

Please write your name and ID number on all the pages, then staple them together.

Answer all the questions.

Note: Bold symbols represent vectors and vector fields.

1. The vector and scalar potentials $\mathbf{A}(\mathbf{r}, t)$ and $\psi(\mathbf{r}, t)$ are obtained from the current and charge density distributions $\mathbf{J}(\mathbf{r}, t)$ and $\rho(\mathbf{r}, t)$ following standard procedure. Accordingly, the potential functions satisfy the Lorentz gauge $\nabla \cdot \mathbf{A}(\mathbf{r}, t) + (1/c^2) \partial \psi / \partial t = 0$. Let $f(\mathbf{r}, t)$ be an arbitrary, real-valued, scalar function defined over the relevant region of space-time, then define new potential functions $\mathbf{A}'(\mathbf{r}, t) = \mathbf{A}(\mathbf{r}, t) + \nabla f(\mathbf{r}, t)$ and $\psi'(\mathbf{r}, t) = \psi(\mathbf{r}, t) - \partial f(\mathbf{r}, t) / \partial t$.

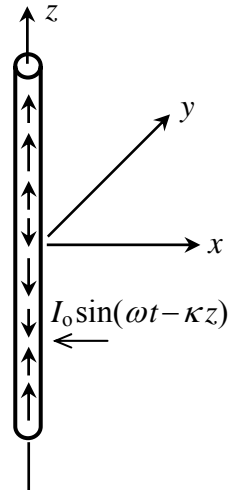
- 1 pt a) Show that the magnetic field $\mathbf{B}(\mathbf{r}, t)$ derived from the new vector potential $\mathbf{A}'(\mathbf{r}, t)$ is the same as that obtained from the old potential $\mathbf{A}(\mathbf{r}, t)$.
- 1 pt b) Show that the electric field $\mathbf{E}(\mathbf{r}, t)$ derived from the new potentials $\mathbf{A}'(\mathbf{r}, t)$ and $\psi'(\mathbf{r}, t)$ is the same as that obtained from the old potentials $\mathbf{A}(\mathbf{r}, t)$ and $\psi(\mathbf{r}, t)$.
- 2 pts c) What restrictions should one impose on $f(\mathbf{r}, t)$ to ensure that the new potentials also satisfy the Lorentz gauge?
- 2 pts d) Under what circumstances will the new potentials satisfy the Coulomb gauge $\nabla \cdot \mathbf{A}'(\mathbf{r}, t) = 0$? In Coulomb gauge, write differential equations that relate $\psi'(\mathbf{r}, t)$ and $\mathbf{A}'(\mathbf{r}, t)$ to $\rho(\mathbf{r}, t)$ and $\mathbf{J}(\mathbf{r}, t)$.

2. An infinitely-long thin wire carries a current $I(z, t) = I_0 \sin(\omega t - \kappa z)$ along the z -axis. For $0 < \kappa < k_0$, where $k_0 = \omega/c$, the scalar and vector potentials in cylindrical coordinates are

$$\psi(\rho, z, t) = -\frac{\kappa I_0}{4\epsilon_0 \omega} \left[Y_0(\sqrt{k_0^2 - \kappa^2} \rho) \sin(\omega t - \kappa z) + J_0(\sqrt{k_0^2 - \kappa^2} \rho) \cos(\omega t - \kappa z) \right],$$

$$\mathbf{A}(\rho, z, t) = -\frac{\mu_0 I_0 \hat{z}}{4} \left[Y_0(\sqrt{k_0^2 - \kappa^2} \rho) \sin(\omega t - \kappa z) + J_0(\sqrt{k_0^2 - \kappa^2} \rho) \cos(\omega t - \kappa z) \right].$$

- 1 pt a) Use the continuity equation to find the charge distribution along the length of the wire.
- 1 pt b) Show that $\mathbf{A}(\rho, z, t)$ and $\psi(\rho, z, t)$ satisfy the Lorentz gauge.
- 2 pts c) Find the \mathbf{E} - and \mathbf{H} -fields in the region outside the wire.
- 2 pts d) Using approximate expressions for the Bessel functions, find the Poynting vector $\mathbf{S}(\rho, z, t)$ in the far-field region.



