}(Y) = e^{i2Y \hat{X}}$$

$\hat{D}(\alpha)$ translates
along X then Y



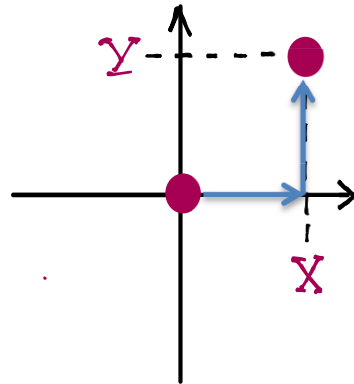
**Discussion –
How to do this?**

Quantum States of the Quantized Field

This gives us

$$\hat{S}(q) = \hat{S}(x) = e^{i2x\hat{Y}}, \quad \hat{S}(p) = \hat{S}(y) = e^{i2y\hat{X}}$$

$\hat{D}(\alpha)$ translates
along X then Y



**Discussion –
How to do this?**

Coherent States from Classical Dipole Radiation

Classical Dipole $d(t) = d_0 \cos(\omega t)$ @ $t=0$

Quantized Field $\hat{E}(z) = \epsilon_{\mathbf{k}} (\hat{a} + \hat{a}^\dagger)$

Dipole-Field Interaction

$$\hat{H} = \hbar\omega_{\mathbf{k}} (\hat{a}^\dagger \hat{a} + 1/2) + \hbar\lambda(t) (\hat{a} + \hat{a}^\dagger)$$

$$\lambda(t) = -\frac{d(t)\epsilon_{\mathbf{k}}}{\hbar} = \lambda_0 \cos(\omega t)$$

**Homework Problem
(voluntary)**

$$\alpha(T) = -i \frac{\lambda_0}{2} e^{-i(\omega - \omega_{\mathbf{k}})T/2} \frac{\sin[(\omega - \omega_{\mathbf{k}})T/2]}{(\omega - \omega_{\mathbf{k}})/2}$$

Quantum States of the Quantized Field

Coherent States from Classical Dipole Radiation

Classical Dipole $d(t) = d_0 \cos(\omega t)$ @ $t=0$

Quantized Field $\hat{E}(z) = \epsilon_k (\hat{a} + \hat{a}^\dagger)$

Dipole-Field Interaction

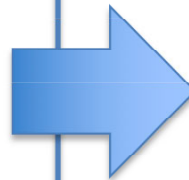
$$\hat{H} = \hbar\omega (\hat{a}^\dagger \hat{a} + 1/2) + \hbar\lambda(t) (\hat{a} + \hat{a}^\dagger)$$

$$\lambda(t) = -\frac{d(t)\epsilon_k}{\hbar} = \lambda_0 \cos(\omega t)$$



**Homework Problem
(voluntary)**

$$\alpha(T) = -i \frac{\lambda_0}{2} e^{-i(\omega - \omega_k)T/2} \frac{\sin[(\omega - \omega_k)T/2]}{(\omega - \omega_k)/2}$$



