

PhD Qualifying Exam, Fall 2021

Opti 501

Solution to Problem 1)

- a) $\boldsymbol{\sigma} \cdot \boldsymbol{\sigma} = 1 \rightarrow (i\sigma_x \hat{\mathbf{x}} + \sigma_z \hat{\mathbf{z}}) \cdot (i\sigma_x \hat{\mathbf{x}} + \sigma_z \hat{\mathbf{z}}) = -\sigma_x^2 + \sigma_z^2 = 1 \rightarrow \sigma_z^2 = 1 + \sigma_x^2.$
- b) $\boldsymbol{\nabla} \cdot \mathbf{E} = 0 \rightarrow \boldsymbol{\sigma} \cdot \mathbf{E}_0 = 0 \rightarrow i\sigma_x E_{x0} + \sigma_z E_{z0} = 0.$
- c) $\boldsymbol{\nabla} \cdot \mathbf{B} = 0 \rightarrow \boldsymbol{\sigma} \cdot \mathbf{H}_0 = 0 \rightarrow i\sigma_x H_{x0} + \sigma_z H_{z0} = 0.$
- d) $\boldsymbol{\nabla} \times \mathbf{E} = -\partial \mathbf{B} / \partial t \rightarrow ik_0 \boldsymbol{\sigma} \times \mathbf{E}_0 = i\omega \mu_0 \mathbf{H}_0 \rightarrow (i\sigma_x \hat{\mathbf{x}} + \sigma_z \hat{\mathbf{z}}) \times (E_{x0} \hat{\mathbf{x}} + E_{y0} \hat{\mathbf{y}} + E_{z0} \hat{\mathbf{z}}) = Z_0 \mathbf{H}_0$
 $\rightarrow i\sigma_x E_{y0} \hat{\mathbf{z}} + (\sigma_z E_{x0} - i\sigma_x E_{z0}) \hat{\mathbf{y}} - \sigma_z E_{y0} \hat{\mathbf{x}} = Z_0 \mathbf{H}_0$
 $\rightarrow Z_0 H_{0x} = -\sigma_z E_{y0}; \quad Z_0 H_{0y} = \sigma_z E_{x0} - i\sigma_x E_{z0}; \quad Z_0 H_{0z} = i\sigma_x E_{y0}.$
- e) If $E_{z0} = 0$, then from (b) we have $E_{x0} = 0$, and from (d) we find $H_{y0} = 0$.
- f) If $H_{z0} = 0$, then from (c) we have $H_{x0} = 0$, and from (d) we find $E_{y0} = 0$.
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Solution to Problem 2)

- a) $\mathbf{E}(z, t) = E_{x0} \cos(k_0 z - \omega t + \varphi_0) \hat{\mathbf{x}}$ and $\mathbf{H}(z, t) = (E_{x0}/Z_0) \cos(k_0 z - \omega t + \varphi_0) \hat{\mathbf{y}}$. Note that \mathbf{E} and \mathbf{H} have the same phase φ_0 .

$\mathbf{E}'(z, t) = E_{x0} \cos(k_0 z + \omega t + \varphi'_0) \hat{\mathbf{x}}$ and $\mathbf{H}'(z, t) = -(E_{x0}/Z_0) \cos(k_0 z + \omega t + \varphi'_0) \hat{\mathbf{y}}$. Again, \mathbf{E}' and \mathbf{H}' have the same phase φ'_0 , although it could differ from φ_0 .

$$\begin{aligned} \text{Total } E\text{-field: } \mathbf{E}(z, t) + \mathbf{E}'(z, t) &= E_{x0} [\cos(k_0 z - \omega t + \varphi_0) + \cos(k_0 z + \omega t + \varphi'_0)] \hat{\mathbf{x}} \\ &= 2E_{x0} \cos[k_0 z + \frac{1}{2}(\varphi'_0 + \varphi_0)] \cos[\omega t + \frac{1}{2}(\varphi'_0 - \varphi_0)] \hat{\mathbf{x}}. \end{aligned}$$

$$\begin{aligned} \text{Total } H\text{-field: } \mathbf{H}(z, t) + \mathbf{H}'(z, t) &= (E_{x0}/Z_0) [\cos(k_0 z - \omega t + \varphi_0) - \cos(k_0 z + \omega t + \varphi'_0)] \hat{\mathbf{y}} \\ &= 2(E_{x0}/Z_0) \sin[k_0 z + \frac{1}{2}(\varphi'_0 + \varphi_0)] \sin[\omega t + \frac{1}{2}(\varphi'_0 - \varphi_0)] \hat{\mathbf{y}}. \end{aligned}$$