3. In empty free space where $\rho_{\text{free}} = 0$, $\mathbf{J}_{\text{free}} = 0$, $\mathbf{P} = 0$, $\mathbf{M} = 0$, a plane-wave has potential functions $\mathbf{A}(\mathbf{r}, t) = \mathbf{A}_0 \exp[i(k_0 \boldsymbol{\sigma} \cdot \mathbf{r} - \omega t)]$ and $\psi(\mathbf{r}, t) = \psi_0 \exp[i(k_0 \boldsymbol{\sigma} \cdot \mathbf{r} - \omega t)]$. Here $k_0 = \omega/c$, $\boldsymbol{\sigma} \cdot \boldsymbol{\sigma} = 1$, and the potentials satisfy the Lorentz gauge $\nabla \cdot \mathbf{A}(\mathbf{r}, t) + (1/c^2) \partial \psi / \partial t = 0$.

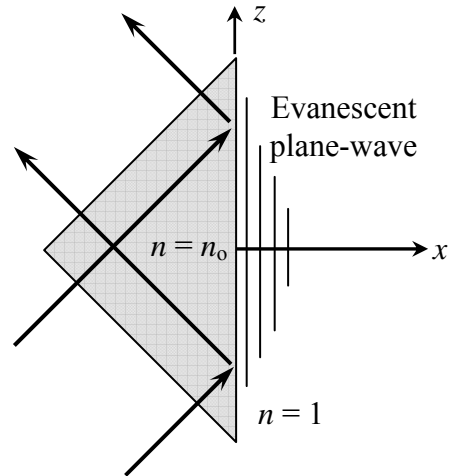
- 2 pts a) What constraint does the Lorentz gauge impose on \mathbf{A}_0 , ψ_0 , and $\boldsymbol{\sigma}$?
- 2 pts b) Derive expressions for the plane-wave's \mathbf{E} -field and \mathbf{H} -field in terms of the defining parameters of the potential functions.
- 2 pts c) Show that the \mathbf{E} - and \mathbf{H} -fields obtained in part (b) satisfy all four Maxwell's equations.

4. A homogeneous plane-wave undergoes total internal reflection within a glass prism of refractive index $n=n_0$, as shown. The resulting evanescent plane-wave in the free-space region beyond the prism (i.e., $x \geq 0$) has the following E - and H -fields:

$$\mathbf{E}(\mathbf{r}, t) = \mathbf{E}_0 \exp[i(k_0 \boldsymbol{\sigma} \cdot \mathbf{r} - \omega t)],$$

$$\mathbf{H}(\mathbf{r}, t) = \mathbf{H}_0 \exp[i(k_0 \boldsymbol{\sigma} \cdot \mathbf{r} - \omega t)].$$

Let $\boldsymbol{\sigma} = i\sigma_x \hat{\mathbf{x}} + \sigma_z \hat{\mathbf{z}}$, while $\mathbf{E}_0 = E_{x0} \hat{\mathbf{x}} + E_{y0} \hat{\mathbf{y}} + E_{z0} \hat{\mathbf{z}}$ and $\mathbf{H}_0 = H_{x0} \hat{\mathbf{x}} + H_{y0} \hat{\mathbf{y}} + H_{z0} \hat{\mathbf{z}}$. In general, σ_x and σ_z are real-valued, while the components of \mathbf{E}_0 and \mathbf{H}_0 are complex.



- 1 pt a) What is the relationship between σ_x and σ_z ?
- 1 pt b) What does Maxwell's first equation say about the relation between E_{x0} and E_{z0} ?
- 1 pt c) What does Maxwell's fourth equation say about the relation between H_{x0} and H_{z0} ?
- 2 pts d) What is the relation between (E_{x0}, E_{y0}, E_{z0}) and (H_{x0}, H_{y0}, H_{z0}) based on Maxwell's third equation?
- 1 pt e) If $E_{z0} = 0$, which components of \mathbf{E}_0 and \mathbf{H}_0 will vanish as well?
- 1 pt f) If $H_{z0} = 0$, which components of \mathbf{E}_0 and \mathbf{H}_0 will vanish as well?
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