Recall from Semi-Classical Laser Theory

$$\langle \hat{j}(t) \rangle \text{ drives } \hat{E}(t)$$

classical dipole
+ quantum
fluctuations

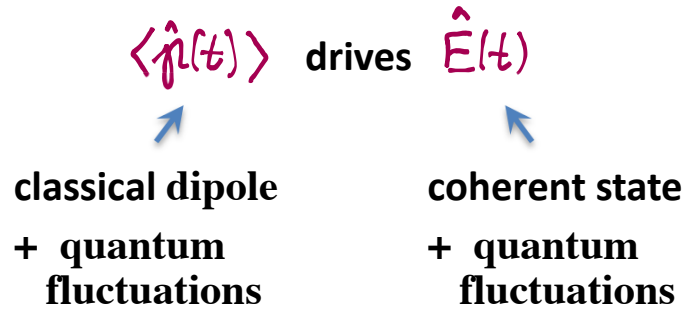
coherent state
+ quantum
fluctuations

Drive from $0 < t < T$ →

$$\alpha(t) = \alpha(T) e^{-i\omega_k(t-T)}$$

Quantum States of the Quantized Field

Recall from Semi-Classical Laser Theory

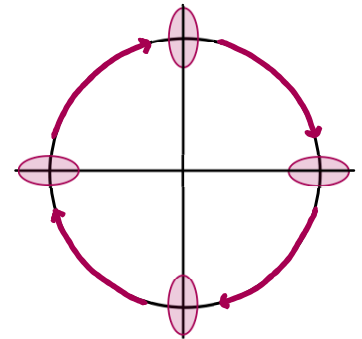


Squeezed States

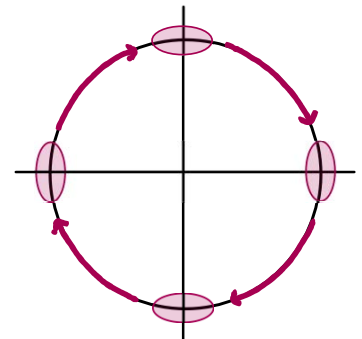
Minimum uncertainty states w/asymmetry

$$\Delta X \Delta Y = 1/4, \quad \Delta X(t) \neq \Delta Y(t)$$

Phase Squeezing



Amplitude Squeezing



Requires interaction with Nonlinear medium

Quantum States of the Quantized Field

Odds and Ends – Thermal States

$$\hat{\rho} = \sum_n P(n) |n\rangle\langle n| = \frac{1}{z} \sum_n e^{-E_n/k_B T} |n\rangle\langle n|$$

$$= (1-q) \sum_n q^n |n\rangle\langle n|, \quad q = e^{-\hbar\omega/k_B T}$$

$z = \text{Tr}[e^{-\hat{H}/k_B T}]$

Mean Photon Number:

$$\bar{n} = \text{Tr}(\hat{\rho} \hat{N}) = \sum_{k,n} \langle k | (1-q) q^n |n\rangle\langle n| \hat{N} |k\rangle$$

$$= (1-q) \sum_n n q^n = \frac{q}{1-q}$$

Photon Number Uncertainty:

$$\langle \hat{N}^2 \rangle = (1-q) \sum_n n^2 q^n = \frac{q^2 + q}{1-q}$$



$$\Delta n^2 = \langle \hat{N}^2 \rangle - \langle \hat{N} \rangle^2$$

$$= \frac{q^2 + q}{1-q} - \frac{q^2}{(1-q)^2} = \frac{q}{(1-q)^2}$$



$$\bar{n} = \frac{q}{1-q}$$

Coherent State limit

$$\Delta n = \frac{\sqrt{q}}{1-q} = \sqrt{\bar{n}(\bar{n}+1)} \geq \sqrt{\bar{n}}$$

Optical Frequencies, Room Temperature:

$$\lambda = 1 \mu\text{m}, \quad T = 300 \text{ K}$$

$$q = 6.5 \times 10^{-6}, \quad \bar{n} \sim 10^{-6}$$

Quantum States of the Quantized Field

Odds and Ends – Quantum-Classical Correspondence

Define a Translation Operator

$$\hat{T}_\alpha(t) = e^{\alpha^* e^{i\omega t} \hat{a} - \alpha e^{-i\omega t} \hat{a}^\dagger} = \hat{D}(-\alpha e^{-i\omega t})$$

Use $[\hat{a}, \hat{F}(\hat{a}^\dagger)] = dF(\hat{a}^\dagger)/d\hat{a}^\dagger$ to show

$$[\hat{a}, \hat{T}_\alpha] = \hat{a} \hat{T}_\alpha - \hat{T}_\alpha \hat{a} = -\alpha e^{-i\omega t} \hat{T}_\alpha$$

$$\Rightarrow \hat{T}_\alpha \hat{a} = \hat{a} \hat{T}_\alpha + \alpha e^{-i\omega t} \hat{T}_\alpha$$

$$\Rightarrow \hat{T}_\alpha \hat{a} \hat{T}_\alpha^\dagger = \hat{a} + \alpha e^{-i\omega t}$$

From this we get

$$\begin{aligned} \hat{E}'_\perp &= \hat{T}_\alpha \hat{E}_\perp \hat{T}_\alpha^\dagger = \hat{T}_\alpha (\epsilon_{\mathbf{k}} \hat{a} e^{i\vec{k} \cdot \vec{r}} + \text{H.C.}) \hat{T}_\alpha^\dagger \\ &= \epsilon_{\mathbf{k}} \hat{a} e^{i\vec{k} \cdot \vec{r}} + \text{H.C.} + \epsilon_{\mathbf{k}} \alpha e^{-i(\omega t - \vec{k} \cdot \vec{r})} + \text{C.C.} \\ &= \hat{E}_\perp + \hat{E}_\perp^{\text{Cl}}(\alpha, t) \end{aligned}$$

