

Quantum Electrodynamics – QED

Starting point: Maxwells Equations

- (1) $\nabla \cdot \vec{E}(\vec{r}, t) = \frac{1}{\epsilon_0} \rho(\vec{r}, t)$
- (2) $\nabla \cdot \vec{B}(\vec{r}, t) = 0$
- (3) $\nabla \times \vec{E}(\vec{r}, t) = -\frac{\partial}{\partial t} \vec{B}(\vec{r}, t)$
- (4) $\nabla \times \vec{B}(\vec{r}, t) = \frac{1}{c^2} \frac{\partial}{\partial t} \vec{E}(\vec{r}, t) + \frac{1}{\epsilon_0 c^2} \vec{j}(\vec{r}, t)$

Implicit: Charges & Fields in Vacuum
No “medium response”

Same issue as with our introductory example:
Maxwells eqs are non-local



We need to put the classical description
in proper form -> Normal Mode expansion

Free Fields - Switch to Fourier Domain

- (1) $i\vec{k} \cdot \vec{E}(\vec{k}, t) = \frac{1}{\epsilon_0} \rho(\vec{k}, t)$
- (2) $i\vec{k} \cdot \vec{B}(\vec{k}, t) = 0$
- (3) $i\vec{k} \times \vec{E}(\vec{k}, t) = -\frac{\partial}{\partial t} \vec{B}(\vec{k}, t)$
- (4) $i\vec{k} \times \vec{B}(\vec{k}, t) = \frac{1}{c^2} \frac{\partial}{\partial t} \vec{E}(\vec{k}, t) + \frac{1}{\epsilon_0 c^2} \vec{j}(\vec{k}, t)$

Fourier Transform: $\left\{ \begin{array}{l} \nabla \cdot \vec{G} \Leftrightarrow i\vec{k} \cdot \vec{G} \\ \nabla \times \vec{G} \Leftrightarrow i\vec{k} \times \vec{G} \end{array} \right.$

Note: This is a Normal Mode decomposition

No charges -> No coupling between modes
with different \vec{k}

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Fourier Transform:
$$\begin{cases} \nabla \cdot \vec{G} \Leftrightarrow i\vec{k} \cdot \vec{h} \\ \nabla \times \vec{G} \Leftrightarrow i\vec{k} \times \vec{h} \end{cases}$$

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No charges \rightarrow No coupling between modes with different \vec{k}

Separate into Transverse & Longitudinal Fields

$$\vec{E}(\vec{k}, t) = \vec{E}_{||}(\vec{k}, t) + \vec{E}_{\perp}(\vec{k}, t)$$

$$\vec{B}(\vec{k}, t) = \cancel{\vec{B}_{||}(\vec{k}, t)} + \vec{B}_{\perp}(\vec{k}, t) \quad \text{MEq (2)}$$

\uparrow Entirely Transverse

Note:
$$\begin{cases} -\frac{i}{k} i\vec{k} \cdot \vec{E} \text{ is the projection of } \vec{E} \text{ onto } \vec{k} \\ \vec{E}_{||} \text{ is } \frac{\vec{k}}{k} \times \text{ the projection of } \vec{E} \text{ onto } \vec{k} \end{cases}$$

\Downarrow MEq (1)

$$\vec{E}_{||} = \frac{\vec{k}}{k} \epsilon_{||} = \frac{\vec{k}}{k} \left(-\frac{i}{k} i\vec{k} \cdot \vec{E} \right) = -i \frac{\vec{k}}{\epsilon_0 k^2} \rho(\vec{k}, t)$$

Coulomb field from the charges

\Downarrow

Only \vec{E}_{\perp} and \vec{B}_{\perp} are new degrees of freedom beyond the particles \rightarrow Free Fields

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Eqs for Transverse Fields, from MEqs (3) & (4)

$$(5a) \quad \frac{\partial}{\partial t} \vec{B}(\vec{k}, t) = -i \vec{k} \times \vec{E}_{\perp}(\vec{k}, t)$$

$$(6a) \quad \frac{\partial}{\partial t} \vec{E}_{\perp}(\vec{k}, t) = c^2 i \vec{k} \times \vec{B}(\vec{k}, t) - \frac{1}{\epsilon_0} \vec{j}_{\perp}(\vec{k}, t)$$

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inverse FT

$$(5b) \quad \frac{\partial}{\partial t} \vec{B}(\vec{r}, t) = -\nabla \times \vec{E}_{\perp}(\vec{r}, t)$$

$$(6b) \quad \frac{\partial}{\partial t} \vec{E}_{\perp}(\vec{r}, t) = c^2 \nabla \times \vec{B}(\vec{r}, t) - \frac{1}{\epsilon_0} \vec{j}_{\perp}(\vec{r}, t)$$

combine (5b) & (6b)

Wave Equation for the Free Fields

$$\left(\nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) \vec{E}_{\perp}(\vec{r}, t) = \frac{1}{\epsilon_0 c^2} \frac{\partial}{\partial t^2} \vec{j}_{\perp}(\vec{r}, t)$$

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Wave Equation for the Free Fields

