

**Problem 7.53** a) We use the dispersion relation to find  $k_z$  in terms of  $k_x$ ,  $\omega$ , and material parameters for each plane-wave. We then proceed to relate the various components of the  $E$ - and  $H$ -fields to each other and to the  $k$ -vector through the use of Maxwell's equations. The dispersion relation is

$$k^2 = k_x^2 + k_y^2 + k_z^2 = (\omega/c)^2 \mu(\omega) \varepsilon(\omega) \rightarrow k_z = \pm \sqrt{(\omega/c)^2 \mu(\omega) \varepsilon(\omega) - k_x^2 - k_y^2}. \quad (1)$$

Considering that  $k_y=0$ , and using the relevant parameters for each of the two media, we find

$$k_z^i = -(\omega/c) \sqrt{\mu_a(\omega) \varepsilon_a(\omega) - (ck_x/\omega)^2}; \quad (2a)$$

$$k_z^r = (\omega/c) \sqrt{\mu_a(\omega) \varepsilon_a(\omega) - (ck_x/\omega)^2}; \quad (2b)$$

$$k_z^t = -(\omega/c) \sqrt{\mu_b(\omega) \varepsilon_b(\omega) - (ck_x/\omega)^2}; \quad (2c)$$

The choice of sign for the square root must be made such that the imaginary part of  $k_z$  is positive for upward-propagating waves, and negative for downward-propagating waves.

For  $p$ -polarized light, Maxwell's 1<sup>st</sup> equation yields

$$\nabla \cdot \mathbf{D} = 0 \rightarrow \mathbf{k} \cdot \mathbf{E}_p = 0 \rightarrow k_x E_{xp} + k_z E_{zp} = 0 \rightarrow \begin{cases} E_{zp}^i = -(k_x/k_z^i) E_{xp}^i \\ E_{zp}^r = -(k_x/k_z^r) E_{xp}^r \\ E_{zp}^t = -(k_x/k_z^t) E_{xp}^t \end{cases} \quad (3)$$

As for the  $H$ -field of the various  $p$ -polarized beams, we use Maxwell's 3<sup>rd</sup> equation to write

$$\begin{aligned} \nabla \times \mathbf{E} = -\partial \mathbf{B} / \partial t &\rightarrow \mathbf{k} \times \mathbf{E} = \mu_0 \mu(\omega) \omega \mathbf{H} \rightarrow k_z E_{xp} - k_x E_{zp} = \mu_0 \mu(\omega) \omega H_{yp} \\ &\rightarrow k_z E_{xp} + (k_x^2/k_z) E_{xp} = \mu_0 \mu(\omega) \omega H_{yp} \rightarrow (k_x^2 + k_z^2) E_{xp} = \mu_0 \mu(\omega) \omega k_z H_{yp} \\ &\rightarrow (\omega/c)^2 \mu(\omega) \varepsilon(\omega) E_{xp} = \mu_0 \mu(\omega) \omega k_z H_{yp} \rightarrow (\omega/c) \varepsilon(\omega) E_{xp} = \mu_0 c k_z H_{yp} \end{aligned}$$

$$\rightarrow H_{yp} = \frac{(\omega/c) \varepsilon(\omega)}{Z_0 k_z} E_{xp} \rightarrow \begin{cases} H_{yp}^i = \frac{(\omega/c) \varepsilon_a(\omega)}{Z_0 k_z^i} E_{xp}^i \\ H_{yp}^r = \frac{(\omega/c) \varepsilon_a(\omega)}{Z_0 k_z^r} E_{xp}^r \\ H_{yp}^t = \frac{(\omega/c) \varepsilon_b(\omega)}{Z_0 k_z^t} E_{xp}^t \end{cases} \quad (4)$$

For the  $s$ -polarized light, we use Maxwell's 4<sup>th</sup> equation to relate  $H_z$  to  $H_x$ , as follows:

$$\mathbf{k} \cdot \mathbf{H} = 0 \rightarrow k_x H_{xs} + k_z H_{zs} = 0 \rightarrow \begin{cases} H_{zs}^i = -(k_x/k_z^i) H_{xs}^i \\ H_{zs}^r = -(k_x/k_z^r) H_{xs}^r \\ H_{zs}^t = -(k_x/k_z^t) H_{xs}^t \end{cases} \quad (5)$$

The  $E$ -field of the  $s$ -polarized beam is readily obtained from Maxwell's 2<sup>nd</sup> equation, that is,

$$\begin{aligned}
\nabla \times \mathbf{H} &= \partial \mathbf{D} / \partial t \rightarrow \mathbf{k} \times \mathbf{H} = -\varepsilon_0 \varepsilon(\omega) \omega \mathbf{E} \rightarrow k_z H_{xs} - k_x H_{zs} = -\varepsilon_0 \varepsilon(\omega) \omega E_{ys} \\
&\rightarrow k_z H_{xs} + (k_x^2 / k_z) H_{xs} = -\varepsilon_0 \varepsilon(\omega) \omega E_{ys} \rightarrow (k_x^2 + k_z^2) H_{xs} = -\varepsilon_0 \varepsilon(\omega) \omega k_z E_{ys} \\
&\rightarrow (\omega / c)^2 \mu(\omega) \varepsilon(\omega) H_{xs} = -\varepsilon_0 \varepsilon(\omega) \omega k_z E_{ys} \rightarrow (\omega / c) \mu(\omega) H_{xs} = -\varepsilon_0 c k_z E_{ys} \\
&\rightarrow E_{ys} = -\frac{(\omega / c) \mu(\omega)}{k_z} Z_0 H_{xs} \rightarrow \begin{cases} E_{ys}^i = -\frac{(\omega / c) \mu_a(\omega)}{k_z^i} Z_0 H_{xs}^i \\ E_{ys}^r = -\frac{(\omega / c) \mu_a(\omega)}{k_z^r} Z_0 H_{xs}^r \\ E_{ys}^t = -\frac{(\omega / c) \mu_b(\omega)}{k_z^t} Z_0 H_{xs}^t \end{cases} \quad (6)
\end{aligned}$$

