

Problem 7.62) a) The expression for the E -field is $\mathbf{E}(\mathbf{r}, t) = \mathbf{E}_0 \exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)]$. The k -vector is, in general, complex-valued, meaning that $\mathbf{k} = \mathbf{k}' + i\mathbf{k}''$. The propagation direction is given by \mathbf{k}' , while \mathbf{k}'' specifies the direction along which the beam is attenuated (whenever $\mathbf{k}'' \neq 0$). The E -field amplitude is given by the complex-valued vector $\mathbf{E}_0 = \mathbf{E}'_0 + i\mathbf{E}''_0$. In the MKSA system of units, \mathbf{E} and \mathbf{E}_0 have units of *volt/meter*, \mathbf{k} has units of m^{-1} , and ω has units of sec^{-1} (or *radians/sec*).

b) If the real-valued vectors \mathbf{E}'_0 and \mathbf{E}''_0 are aligned with each other, or if one of them happens to be zero, then the E -field is said to be linearly-polarized. When both \mathbf{E}'_0 and \mathbf{E}''_0 are non-zero and also have different orientations in space, the E -field is circularly or elliptically polarized. (For circular polarization, \mathbf{E}'_0 and \mathbf{E}''_0 must have equal magnitudes and be perpendicular to each other.)

c) The expression for the H -field is $\mathbf{H}(\mathbf{r}, t) = \mathbf{H}_0 \exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)]$. The H -field amplitude is given by the complex-valued vector $\mathbf{H}_0 = \mathbf{H}'_0 + i\mathbf{H}''_0$. In the MKSA system of units, \mathbf{H} and \mathbf{H}_0 have units of *ampere/meter*.

d) In the absence of $\mathbf{P}(\mathbf{r}, t)$ and $\rho_{\text{free}}(\mathbf{r}, t)$, we will have $\mathbf{D}(\mathbf{r}, t) = \epsilon_0 \mathbf{E}(\mathbf{r}, t)$, and Maxwell's 1st equation reduces to $\nabla \cdot \mathbf{E}(\mathbf{r}, t) = 0$. Substituting the E -field distribution of part (a) in this equation then yields $\mathbf{k} \cdot \mathbf{E}_0 = 0$, which is the constraint imposed on \mathbf{k} and \mathbf{E}_0 by Maxwell's 1st equation.

e) Using the E - and H -field distributions given in (a) and (c), Maxwell's 2nd equation yields: $\mathbf{k} \times \mathbf{H}_0 = -\omega \epsilon_0 \mathbf{E}_0$, which is the only constraint imposed by the 2nd equation on \mathbf{k} , ω , \mathbf{E}_0 , and \mathbf{H}_0 .

f) Using the E - and H -field distributions given in (a) and (c), Maxwell's 3rd equation yields: $\mathbf{k} \times \mathbf{E}_0 = \omega \mu_0 \mathbf{H}_0$, which is the only constraint imposed by the 3rd equation on \mathbf{k} , ω , \mathbf{E}_0 , and \mathbf{H}_0 .

g) In the absence of $\mathbf{M}(\mathbf{r}, t)$ we will have $\mathbf{B}(\mathbf{r}, t) = \mu_0 \mathbf{H}(\mathbf{r}, t)$, and Maxwell's 4th equation reduces to $\nabla \cdot \mathbf{H}(\mathbf{r}, t) = 0$. Substituting the H -field distribution of part (c) in this equation then yields $\mathbf{k} \cdot \mathbf{H}_0 = 0$, which is the sole constraint imposed on \mathbf{k} and \mathbf{H}_0 by Maxwell's 4th equation.

h) In part (f) we found $\mathbf{H}_0 = (\omega \mu_0)^{-1} \mathbf{k} \times \mathbf{E}_0$. Substituting this expression for \mathbf{H}_0 into the constraint obtained in part (e) yields: $\mathbf{k} \times [(\omega \mu_0)^{-1} \mathbf{k} \times \mathbf{E}_0] = -\omega \epsilon_0 \mathbf{E}_0$. Using the vector identity $\mathbf{A} \times (\mathbf{B} \times \mathbf{C}) = (\mathbf{A} \cdot \mathbf{C})\mathbf{B} - (\mathbf{A} \cdot \mathbf{B})\mathbf{C}$ we write the preceding equation as $(\mathbf{k} \cdot \mathbf{E}_0)\mathbf{k} - (\mathbf{k} \cdot \mathbf{k})\mathbf{E}_0 = -\mu_0 \epsilon_0 \omega^2 \mathbf{E}_0$. From part (d) we know that $\mathbf{k} \cdot \mathbf{E}_0 = 0$; therefore, $(\mathbf{k} \cdot \mathbf{k})\mathbf{E}_0 = \mu_0 \epsilon_0 \omega^2 \mathbf{E}_0$. Dropping \mathbf{E}_0 from both sides of this equation and using the fact that $\mu_0 \epsilon_0 = 1/c^2$ now yields $k^2 = (\omega/c)^2$, which is the desired dispersion relation.