We also have $|2\rangle'(t) = \hat{T}_\alpha |\alpha(t)\rangle = |0\rangle$

Action of the unitary transformation $\hat{T}_\alpha(t)$

$$\hat{E}'_\perp = \hat{T}_\alpha(t) \hat{E}_\perp \hat{T}_\alpha(t)^\dagger = \hat{E}_\perp + \hat{E}_\perp^{\text{Cl}}(\alpha, t)$$

$$|2\rangle'(t) = \hat{T}_\alpha(t) |\alpha(t)\rangle = |0\rangle$$



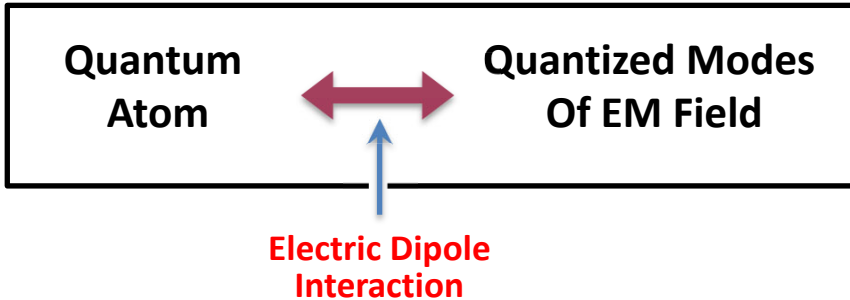
We can work with

$$\hat{E}_\perp, |\alpha(t)\rangle \quad \text{or} \quad \hat{E}_\perp + \hat{E}_\perp^{\text{Cl}}(\alpha, t), |0\rangle$$

**Validates Semiclassical Optics
for strong Coherent Fields!**

Quantized Light – Matter Interactions

General Problem:



Starting Point: System Hamiltonian

$$\hat{H} = \hat{H}_F + \hat{H}_A + \hat{H}_{AF} \quad (1)$$

$$\hat{H}_F = \sum_{\vec{k}} \hbar \omega_{\vec{k}} \left(\hat{a}_{\vec{k}}^\dagger \hat{a}_{\vec{k}} + \frac{1}{2} \right) \quad \text{Field}$$

$$\hat{H}_A = \sum_i E_i |i\rangle \langle i| = \sum_i E_i \hat{\sigma}_{ii} \quad \text{Atom}$$

$$\hat{H}_{AF} = -\hat{\vec{p}} \cdot \hat{\vec{E}}(\vec{r}, t) \quad \text{ED interaction}$$

$E_i, |i\rangle$: energies, energy levels of the atom

Dipole Operator:

$$(2) \quad \hat{\vec{p}} = \sum_{ij} \vec{p}_{ij} |i\rangle \langle j| = \sum_{ij} \vec{p}_{ij} \hat{\sigma}_{ij}$$

Field Operator:

$$\hat{\vec{E}}(\vec{r}, t) = \sum_{\vec{k}} \vec{\epsilon}_{\vec{k}} \mathcal{E}_{\vec{k}} \hat{a}_{\vec{k}} u_{\vec{k}}(\vec{r}) + \text{H.c.}, \quad \mathcal{E}_{\vec{k}} = \sqrt{\frac{\hbar \omega_{\vec{k}}}{2 \epsilon_0 V}}$$

← 2 polarization modes implicit

Pin down atom where $u_{\vec{k}}(\vec{r}) = 1$

- anywhere if $u_{\vec{k}}(\vec{r}) = e^{i\vec{k} \cdot \vec{r}}$
- if $u_{\vec{k}}(\vec{r}) = \sin(kz)$ then where $\sin(kz) = 1$



$$(3) \quad \hat{\vec{E}}(\vec{r}, t) = \hat{\vec{E}}(t) = \sum_{\vec{k}} \vec{\epsilon}_{\vec{k}} \mathcal{E}_{\vec{k}} (\hat{a}_{\vec{k}} + \hat{a}_{\vec{k}}^\dagger)$$

Quantized Light – Matter Interactions

Dipole Operator:

$$(2) \quad \hat{\vec{p}} = \sum_{i,j} \vec{p}_{ij} |i\rangle\langle j| = \sum_{i,j} \vec{p}_{ij} \hat{\sigma}_{ij}$$

Field Operator:

$$\vec{E}(\vec{r}, t) = \sum_{\vec{k}} \vec{\epsilon}_{\vec{k}} \mathcal{E}_{\vec{k}} \hat{a}_{\vec{k}} u_{\vec{k}}(\vec{r}) + \text{H.c.}, \quad \mathcal{E}_{\vec{k}} = \sqrt{\frac{\hbar \omega_{\vec{k}}}{2\epsilon_0 V}}$$

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$$(3) \quad \vec{E}(\vec{r}, t) = \vec{E}(t) = \sum_{\vec{k}} \vec{\epsilon}_{\vec{k}} \mathcal{E}_{\vec{k}} (\hat{a}_{\vec{k}} + \hat{a}_{\vec{k}}^\dagger)$$

