

Problem 1) a) The standard Maxwell's equations are

$$\nabla \cdot \mathbf{D}(\mathbf{r}, t) = \rho_{\text{free}}(\mathbf{r}, t), \quad (1a)$$

$$\nabla \times \mathbf{H}(\mathbf{r}, t) = \mathbf{J}_{\text{free}}(\mathbf{r}, t) + \partial \mathbf{D}(\mathbf{r}, t) / \partial t, \quad (1b)$$

$$\nabla \times \mathbf{E}(\mathbf{r}, t) = -\partial \mathbf{B}(\mathbf{r}, t) / \partial t, \quad (1c)$$

$$\nabla \cdot \mathbf{B}(\mathbf{r}, t) = 0. \quad (1d)$$

To eliminate \mathbf{E} and \mathbf{B} , one need only modify the third and fourth equations, as follows:

$$\begin{aligned} \text{Eq. (1c):} \quad \varepsilon_0 \nabla \times \mathbf{E}(\mathbf{r}, t) + \nabla \times \mathbf{P}(\mathbf{r}, t) &= -\varepsilon_0 \partial \mathbf{M}(\mathbf{r}, t) / \partial t - \varepsilon_0 \mu_0 \partial \mathbf{H}(\mathbf{r}, t) / \partial t + \nabla \times \mathbf{P}(\mathbf{r}, t) \\ \Rightarrow \nabla \times \mathbf{D}(\mathbf{r}, t) &= -\varepsilon_0 [\partial \mathbf{M}(\mathbf{r}, t) / \partial t - \varepsilon_0^{-1} \nabla \times \mathbf{P}(\mathbf{r}, t)] - \varepsilon_0 \mu_0 \partial \mathbf{H}(\mathbf{r}, t) / \partial t. \end{aligned} \quad (1c')$$

$$\text{Eq. (1d):} \quad \mu_0 \nabla \cdot \mathbf{H}(\mathbf{r}, t) = -\nabla \cdot \mathbf{M}(\mathbf{r}, t). \quad (1d')$$

b) Transforming the modified equations to the Fourier domain yields,

$$i\mathbf{k} \cdot \mathbf{D}(\mathbf{k}, \omega) = \rho_{\text{free}}(\mathbf{k}, \omega), \quad (2a)$$

$$i\mathbf{k} \times \mathbf{H}(\mathbf{k}, \omega) = \mathbf{J}_{\text{free}}(\mathbf{k}, \omega) - i\omega \mathbf{D}(\mathbf{k}, \omega), \quad (2b)$$

$$\mathbf{k} \times \mathbf{D}(\mathbf{k}, \omega) = \varepsilon_0 \omega \mathbf{M}(\mathbf{k}, \omega) + \mathbf{k} \times \mathbf{P}(\mathbf{k}, \omega) + (\omega/c^2) \mathbf{H}(\mathbf{k}, \omega), \quad (2c)$$

$$\mu_0 \mathbf{k} \cdot \mathbf{H}(\mathbf{k}, \omega) = -\mathbf{k} \cdot \mathbf{M}(\mathbf{k}, \omega). \quad (2d)$$

c) Cross-multiplying “ $-i\mathbf{k}$ ” into Eq.(2b), one arrives at

$$\mathbf{k} \times [\mathbf{k} \times \mathbf{H}(\mathbf{k}, \omega)] = -i\mathbf{k} \times \mathbf{J}_{\text{free}}(\mathbf{k}, \omega) - \omega \mathbf{k} \times \mathbf{D}(\mathbf{k}, \omega). \quad (3)$$

Using the vector identity $\mathbf{k} \times (\mathbf{k} \times \mathbf{H}) = (\mathbf{k} \cdot \mathbf{H})\mathbf{k} - k^2 \mathbf{H}$ in the preceding equation, one obtains

$$[\mathbf{k} \cdot \mathbf{H}(\mathbf{k}, \omega)]\mathbf{k} - k^2 \mathbf{H}(\mathbf{k}, \omega) = -i\mathbf{k} \times \mathbf{J}_{\text{free}}(\mathbf{k}, \omega) - \omega \mathbf{k} \times \mathbf{D}(\mathbf{k}, \omega). \quad (4)$$

