

**Problem 6)** This problem is similar to that discussed in Sec.6, the only difference being that the radial solution  $f(r)$  is now allowed to contain Bessel functions of the 2<sup>nd</sup> kind in addition to those of the 1<sup>st</sup> kind. We thus have

$$f(r) = J_m(cr/v) + \alpha Y_m(cr/v), \quad (1)$$

where  $c$  is the separation constant, as before, whereas  $\alpha$  is a new constant coefficient which specifies the relative contributions of the two Bessel functions in each vibrational mode,  $m$ . The boundary conditions now demand that  $f(R_1) = 0$  and also  $f(R_2) = 0$ . Consequently,

$$\alpha = -\frac{J_m(cR_1/v)}{Y_m(cR_1/v)} = -\frac{J_m(cR_2/v)}{Y_m(cR_2/v)}. \quad (2)$$

The separation constant  $c$  must ensure the equality of the two  $J_m/Y_m$  ratios appearing in the above equation. The acceptable values of  $c$  are, therefore, determined numerically, by searching for those values that satisfy the equality of these two ratios. There will be an infinite number of (discrete) values,  $c_{mn}$ , that will satisfy the above requirement for each and every mode,  $m$ . The corresponding values of  $\alpha$  must then be designated  $\alpha_{mn}$ , and obtained from Eq.(2). The final solution for the annular membrane will, therefore, have the same general form as that given by Eq.(28), except for  $J_m(r_{mn}r/R)$  being replaced by  $J_m(c_{mn}r/v) + \alpha_{mn}Y_m(c_{mn}r/v)$ . The oscillation frequency of the  $mn^{\text{th}}$  mode will then be given by  $\omega_{mn} = \sqrt{c_{mn}^2 - (\gamma/2)^2}$ .

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