Since the total E -field at the mirror surfaces must be zero, we will have

$$\text{First mirror surface at } z = 0: \quad 2E_{x0} \cos[\frac{1}{2}(\varphi'_0 + \varphi_0)] \cos[\omega t + \frac{1}{2}(\varphi'_0 - \varphi_0)] \hat{\mathbf{x}} = 0.$$

$$\text{Second mirror surface at } z = d: \quad 2E_{x0} \cos[k_0 d + \frac{1}{2}(\varphi'_0 + \varphi_0)] \cos[\omega t + \frac{1}{2}(\varphi'_0 - \varphi_0)] \hat{\mathbf{x}} = 0.$$

Consequently,

$$\begin{cases} \cos[\frac{1}{2}(\varphi'_0 + \varphi_0)] = 0 \\ \cos[k_0 d + \frac{1}{2}(\varphi'_0 + \varphi_0)] = 0 \end{cases} \rightarrow \begin{cases} \varphi_0 + \varphi'_0 = (2n + 1)\pi & \text{(odd multiple of } \pi); \\ k_0 d = m\pi \rightarrow d = m\lambda_0/2 & \text{(integer multiple of } \lambda_0/2). \end{cases}$$

The total fields in the cavity are thus found to be

$$\begin{aligned} \mathbf{E}(z, t) &= 2E_{x0} \sin(k_0 z) \cos[\omega t + \frac{1}{2}(\varphi'_0 - \varphi_0)] \hat{\mathbf{x}}. \\ \mathbf{H}(z, t) &= -2(E_{x0}/Z_0) \cos(k_0 z) \sin[\omega t + \frac{1}{2}(\varphi'_0 - \varphi_0)] \hat{\mathbf{y}}. \end{aligned}$$

b) The surface-current-density equals the total tangential H -field at each mirror surface. For the first mirror at $z = 0$, $\cos(k_0 z) = 1$. For the second mirror at $z = d$, $\cos(k_0 z) = \cos(m\pi) = \pm 1$. The magnitude of the surface-current-density on both mirrors is, therefore, $J_{s0} = 2E_{x0}/Z_0$.

c) The trapped energy per unit cross-sectional area is given by

Trapped E -field energy:

$$\begin{aligned} \frac{1}{2}\epsilon_0 \int_0^d |\mathbf{E}(z, t)|^2 dz &= 2\epsilon_0 E_{x0}^2 \cos^2[\omega t + \frac{1}{2}(\varphi'_0 - \varphi_0)] \int_0^d \sin^2(k_0 z) dz \\ &= \epsilon_0 E_{x0}^2 d \cos^2[\omega t + \frac{1}{2}(\varphi'_0 - \varphi_0)]. \end{aligned}$$

Trapped H -field energy:

$$\begin{aligned} \frac{1}{2}\mu_0 \int_0^d |\mathbf{H}(z, t)|^2 dz &= 2\mu_0 (E_{x0}/Z_0)^2 \sin^2[\omega t + \frac{1}{2}(\varphi'_0 - \varphi_0)] \int_0^d \cos^2(k_0 z) dz \\ &= \epsilon_0 E_{x0}^2 d \sin^2[\omega t + \frac{1}{2}(\varphi'_0 - \varphi_0)]. \end{aligned}$$

The peak values of E -field and H -field energies (per unit cross-sectional area) are thus equal to $\epsilon_0 E_{x0}^2 d$. However, there exists a phase difference between these two entities: When the E -field energy is zero, the H -field energy is at a maximum, and vice-versa. At one instant, all the energy is in the E -field; a quarter of a period later, all the energy is in the H -field. The energy thus swings back and forth from one form to the other.

d) $\mathbf{S}(z, t) = \mathbf{E}(z, t) \times \mathbf{H}(z, t)$

$$\begin{aligned} &= -(4E_{x0}^2/Z_0) \sin(k_0 z) \cos(k_0 z) \sin[\omega t + \frac{1}{2}(\varphi'_0 - \varphi_0)] \cos[\omega t + \frac{1}{2}(\varphi'_0 - \varphi_0)] \hat{\mathbf{z}} \\ &= -(E_{x0}^2/Z_0) \sin(2k_0 z) \sin[2\omega t + (\varphi'_0 - \varphi_0)] \hat{\mathbf{z}}. \end{aligned}$$

At the nodes of the E -field, as well as those of the H -field, the Poynting vector is zero. No energy, therefore, crosses these nodes. In between the nodes, the energy flows to the right for one quarter of one oscillation period ($T = 2\pi/\omega$), then flows to the left during the next quarter. The process is then repeated.