Substitution from Eqs.(2c) and (2d) into Eq.(4) then yields

$$\begin{aligned} -\mu_0^{-1} [\mathbf{k} \cdot \mathbf{M}(\mathbf{k}, \omega)]\mathbf{k} - k^2 \mathbf{H}(\mathbf{k}, \omega) &= -i\mathbf{k} \times \mathbf{J}_{\text{free}}(\mathbf{k}, \omega) - \varepsilon_0 \omega^2 \mathbf{M}(\mathbf{k}, \omega) - \omega \mathbf{k} \times \mathbf{P}(\mathbf{k}, \omega) - (\omega^2/c^2) \mathbf{H}(\mathbf{k}, \omega) \\ \Rightarrow \mathbf{H}(\mathbf{k}, \omega) &= \{ i\mathbf{k} \times \mathbf{J}_{\text{free}}(\mathbf{k}, \omega) + \omega \mathbf{k} \times \mathbf{P}(\mathbf{k}, \omega) + \varepsilon_0 \omega^2 \mathbf{M}(\mathbf{k}, \omega) - \mu_0^{-1} [\mathbf{k} \cdot \mathbf{M}(\mathbf{k}, \omega)]\mathbf{k} \} / (k^2 - \omega^2/c^2). \end{aligned} \quad (5)$$

To determine $\mathbf{D}(\mathbf{k}, \omega)$, proceed along similar lines, namely, cross-multiply \mathbf{k} into Eq.(2c), use the vector identity $\mathbf{k} \times (\mathbf{k} \times \mathbf{D}) = (\mathbf{k} \cdot \mathbf{D})\mathbf{k} - k^2 \mathbf{D}$, then substitute from Eqs.(2a) and (2b) into the resulting equation to obtain

$$\begin{aligned} \mathbf{k} \times [\mathbf{k} \times \mathbf{D}(\mathbf{k}, \omega)] &= \varepsilon_0 \omega \mathbf{k} \times \mathbf{M}(\mathbf{k}, \omega) + \mathbf{k} \times [\mathbf{k} \times \mathbf{P}(\mathbf{k}, \omega)] + (\omega/c^2) \mathbf{k} \times \mathbf{H}(\mathbf{k}, \omega) \\ \Rightarrow [\mathbf{k} \cdot \mathbf{D}(\mathbf{k}, \omega)]\mathbf{k} - k^2 \mathbf{D}(\mathbf{k}, \omega) &= \varepsilon_0 \omega \mathbf{k} \times \mathbf{M}(\mathbf{k}, \omega) + [\mathbf{k} \cdot \mathbf{P}(\mathbf{k}, \omega)]\mathbf{k} - k^2 \mathbf{P}(\mathbf{k}, \omega) \\ &\quad - i(\omega/c^2) \mathbf{J}_{\text{free}}(\mathbf{k}, \omega) - (\omega^2/c^2) \mathbf{D}(\mathbf{k}, \omega) \\ \Rightarrow \mathbf{D}(\mathbf{k}, \omega) &= \{ -i\rho_{\text{free}}(\mathbf{k}, \omega)\mathbf{k} + i(\omega/c^2) \mathbf{J}_{\text{free}}(\mathbf{k}, \omega) + k^2 \mathbf{P}(\mathbf{k}, \omega) - [\mathbf{k} \cdot \mathbf{P}(\mathbf{k}, \omega)]\mathbf{k} \\ &\quad - \varepsilon_0 \omega \mathbf{k} \times \mathbf{M}(\mathbf{k}, \omega) \} / (k^2 - \omega^2/c^2). \end{aligned} \quad (6)$$

Problem 2)

$$a) \quad \rho(\mathbf{r}, t) = q\delta(x-Vt)\delta(y)\delta(z); \quad \mathbf{J}(\mathbf{r}, t) = qV\delta(x-Vt)\delta(y)\delta(z)\hat{\mathbf{x}}.$$

$$b) \quad \rho(\mathbf{k}, \omega) = \iiint_{-\infty}^{\infty} \rho(\mathbf{r}, t) \exp[-i(\mathbf{k}\cdot\mathbf{r} - \omega t)] d\mathbf{r} dt = q \iint_{-\infty}^{\infty} \delta(x-Vt) \exp[-i(k_x x - \omega t)] dx dt \\ = q \int_{-\infty}^{\infty} \exp[i(\omega - V k_x)t] dt = 2\pi q \delta(\omega - V k_x).$$

$$\mathbf{J}(\mathbf{k}, \omega) = \iiint_{-\infty}^{\infty} \mathbf{J}(\mathbf{r}, t) \exp[-i(\mathbf{k}\cdot\mathbf{r} - \omega t)] d\mathbf{r} dt = qV\hat{\mathbf{x}} \iint_{-\infty}^{\infty} \delta(x-Vt) \exp[-i(k_x x - \omega t)] dx dt \\ = qV\hat{\mathbf{x}} \int_{-\infty}^{\infty} \exp[i(\omega - V k_x)t] dt = 2\pi qV\delta(\omega - V k_x)\hat{\mathbf{x}}.$$