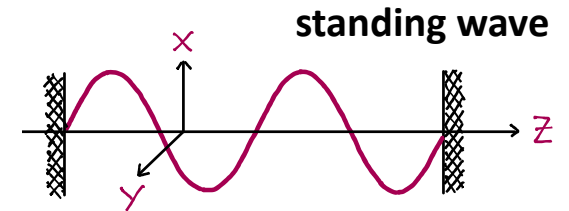
$$\left(\nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) \vec{E}_{\perp}(\vec{r}, t) = \frac{1}{\epsilon_0 c^2} \frac{\partial}{\partial t^2} \vec{j}_{\perp}(\vec{r}, t)$$

Normal Modes in a 1D Cavity

Length L

Cross section A

Volume $V = LA$



Normal Modes are Standing Waves

Let $\vec{E}(z, t) = \vec{e}_x E_x(z, t)$ and expand

fiducial mass

$$(7) \quad E_x(z, t) = \sum_j A_j q_j(t) \sin(k_j z), \quad A_j = \sqrt{\frac{\omega_j^2 m_j}{2 \epsilon_0 V}}$$

MEq (4) w/no charges

$$\nabla \times \vec{B} = \frac{1}{c^2} \frac{\partial}{\partial t} \vec{E}_{\perp}(\vec{r}, t) = \vec{e}_x \frac{1}{c^2} \sum_j A_j \dot{q}_j(t) \sin(k_j z)$$

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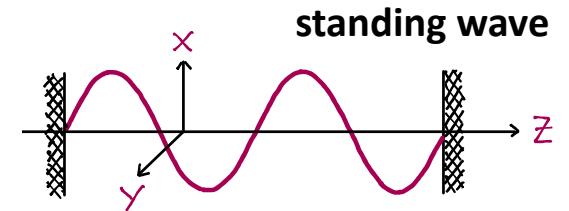
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\vec{B} transverse $\Rightarrow B_z = 0$

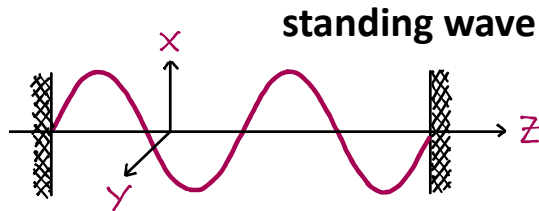
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From Eq. (5a) we see that

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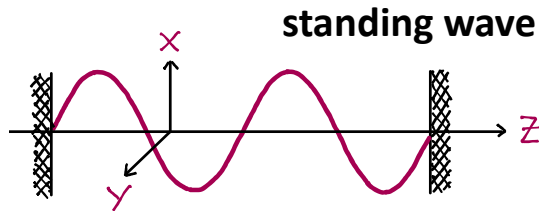
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Putting this together we get

$$\frac{\partial B_y}{\partial z} = - \sum_j \frac{A_j}{c^2} \dot{q}_j(t) \sin(k_j z)$$



$$(8) B_y(z,t) = \sum_j \frac{A_j}{k_j c^2} \dot{q}_j(t) \cos(k_j z)$$

Hamiltonian (Energy) for the Classical Field

$$\begin{aligned} \mathcal{H} &= \frac{\epsilon_0 A}{2} \int_0^L dz (|\vec{E}|^2 + c^2 |\vec{B}|^2) = \\ &= \frac{\epsilon_0 A}{2} \int_0^L dz \sum_j \left[A_j^2 \dot{q}_j(t)^2 \sin^2(k_j z) + \frac{A_j^2}{k_j^2} \dot{q}_j(t)^2 \cos^2(k_j z) \right] \end{aligned}$$

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Integrating over the Cavity volume

$$A \int_0^L dz \sin^2(k_j z) = A \int_0^L dz \cos^2(k_j z) = V/2$$

and substituting $A_j^2 = \frac{\omega_j^2 m_j}{2\epsilon_0 V}$ we finally get

$$\mathcal{H} = \sum_j \left[\frac{1}{2} m_j \omega_j^2 q_j^2 + \frac{1}{2} m_j \dot{q}_j^2 \right]$$

Lagrangian for the Classical Field

$$\mathcal{L} = \frac{\epsilon_0 A}{2} \int_0^L dz (c^2 |\vec{B}|^2 - |\vec{E}|^2) \quad \checkmark$$

$$= \sum_j \left[\frac{1}{2} m_j \dot{q}_j^2 - \frac{1}{2} m_j \omega_j^2 q_j^2 \right]$$

Check $\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}_j} - \frac{\partial \mathcal{L}}{\partial q_j} = 0 \rightarrow \ddot{q}_j + \omega_j^2 q_j = 0$

$(\nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2}) \vec{E}_\perp(\vec{r}, t) = 0 \rightarrow \ddot{q}_j + \omega_j^2 q_j = 0$

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And Finally:

$$p_j = \frac{\partial \mathcal{L}}{\partial \dot{q}_j} = m_j \dot{q}_j$$

**As before, a collection
of Harmonic Oscillators,
ready for quantization!**

End 03-31-2021