b) For  $p$ -polarized light, the continuity of  $E_x$  and  $D_z$  at the  $z=0$  interface yields

$$\begin{cases} E_{xp}^i + E_{xp}^r = E_{xp}^t \\ D_{zp}^i + D_{zp}^r = D_{zp}^t \end{cases} \rightarrow \begin{cases} E_{xp}^i + E_{xp}^r = E_{xp}^t \\ \varepsilon_0 \varepsilon_a E_{zp}^i + \varepsilon_0 \varepsilon_a E_{zp}^r = \varepsilon_0 \varepsilon_b E_{zp}^t \end{cases} \quad (7a)$$

Use Eq. (3) in Eq. (7b), then substitute for  $E_{xp}^t$  from Eq. (7a).

$$\begin{aligned}
&\rightarrow (\varepsilon_a k_x / k_z^i) E_{xp}^i + (\varepsilon_a k_x / k_z^r) E_{xp}^r = (\varepsilon_b k_x / k_z^t) (E_{xp}^i + E_{xp}^r) \\
&\rightarrow [(\varepsilon_a / k_z^r) - (\varepsilon_b / k_z^t)] E_{xp}^r = [(\varepsilon_b / k_z^t) - (\varepsilon_a / k_z^i)] E_{xp}^i
\end{aligned} \quad (7b)$$

Use Eqs. (2a, 2b) to set  $k_z^r = -k_z^i$ .

$$\rightarrow \rho_p = \frac{E_{xp}^r}{E_{xp}^i} = \frac{(\varepsilon_b / k_z^t) - (\varepsilon_a / k_z^i)}{(\varepsilon_a / k_z^r) - (\varepsilon_b / k_z^t)} = \frac{\varepsilon_a k_z^t - \varepsilon_b k_z^i}{\varepsilon_a k_z^t + \varepsilon_b k_z^i}. \quad (8)$$

The transmission coefficient  $\tau_p$  is found from Eqs. (7a) and (8), as follows:

$$\tau_p = E_{xp}^t / E_{xp}^i = 1 + (E_{xp}^r / E_{xp}^i) = 1 + \rho_p = \frac{2\varepsilon_a k_z^t}{\varepsilon_a k_z^t + \varepsilon_b k_z^i}. \quad (9)$$

c) For  $s$ -polarized light, the continuity of  $H_x$  and  $B_z$  at the  $z=0$  interface yields

$$\begin{cases} H_{xs}^i + H_{xs}^r = H_{xs}^t \\ B_{zs}^i + B_{zs}^r = B_{zs}^t \end{cases} \rightarrow \begin{cases} H_{xs}^i + H_{xs}^r = H_{xs}^t \\ \mu_0 \mu_a H_{zs}^i + \mu_0 \mu_a H_{zs}^r = \mu_0 \mu_a H_{zs}^t \end{cases} \quad (10a)$$

Use Eq. (5) in Eq. (10b), then substitute for  $H_{xs}^t$  from Eq. (10a).

$$\begin{aligned}
&\rightarrow (\mu_a k_x / k_z^i) H_{xs}^i + (\mu_a k_x / k_z^r) H_{xs}^r = (\mu_b k_x / k_z^t) (H_{xs}^i + H_{xs}^r) \\
&\rightarrow [(\mu_a / k_z^r) - (\mu_b / k_z^t)] H_{xs}^r = [(\mu_b / k_z^t) - (\mu_a / k_z^i)] H_{xs}^i
\end{aligned} \quad (10b)$$

Use Eqs. (2a, 2b) to set  $k_z^r = -k_z^i$ .

$$\rightarrow \frac{H_{xs}^r}{H_{xs}^i} = \frac{(\mu_b / k_z^t) - (\mu_a / k_z^i)}{(\mu_a / k_z^r) - (\mu_b / k_z^t)} = \frac{\mu_a k_z^t - \mu_b k_z^i}{\mu_a k_z^t + \mu_b k_z^i}. \quad (11)$$

The Fresnel reflection coefficient for  $s$ -polarized light is defined as  $\rho_s = E_{y_s}^r / E_{y_s}^i$ . From Eq.(6), it is clear that  $\rho_s = -H_{x_s}^r / H_{x_s}^i$ . Therefore,

$$\rho_s = \frac{\mu_b k_z^i - \mu_a k_z^t}{\mu_b k_z^i + \mu_a k_z^t}. \quad (12)$$

The transmission coefficient for the  $H$ -field is found from Eqs.(10a) and (11), as follows:

$$H_{x_s}^t / H_{x_s}^i = 1 + (H_{x_s}^r / H_{x_s}^i) = \frac{2\mu_a k_z^t}{\mu_a k_z^t + \mu_b k_z^i}. \quad (13)$$

The Fresnel transmission coefficient for  $s$ -polarized light, being defined as  $\tau_s = E_{y_s}^t / E_{y_s}^i$ , may now be found from Eq.(6) as  $\tau_s = (\mu_b k_z^i / \mu_a k_z^t) H_{x_s}^t / H_{x_s}^i$ . Consequently

$$\tau_s = \frac{2\mu_b k_z^i}{\mu_a k_z^t + \mu_b k_z^i}. \quad (14)$$


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