The scalar and vector potentials are thus given by

$$\psi(\mathbf{k}, \omega) = \varepsilon_0^{-1} \rho(\mathbf{k}, \omega) / (k^2 - \omega^2/c^2) = (2\pi q / \varepsilon_0) \delta(\omega - V k_x) / (k^2 - \omega^2/c^2); \\ \mathbf{A}(\mathbf{k}, \omega) = \mu_0 \mathbf{J}(\mathbf{k}, \omega) / (k^2 - \omega^2/c^2) = (2\pi \mu_0 q V \hat{\mathbf{x}}) \delta(\omega - V k_x) / (k^2 - \omega^2/c^2).$$

c) Inverse Fourier transforming the scalar potential, we find

$$\psi(\mathbf{r}, t) = (2\pi)^{-4} \iiint_{-\infty}^{\infty} \psi(\mathbf{k}, \omega) \exp[i(\mathbf{k}\cdot\mathbf{r} - \omega t)] d\mathbf{k} d\omega \\ = (2\pi)^{-3} (q/\varepsilon_0) \iiint_{-\infty}^{\infty} (k^2 - \omega^2/c^2)^{-1} \delta(\omega - V k_x) \exp[i(\mathbf{k}\cdot\mathbf{r} - \omega t)] d\mathbf{k} d\omega \\ = (2\pi)^{-3} (q/\varepsilon_0) \iiint_{-\infty}^{\infty} [(1 - V^2/c^2)k_x^2 + k_y^2 + k_z^2]^{-1} \exp\{i[k_x(x - Vt) + k_y y + k_z z]\} dk_x dk_y dk_z$$

Defining the parameter $\gamma = 1/\sqrt{1 - (V/c)^2}$, then changing the variable from k_x to k_x/γ yields

$$\psi(\mathbf{r}, t) = (2\pi)^{-3} (\gamma q / \varepsilon_0) \iiint_{-\infty}^{\infty} (k_x^2 + k_y^2 + k_z^2)^{-1} \exp\{i[k_x \gamma (x - Vt) + k_y y + k_z z]\} dk_x dk_y dk_z \\ = (2\pi)^{-3} (\gamma q / \varepsilon_0) \iiint_{-\infty}^{\infty} k^{-2} \exp\{i\mathbf{k} \cdot [\gamma(x - Vt)\hat{\mathbf{x}} + y\hat{\mathbf{y}} + z\hat{\mathbf{z}}]\} d\mathbf{k} \\ = (2\pi)^{-2} (\gamma q / \varepsilon_0) \int_0^{\infty} dk \int_0^{\pi} \sin\theta \exp[ik\sqrt{\gamma^2(x - Vt)^2 + y^2 + z^2} \cos\theta] d\theta \\ = (\gamma q / 2\pi^2 \varepsilon_0) \int_0^{\infty} \left\{ \frac{\sin[k\sqrt{\gamma^2(x - Vt)^2 + y^2 + z^2}]}{[k\sqrt{\gamma^2(x - Vt)^2 + y^2 + z^2}]} \right\} dk \\ = \gamma q / [4\pi \varepsilon_0 \sqrt{\gamma^2(x - Vt)^2 + y^2 + z^2}]$$

Similarly, the inverse Fourier transform of $\mathbf{A}(\mathbf{k}, \omega)$ is found to be

$$\mathbf{A}(\mathbf{r}, t) = \mu_0 \gamma q V \hat{\mathbf{x}} / [4\pi \sqrt{\gamma^2(x - Vt)^2 + y^2 + z^2}].$$

d) The fields are found using $\mathbf{E}(\mathbf{r}, t) = -\nabla\psi(\mathbf{r}, t) - \partial\mathbf{A}(\mathbf{r}, t)/\partial t$ and $\mathbf{B}(\mathbf{r}, t) = \nabla \times \mathbf{A}(\mathbf{r}, t)$, as follows:

$$\mathbf{E}(\mathbf{r}, t) = (\gamma q / 4\pi \varepsilon_0) [(x - Vt)\hat{\mathbf{x}} + y\hat{\mathbf{y}} + z\hat{\mathbf{z}}] / [\gamma^2(x - Vt)^2 + y^2 + z^2]^{3/2}; \\ \mathbf{B}(\mathbf{r}, t) = (\mu_0 q \gamma V / 4\pi) (-z\hat{\mathbf{y}} + y\hat{\mathbf{z}}) / [\gamma^2(x - Vt)^2 + y^2 + z^2]^{3/2}.$$

