

Problem 2) For a capacitor  $Q = CV \Rightarrow \frac{dQ}{dt} = C \frac{dV}{dt} \Rightarrow I = C \frac{dV}{dt}$ .

For a resistor  $V = RI$ . Therefore,  $V(t) = RI(t) + V_c(t) \Rightarrow$

$$RC \frac{dV_c(t)}{dt} + V_c(t) = V(t).$$

a) Because  $\text{step}(t) = 0$  from  $t = -\infty$  to  $0$ , the capacitor is empty during the time interval  $t < 0$ . Therefore, we need not worry about the

initial voltage of the capacitor. Fourier Transforming the differential

equation we find:  $RC (i2\pi s) \underset{\text{F.T. of } V_c(t)}{V_c(s)} + V_c(s) = \mathcal{F}\{V(t)\} \Rightarrow$

$$\underset{\sim}{V}_c(s) = \frac{V_0 \mathcal{F}\{\text{step}(t)\}}{1 + i2\pi RCs} = \frac{\frac{1}{2} V_0 \delta(s) - i \frac{V_0}{2\pi s}}{1 + i2\pi RCs} = \frac{V_0 \delta(s)}{2(1 + i2\pi RCs)} - i \frac{V_0}{2\pi s(1 + i2\pi RCs)}$$

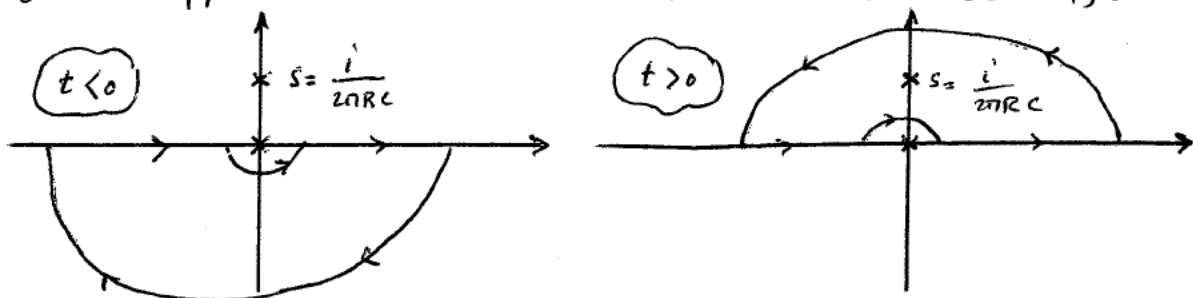
The function that multiplies  $\delta(s)$  may be replaced by its value at  $s = 0$ .

We thus have:

$$\underset{\sim}{V}_c(s) = \frac{1}{2} V_0 \delta(s) - \frac{V_0}{4\pi^2 RC} \frac{1}{s(s - \frac{i}{2\pi RC})}$$

$$\begin{aligned} V_c(t) &= \mathcal{F}^{-1}\{\underset{\sim}{V}_c(s)\} = \int_{-\infty}^{\infty} \underset{\sim}{V}_c(s) e^{+i2\pi st} ds = \frac{1}{2} V_0 \int_{-\infty}^{\infty} \delta(s) e^{i2\pi st} ds - \frac{V_0}{4\pi^2 RC} \int_{-\infty}^{\infty} \frac{e^{i2\pi st} ds}{s(s - \frac{i}{2\pi RC})} \\ &= \frac{1}{2} V_0 - \frac{V_0}{4\pi^2 RC} \int_{-\infty}^{\infty} \frac{e^{i2\pi st}}{s(s - \frac{i}{2\pi RC})} ds. \end{aligned}$$

The integral must be evaluated on the lower semi-circle when  $t < 0$ , and on the upper semi-circle when  $t > 0$  (Jordan's lemma).