Combining (2) & (3):

$$\begin{aligned} \hat{H}_{AF} &= \sum_{i,j} \sum_{\vec{k}} -\vec{p}_{ij} \cdot \vec{\epsilon}_{\vec{k}} \mathcal{E}_{\vec{k}} \hat{\sigma}_{ij} (\hat{a}_{\vec{k}} + \hat{a}_{\vec{k}}^\dagger) \\ &= \sum_{i,j} \sum_{\vec{k}} \hbar g_{\vec{k}}^{(ij)} (\hat{a}_{\vec{k}} + \hat{a}_{\vec{k}}^\dagger) \end{aligned}$$

where $g_{\vec{k}}^{(ij)} = \frac{\vec{p}_{ij} \cdot \vec{\epsilon}_{\vec{k}} \mathcal{E}_{\vec{k}}}{\hbar}$

↑ Rabi Freq., note sign convention

2-level atom $\rightarrow (i,j) = (1,2)$:

$$\hat{H}_{AF} = \hbar \sum_{\vec{k}} (g_{\vec{k}} \hat{\sigma}_{21} + g_{\vec{k}}^* \hat{\sigma}_{12}) (\hat{a}_{\vec{k}} + \hat{a}_{\vec{k}}^\dagger)$$

Define:

$$\hat{\sigma}_+ = \hat{\sigma}_{21} = |2\rangle\langle 1|$$

$$\hat{\sigma}_- = \hat{\sigma}_{12} = |1\rangle\langle 2|$$

$$\hat{\sigma}_z = \hat{\sigma}_{22} - \hat{\sigma}_{11} = |2\rangle\langle 2| - |1\rangle\langle 1|$$

Pauli matrices

$$\hat{\sigma}_x = \frac{1}{2} (\hat{\sigma}_+ + \hat{\sigma}_-)$$

$$\hat{\sigma}_y = \frac{1}{2i} (\hat{\sigma}_+ - \hat{\sigma}_-)$$

$$\hat{\sigma}_z$$

Quantized Light – Matter Interactions

Combining (2) & (3):

$$\hat{H}_{AF} = \sum_{ij} \sum_{\mathbf{k}} -\vec{n}_{ij} \cdot \vec{\epsilon}_{\mathbf{k}} \mathcal{E}_{\mathbf{k}} \hat{\sigma}_{ij} (\hat{a}_{\mathbf{k}} + \hat{a}_{\mathbf{k}}^{\dagger})$$

$$= \sum_{ij} \sum_{\mathbf{k}} \hbar g_{\mathbf{k}}^{(ij)} (\hat{a}_{\mathbf{k}} + \hat{a}_{\mathbf{k}}^{\dagger})$$

where $g_{\mathbf{k}}^{(ij)} = \frac{\vec{n}_{ij} \cdot \vec{\epsilon}_{\mathbf{k}} \mathcal{E}_{\mathbf{k}}}{\hbar}$

Rabi Freq., note sign convention

2-level atom $\rightarrow (i,j) = (1,2) :$

$$\hat{H}_{AF} = \hbar \sum_{\mathbf{k}} (g_{\mathbf{k}} \hat{\sigma}_{21} + g_{\mathbf{k}}^* \hat{\sigma}_{12}) (\hat{a}_{\mathbf{k}} + \hat{a}_{\mathbf{k}}^{\dagger})$$

Define:

Pauli matrices

$$\left. \begin{aligned} \hat{\sigma}_+ &= \hat{\sigma}_{21} = |2\rangle\langle 1| \\ \hat{\sigma}_- &= \hat{\sigma}_{12} = |1\rangle\langle 2| \\ \hat{\sigma}_z &= \hat{\sigma}_{22} - \hat{\sigma}_{11} = |2\rangle\langle 2| - |1\rangle\langle 1| \end{aligned} \right\} \rightarrow \begin{aligned} \hat{\sigma}_x &= \frac{1}{2} (\hat{\sigma}_+ + \hat{\sigma}_-) \\ \hat{\sigma}_y &= \frac{1}{2i} (\hat{\sigma}_+ - \hat{\sigma}_-) \\ \hat{\sigma}_z & \end{aligned}$$

With this notation

(4)

$$\hat{H}_{AF} = \hbar \sum_{\mathbf{k}} (g_{\mathbf{k}} \hat{\sigma}_+ \hat{a}_{\mathbf{k}} + \cancel{g_{\mathbf{k}} \hat{\sigma}_+ \hat{a}_{\mathbf{k}}^{\dagger}} + \cancel{g_{\mathbf{k}} \hat{\sigma}_- \hat{a}_{\mathbf{k}}} + g_{\mathbf{k}} \hat{\sigma}_- \hat{a}_{\mathbf{k}}^{\dagger})$$

Energy conservation?



$$\hat{H}_{AF} = \hbar \sum_{\mathbf{k}} (g_{\mathbf{k}} \hat{\sigma}_+ \hat{a}_{\mathbf{k}} + g_{\mathbf{k}} \hat{\sigma}_- \hat{a}_{\mathbf{k}}^{\dagger})$$