Problem 3)

a) $\mathbf{P}(\mathbf{r}, t) = p_0 \delta(x) \delta(y) \delta(z) [\cos(\omega_0 t) \hat{\mathbf{x}} + \sin(\omega_0 t) \hat{\mathbf{y}}]$.

b) $\rho_b^{(e)}(\mathbf{r}, t) = -\nabla \cdot \mathbf{P}(\mathbf{r}, t) = -\partial P_x / \partial x - \partial P_y / \partial y = -p_0 [\delta'(x) \delta(y) \cos(\omega_0 t) + \delta(x) \delta'(y) \sin(\omega_0 t)] \delta(z)$.

$$\mathbf{J}_b^{(e)}(\mathbf{r}, t) = \partial \mathbf{P}(\mathbf{r}, t) / \partial t = -p_0 \omega_0 \delta(x) \delta(y) \delta(z) [\sin(\omega_0 t) \hat{\mathbf{x}} - \cos(\omega_0 t) \hat{\mathbf{y}}].$$

c) Continuity equation:

$$\begin{aligned} \nabla \cdot \mathbf{J}_b^{(e)} + \partial \rho_b^{(e)} / \partial t &= -p_0 \omega_0 [\delta'(x) \delta(y) \delta(z) \sin(\omega_0 t) - \delta(x) \delta'(y) \delta(z) \cos(\omega_0 t)] \\ &\quad - p_0 \omega_0 [-\delta'(x) \delta(y) \sin(\omega_0 t) + \delta(x) \delta'(y) \cos(\omega_0 t)] \delta(z) = 0. \end{aligned}$$

d) For the electric point-dipole $p_0 \cos(\omega_0 t) \hat{\mathbf{z}}$ aligned with the z -axis, the potentials are given in a spherical coordinate system. In the present problem, however, we have two oscillating dipoles, one aligned with the x -axis, having magnitude $p_0 \cos(\omega_0 t)$, the other aligned with the y -axis and having magnitude $p_0 \sin(\omega_0 t)$. Retaining the same spherical coordinate system in which θ is measured from the z -axis, we recognize that, for the first dipole, $\cos\theta$ in the expression of the scalar potential must be replaced with $\sin\theta \cos\phi$, while for the second dipole it must be replaced with $\sin\theta \sin\phi$. Also, for the second dipole, the origin of time t must be shifted by one quarter of one period such that $\cos(\omega_0 t)$ is turned into $\sin(\omega_0 t)$, in which case $\sin(\omega_0 t)$ appearing in the expressions of $\psi(\mathbf{r}, t)$ and $\mathbf{A}(\mathbf{r}, t)$ must undergo a corresponding shift to become $-\cos(\omega_0 t)$. Subsequently we add the respective potentials of the two dipoles to obtain

$$\mathbf{A}(\mathbf{r}, t) = -(\mu_0 p_0 \omega_0 / 4\pi r) \{ \sin[\omega_0(t-r/c)] \hat{\mathbf{x}} - \cos[\omega_0(t-r/c)] \hat{\mathbf{y}} \};$$

$$\begin{aligned} \psi(\mathbf{r}, t) &= (p_0 \sin\theta \cos\phi / 4\pi \epsilon_0 r^2) \{ \cos[\omega_0(t-r/c)] - (\omega_0 r/c) \sin[\omega_0(t-r/c)] \} \\ &\quad + (p_0 \sin\theta \sin\phi / 4\pi \epsilon_0 r^2) \{ \sin[\omega_0(t-r/c)] + (\omega_0 r/c) \cos[\omega_0(t-r/c)] \} \\ &= (p_0 \sin\theta / 4\pi \epsilon_0 r^2) \{ \cos[\omega_0(t-r/c) - \phi] - (\omega_0 r/c) \sin[\omega_0(t-r/c) - \phi] \}. \end{aligned}$$