For  $t < 0$  the integral will be <sup>minus</sup> one-half of the residue at  $s=0$ . That is,

$$\int_{-\infty}^{\infty} \frac{e^{i2\pi st}}{s(s - \frac{i}{2\pi RC})} ds = -\frac{1}{2} (2\pi i) \frac{1}{0 - \frac{i}{2\pi RC}} = 2\pi^2 RC \quad \leftarrow t < 0$$

because  $s=0$   
is on-axis

Therefore,  $V_c(t) = \frac{1}{2} V_0 - \frac{V_0}{4\pi^2 RC} (2\pi^2 RC) = \frac{1}{2} V_0 - \frac{1}{2} V_0 = 0 \quad \leftarrow t < 0$

For  $t > 0$ , the integral will be equal to one-half of the residue at  $s=0$ , plus the residue at  $s = \frac{i}{2\pi RC}$ , that is,

$$\int_{-\infty}^{\infty} \frac{e^{i2\pi st}}{s(s - \frac{i}{2\pi RC})} ds = 2\pi i \left\{ \frac{1}{2} \frac{1}{0 - \frac{i}{2\pi RC}} + \frac{e^{-t/RC}}{i/2\pi RC} \right\}$$

$$= -2\pi^2 RC + 4\pi^2 RC e^{-t/RC} \quad \leftarrow t \geq 0$$

Therefore,  $V_c(t) = \frac{1}{2} V_0 - \frac{V_0}{4\pi^2 RC} [-2\pi^2 RC + 4\pi^2 RC e^{-t/RC}] \Rightarrow$

$$V_c(t) = V_0 (1 - e^{-t/RC}) \quad \leftarrow t \geq 0$$

Combining the above results we'll have  $V_c(t) = V_0 (1 - e^{-t/RC}) \text{step}(t)$ .

b) When  $V(t) = \delta(t)$  we'll have  $\mathcal{F}\{V(t)\} = 1$ . Therefore,

$$V_c(s) = \frac{1}{1 + i2\pi RCs} \Rightarrow V_c(t) = \mathcal{F}^{-1}\{V_c(s)\} = \int_{-\infty}^{\infty} \frac{e^{i2\pi st}}{1 + i2\pi RCs} ds$$

$$= \frac{1}{i2\pi RC} \int_{-\infty}^{\infty} \frac{e^{i2\pi st}}{s - \frac{i}{2\pi RC}} ds$$

The only pole here is at  $s = \frac{i}{2\pi RC}$ . For  $t < 0$  the integral must be

done on the semi-circle in the lower half-plane (to satisfy Jordan's lemma). But here there are no poles and, therefore, the integral will be zero. For  $t > 0$ , the integral is done over the semi-circle in the upper half-plane, where the only enclosed pole is  $s = \frac{-t}{RC}$ . Consequently,

$$V_c(t) = \frac{1}{i2\pi RC} (2\pi i) \frac{e^{-t/RC}}{1} = +\frac{1}{RC} e^{-t/RC}; \quad t > 0.$$

Combining the above results, we'll have  $V_c(t) = \frac{1}{RC} e^{-t/RC} \text{Step}(t)$ .

Alternatively, since  $\delta(t)$  is the derivative of  $\text{Step}(t)$ , the impulse-response could be found directly by differentiating the step-response (setting  $V_0 = 1$ ). Therefore,

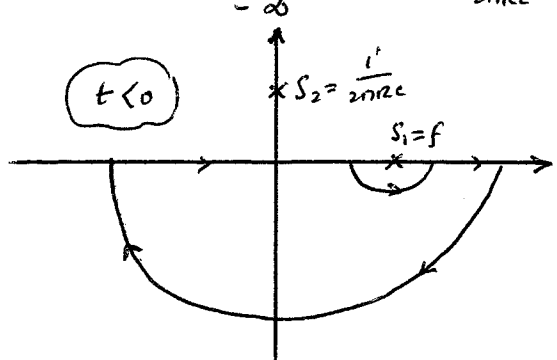
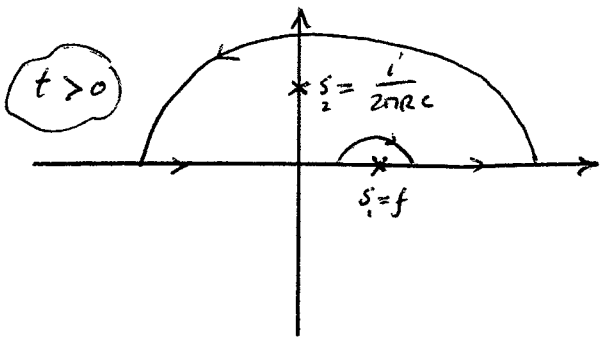
$$\text{Impulse response} = \frac{d}{dt} [(1 - e^{-t/RC}) \text{Step}(t)] = \frac{1}{RC} e^{-t/RC} \text{Step}(t) + (1 - e^{-t/RC}) \delta(t).$$

