

Problem 7.80 a) In the transparent dielectric medium, the k -vector is $\pm(n_0\omega/c)\hat{\mathbf{z}}$, and the \mathbf{H} field magnitude is n_0E_0/Z_0 . Consequently, the incident and reflected fields are given by

$$\mathbf{E}^{(i)}(\mathbf{r}, t) = E_0 \exp[-i(\omega/c)(n_0z + ct)]\hat{\mathbf{x}}, \quad (1a)$$

$$\mathbf{H}^{(i)}(\mathbf{r}, t) = -(n_0E_0/Z_0)\exp[-i(\omega/c)(n_0z + ct)]\hat{\mathbf{y}}. \quad (1b)$$

$$\mathbf{E}^{(r)}(\mathbf{r}, t) = \rho E_0 \exp[i(\omega/c)(n_0z - ct)]\hat{\mathbf{x}}, \quad (2a)$$

$$\mathbf{H}^{(r)}(\mathbf{r}, t) = (n_0\rho E_0/Z_0)\exp[i(\omega/c)(n_0z - ct)]\hat{\mathbf{y}}. \quad (2b)$$

In the absorptive substrate, the k -vector is complex-valued, that is, $\mathbf{k} = -(n + i\kappa)(\omega/c)\hat{\mathbf{z}}$, and the \mathbf{H} field amplitude is $(n + i\kappa)/Z_0$ times that of the \mathbf{E} field. Therefore,

$$\mathbf{E}^{(t)}(\mathbf{r}, t) = \tau E_0 \exp\{-i(\omega/c)[(n + i\kappa)z + ct]\}\hat{\mathbf{x}}, \quad (3a)$$

$$\mathbf{H}^{(t)}(\mathbf{r}, t) = -(n + i\kappa)(\tau E_0/Z_0)\exp\{-i(\omega/c)[(n + i\kappa)z + ct]\}\hat{\mathbf{y}}. \quad (3b)$$

b) Matching the boundary conditions means enforcing the continuity of the \mathbf{E}_{\parallel} and \mathbf{H}_{\parallel} at $z = 0$. We will have

$$E_x^{(i)}(z = 0^+) + E_x^{(r)}(z = 0^+) = E_x^{(t)}(z = 0^-) \quad \rightarrow \quad E_0 + \rho E_0 = \tau E_0, \quad (4a)$$

$$H_y^{(i)}(z = 0^+) + H_y^{(r)}(z = 0^+) = H_y^{(t)}(z = 0^-) \quad \rightarrow \quad n_0 E_0 - n_0 \rho E_0 = \tau (n + i\kappa) E_0. \quad (4b)$$

Solving the above equations for ρ and τ , we find

$$\rho = \frac{n_0 - (n + i\kappa)}{n_0 + (n + i\kappa)}, \quad (5a)$$

$$\tau = 1 + \rho = \frac{2n_0}{n_0 + (n + i\kappa)}. \quad (5b)$$

c) The time-averaged rate of EM energy flow per unit area per unit time is $\langle \mathbf{S}(\mathbf{r}, t) \rangle = \frac{1}{2} \text{Re}(\mathbf{E} \times \mathbf{H}^*)$. The corresponding entities for the incident, reflected, and transmitted beams are

$$\begin{aligned} \langle \mathbf{S}^{(i)}(\mathbf{r}, t) \rangle &= -\frac{1}{2} \text{Re}\{E_0 \exp[-i(\omega/c)(n_0z + ct)] (n_0 E_0^*/Z_0) \exp[i(\omega/c)(n_0z + ct)]\}\hat{\mathbf{z}} \\ &= -\frac{1}{2} n_0 Z_0^{-1} |E_0|^2 \hat{\mathbf{z}}. \end{aligned} \quad (6)$$

$$\langle \mathbf{S}^{(r)}(\mathbf{r}, t) \rangle = \frac{1}{2} n_0 Z_0^{-1} |\rho E_0|^2 \hat{\mathbf{z}}. \quad (7)$$

$$\begin{aligned} \langle \mathbf{S}^{(t)}(\mathbf{r}, t) \rangle &= -\frac{1}{2} \text{Re}\{\tau E_0 \exp(-i(\omega/c)[(n + i\kappa)z + ct]) \\ &\quad \times (n - i\kappa) (\tau^* E_0^*/Z_0) \exp(i(\omega/c)[(n - i\kappa)z + ct])\}\hat{\mathbf{z}} \\ &= -\frac{1}{2} n Z_0^{-1} |\tau E_0|^2 \exp(2\kappa\omega z/c) \hat{\mathbf{z}}. \end{aligned} \quad (8)$$

d) The energy balance equation at $z = 0$ may thus be written as follows:

$$n_0 - n_0 |\rho|^2 = n |\tau|^2 \quad \rightarrow \quad |\rho|^2 + \left(\frac{n}{n_0}\right) |\tau|^2 = 1. \quad (9)$$

To confirm the above energy balance equation, note that

$$|\rho|^2 = \rho\rho^* = \frac{(n_0-n-i\kappa)(n_0-n+i\kappa)}{(n_0+n+i\kappa)(n_0+n-i\kappa)} = \frac{(n_0-n)^2+\kappa^2}{(n_0+n)^2+\kappa^2}, \quad (10a)$$

$$|\tau|^2 = \tau\tau^* = \frac{4n_0^2}{(n_0+n)^2+\kappa^2}. \quad (10b)$$

Consequently,

$$|\rho|^2 + \left(\frac{n}{n_0}\right) |\tau|^2 = \frac{(n_0-n)^2+\kappa^2}{(n_0+n)^2+\kappa^2} + \frac{4nn_0}{(n_0+n)^2+\kappa^2} = \frac{(n_0+n)^2+\kappa^2}{(n_0+n)^2+\kappa^2} = 1. \quad (11)$$

The energy absorbed in the substrate is thus seen to be precisely equal to the difference between the incident and reflected energies.
