

Problem 7.70)

a) At the interface of the perfect conductor with the dielectric layer, the tangential E -field must vanish. Therefore, $E_x(x, y, z = 0, t) = E_1 \sin(\varphi_1) \cos(\omega_0 t) = 0$, which leads to $\varphi_1 = 0$ or π . In what follows, we shall set $\varphi_1 = 0$.

$$\begin{aligned} \text{b) } \nabla \times \mathbf{E}(\mathbf{r}, t) &= -\frac{\partial \mathbf{B}(\mathbf{r}, t)}{\partial t} \rightarrow \left(\frac{\partial E_x}{\partial z}\right) \hat{\mathbf{y}} = -\mu_0 \mu(\omega_0) \frac{\partial \mathbf{H}}{\partial t} \\ &\rightarrow E_1 k_1 \cos(k_1 z) \cos(\omega_0 t) = -\mu_0 \frac{\partial H_y}{\partial t} \rightarrow \mathbf{H}(\mathbf{r}, t) = -\left(\frac{E_1 k_1}{\mu_0 \omega_0}\right) \hat{\mathbf{y}} \cos(k_1 z) \sin(\omega_0 t). \end{aligned}$$

$$\begin{aligned} \text{c) } \nabla \times \mathbf{H}(\mathbf{r}, t) &= \mathbf{J}_{\text{free}}(\mathbf{r}, t) + \frac{\partial \mathbf{D}(\mathbf{r}, t)}{\partial t} \rightarrow -\left(\frac{\partial H_y}{\partial z}\right) \hat{\mathbf{x}} = \varepsilon_0 \varepsilon(\omega_0) \frac{\partial \mathbf{E}(\mathbf{r}, t)}{\partial t} \\ &\rightarrow -\left(\frac{E_1 k_1^2}{\mu_0 \omega_0}\right) \sin(k_1 z) \sin(\omega_0 t) = -\varepsilon_0 \varepsilon(\omega_0) \omega_0 E_1 \sin(k_1 z) \sin(\omega_0 t) \\ &\rightarrow k_1^2 = \mu_0 \varepsilon_0 \varepsilon(\omega_0) \omega_0^2 \rightarrow k_1 = (\omega_0/c) n(\omega_0). \end{aligned}$$

d) Considering that $\rho_{\text{free}} = 0$ and $\mathbf{D} = \varepsilon_0 \varepsilon \mathbf{E}$, the first Maxwell equation becomes $\nabla \cdot \mathbf{E} = 0$. We thus have $\nabla \cdot \mathbf{E} = \partial E_x / \partial x = 0$. Similarly, since $\mathbf{B} = \mu_0 \mu \mathbf{H}$, Maxwell's fourth equation becomes $\nabla \cdot \mathbf{H} = 0$, which is automatically satisfied given that $\nabla \cdot \mathbf{H} = \partial H_y / \partial y = 0$.

$$\begin{aligned} \text{e) } \nabla \times \mathbf{E}(\mathbf{r}, t) &= -\frac{\partial \mathbf{B}(\mathbf{r}, t)}{\partial t} \rightarrow \left(\frac{\partial E_x}{\partial z}\right) \hat{\mathbf{y}} = -\mu_0 \frac{\partial \mathbf{H}}{\partial t} \\ &\rightarrow E_0 k_0 \cos(k_0 z + \varphi_0) \cos(\omega_0 t) = -\mu_0 \frac{\partial H_y}{\partial t} \\ &\rightarrow \mathbf{H}(\mathbf{r}, t) = -\left(\frac{E_0 k_0}{\mu_0 \omega_0}\right) \hat{\mathbf{y}} \cos(k_0 z + \varphi_0) \sin(\omega_0 t). \end{aligned}$$

Substitution for $\mathbf{E}(\mathbf{r}, t)$ and $\mathbf{H}(\mathbf{r}, t)$ in Maxwell's second equation now yields

$$\begin{aligned} \nabla \times \mathbf{H}(\mathbf{r}, t) &= \mathbf{J}_{\text{free}}(\mathbf{r}, t) + \frac{\partial \mathbf{D}(\mathbf{r}, t)}{\partial t} \rightarrow -\left(\frac{\partial H_y}{\partial z}\right) \hat{\mathbf{x}} = \varepsilon_0 \frac{\partial \mathbf{E}(\mathbf{r}, t)}{\partial t} \\ &\rightarrow -\left(\frac{E_0 k_0^2}{\mu_0 \omega_0}\right) \sin(k_0 z + \varphi_0) \sin(\omega_0 t) = -\varepsilon_0 \omega_0 E_0 \sin(k_0 z + \varphi_0) \sin(\omega_0 t) \\ &\rightarrow k_0^2 = \mu_0 \varepsilon_0 \omega_0^2 \rightarrow k_0 = \omega_0/c. \end{aligned}$$

f) At the interface located at $z = d$, both E_x and H_y must be continuous. We thus have

$$\begin{aligned} &\begin{cases} E_1 \sin(k_1 d) \cos(\omega_0 t) = E_0 \sin(k_0 d + \varphi_0) \cos(\omega_0 t) \\ -\left(\frac{E_1 k_1}{\mu_0 \omega_0}\right) \cos(k_1 d) \sin(\omega_0 t) = -\left(\frac{E_0 k_0}{\mu_0 \omega_0}\right) \cos(k_0 d + \varphi_0) \sin(\omega_0 t) \end{cases} \\ &\rightarrow \begin{cases} E_1 \sin(k_1 d) = E_0 \sin(k_0 d + \varphi_0) \\ E_1 k_1 \cos(k_1 d) = E_0 k_0 \cos(k_0 d + \varphi_0) \end{cases} \rightarrow \begin{cases} \tan(k_1 d) = n(\omega_0) \tan(k_0 d + \varphi_0) \\ E_1/E_0 = \sin(k_0 d + \varphi_0) / \sin(k_1 d). \end{cases} \end{aligned}$$

The above equations may now be solved for the values of φ_0 and E_1/E_0 .