The coefficient of  $\delta(t)$  in the above expression is equal to zero at  $t = 0$ . Therefore, the last term drops out, and what's left is identical with the result obtained previously.

$$\begin{aligned} \text{c) } \mathcal{F}\{V_c(t)\} &= \mathcal{F}\{V_0 e^{i2\pi ft} \text{Step}(t)\} = V_0 \delta(s-f) * \left\{ \frac{1}{2} \delta(s) - \frac{i}{2\pi s} \right\} \\ &= \frac{1}{2} V_0 \delta(s-f) - \frac{iV_0}{2\pi(s-f)}. \end{aligned}$$

$$V_c(s) = \frac{\frac{1}{2} V_0 \delta(s-f) - \frac{iV_0}{2\pi(s-f)}}{1 + i2\pi RCs} = \frac{\frac{1}{2} V_0}{1 + i2\pi RCf} \delta(s-f) - \frac{V_0}{4\pi^2 RC} \frac{1}{(s-f)(s - \frac{i}{2\pi RC})}$$

Therefore,  $V_c(t) = \int_{-\infty}^{\infty} \frac{V_0}{2(1+i2\pi Rcf)} \delta(s-f) e^{+i2\pi st} ds = \frac{V_0}{4\pi^2 RC} \int_{-\infty}^{\infty} \frac{e^{+i2\pi st}}{(s-f)(s-\frac{i}{2\pi RC})} ds$



For  $t < 0$ , the second integral is equal to minus one-half of the residue at  $s_1 = f$ , that is,

$$V_c(t) = \frac{V_0 e^{i2\pi ft}}{2(1+i2\pi Rcf)} + \frac{V_0}{4\pi^2 RC} \frac{1}{2}(2\pi i) \frac{e^{i2\pi ft}}{f - \frac{i}{2\pi RC}}$$

$$= \frac{V_0 e^{i2\pi ft}}{2(1+i2\pi Rcf)} - \frac{V_0 e^{i2\pi ft}}{2(1+i2\pi Rcf)} = 0 \quad \leftarrow t < 0$$

For  $t > 0$ , the second integral is taken over the upper semi-circle. Its value will be given by one-half the residue at  $s_1 = f$  plus the residue at  $s_2 = \frac{i}{2\pi RC}$ . We'll have:

$$V_c(t) = \frac{V_0 e^{i2\pi ft}}{2(1+i2\pi Rcf)} - \frac{V_0}{4\pi^2 RC} (2\pi i) \left\{ \frac{e^{i2\pi ft}}{2(f - \frac{i}{2\pi RC})} + \frac{e^{-t/RC}}{\frac{i}{2\pi RC} - f} \right\}$$

$$= \frac{V_0 e^{i2\pi ft}}{1+i2\pi Rcf} - V_0 \frac{e^{-t/RC}}{1+i2\pi Rcf} = \frac{V_0}{1+i2\pi Rcf} (e^{i2\pi ft} - e^{-t/RC}) \quad \leftarrow t > 0$$

Note that  $V_c(0^+) = 0$ , which is a required boundary condition for the first-order differential equation that governs the evolution of  $V_c(t)$ .

The above results may now be combined to yield:

$$V_c(t) = \frac{V_0}{1+i2\pi Rcf} (e^{i2\pi ft} - e^{-t/RC}) \text{step}(t).$$