

Problem 42) a)

No. At  $t=0$  we have  $\vec{E}(\vec{r}_0, t) = \vec{E}_{OR}$ , while at  $t=T/4$ , where  $T$  is the period of the oscillation, we'll have  $\vec{E}(\vec{r}_0, t) = \vec{E}_{OI}$ . If  $\vec{E}_{OR}$  and  $\vec{E}_{OI}$  are not aligned, the  $\vec{E}$ -field will have two different orientations at  $t=0$  and  $t=T/4$ . For linear polarization, therefore, it is necessary for  $\vec{E}_{OR}$  and  $\vec{E}_{OI}$  to be aligned.

$$b) \vec{E}(\vec{r}_0, t) = \vec{E}_{OR} \cos \omega t + \vec{E}_{OI} \sin \omega t$$

When  $\vec{E}_{OR}$  and  $\vec{E}_{OI}$  are aligned their common direction will be the direction of  $\vec{E}(\vec{r}_0, t)$ . The length of  $\vec{E}(\vec{r}_0, t)$  can then be written in terms of the lengths of the vectors  $\vec{E}_{OR}$  and  $\vec{E}_{OI}$ , namely,  $E_{OR}$  and  $E_{OI}$ , as follows:

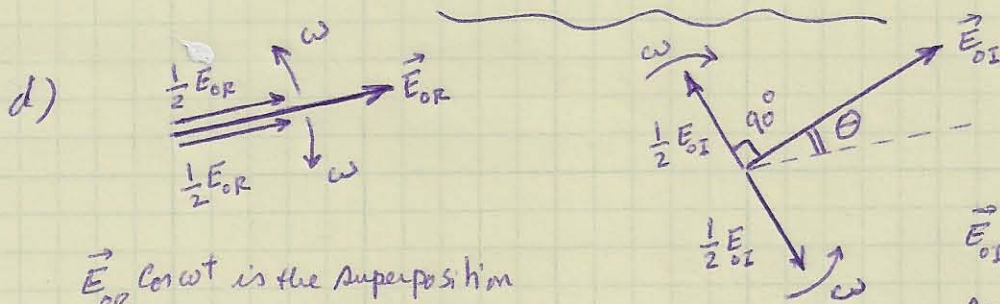
$$E(\vec{r}_0, t) = E_{OR} \cos \omega t + E_{OI} \sin \omega t = \sqrt{E_{OR}^2 + E_{OI}^2} \left( \frac{E_{OR}}{\sqrt{E_{OR}^2 + E_{OI}^2}} \cos \omega t + \frac{E_{OI}}{\sqrt{E_{OR}^2 + E_{OI}^2}} \sin \omega t \right)$$

$\swarrow \cos \phi$   $\searrow \sin \phi$   
 $\searrow \sin \phi$   $\swarrow \cos \phi$

$$= \sqrt{E_{OR}^2 + E_{OI}^2} (\cos \phi \cos \omega t + \sin \phi \sin \omega t) = \sqrt{E_{OR}^2 + E_{OI}^2} \cos(\omega t - \phi)$$

The magnitude of the  $\vec{E}$ -field is, therefore,  $\sqrt{E_{OR}^2 + E_{OI}^2}$ .

c) No. At  $t=0$  we have  $\vec{E}(\vec{r}_0, t) = \vec{E}_{OR}$ , while at  $t=T/4$  we'll have  $\vec{E}(\vec{r}_0, t) = \vec{E}_{OI}$ . A circularly-polarized beam must have the same  $E$ -field magnitude at all times. Therefore, if  $|\vec{E}_{OR}| \neq |\vec{E}_{OI}|$ , the  $E$ -field magnitude at  $t=0$  will differ from its magnitude at  $t=T/4$ , which means that the beam is not circularly polarized.



$\vec{E}_{OR} \cos \omega t$  is the superposition of two circularly-polarized beams

$\vec{E}_{OI} \sin \omega t$  is also the superposition of two circularly-polarized beams.

In general, the two RCP beams combine to produce a single RCP beam. Similarly, the two LCP beams combine to produce a single LCP beam. When  $\theta = \pm 90^\circ$ , however, either the two RCP beams cancel each other out, or the two LCP beams cancel each other out. The result will then be a single circularly polarized beam. When  $\theta \neq 90^\circ$ , both RCP and LCP beams will have non-zero magnitudes, and the resulting beam can't be a pure circularly-polarized beam.

$$e) \vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \Rightarrow ik_0 \vec{\sigma} \times \vec{E}_0 = -(-i\omega) \mu_0 \vec{H}_0 \Rightarrow i\frac{\omega}{c} \vec{\sigma} \times \vec{E}_0 = i\omega \mu_0 \vec{H}_0 \\ \Rightarrow \vec{\sigma} \times \vec{E}_0 = \epsilon_0 \vec{H}_0.$$

For a homogeneous plane-wave  $\vec{\sigma} = \vec{\sigma}_R$ ; therefore,

$$\epsilon_0 \vec{H}_0 = \vec{\sigma}_R \times (\vec{E}_{0R} + i\vec{E}_{0I}) \Rightarrow \begin{cases} \epsilon_0 \vec{H}_{0R} = \vec{\sigma}_R \times \vec{E}_{0R} \\ \epsilon_0 \vec{H}_{0I} = \vec{\sigma}_R \times \vec{E}_{0I} \end{cases}$$

From Maxwell's first equation,  $\vec{\nabla} \cdot \vec{E} = 0$ , we know that  $\vec{\sigma} \cdot \vec{E} = 0$ . Therefore, for a homogeneous plane-wave,  $\vec{\sigma}_R \cdot (\vec{E}_{0R} + i\vec{E}_{0I}) = 0 \\ \Rightarrow \vec{\sigma}_R \cdot \vec{E}_{0R} = 0 \text{ and } \vec{\sigma}_R \cdot \vec{E}_{0I} = 0$ . In other words, the unit-vector  $\vec{\sigma}_R$  is perpendicular to both  $\vec{E}_{0R}$  and  $\vec{E}_{0I}$ . We conclude that  $\vec{\sigma}_R$ ,  $\vec{E}_{0R}$  and  $\vec{H}_{0R}$  are mutually orthogonal, and also  $\vec{\sigma}_R$ ,  $\vec{E}_{0I}$  and  $\vec{H}_{0I}$  are mutually orthogonal. The magnitudes of  $\vec{H}_{0R}$  and  $\vec{H}_{0I}$  are thus given by:  $|\vec{H}_{0R}| = |\vec{E}_{0R}| / \epsilon_0$  and  $|\vec{H}_{0I}| = |\vec{E}_{0I}| / \epsilon_0$ .

$$f) \vec{E}(\vec{r}, t) = \text{Re} \{ \vec{E}_0 \exp[i(k_0 \vec{\sigma}_R \cdot \vec{r} - \omega t)] \} = \vec{E}_{0R} \cos(k_0 \vec{\sigma}_R \cdot \vec{r} - \omega t) - \vec{E}_{0I} \sin(k_0 \vec{\sigma}_R \cdot \vec{r} - \omega t) \\ \vec{H}(\vec{r}, t) = \vec{H}_{0R} \cos(k_0 \vec{\sigma}_R \cdot \vec{r} - \omega t) - \vec{H}_{0I} \sin(k_0 \vec{\sigma}_R \cdot \vec{r} - \omega t) \\ \vec{S}(\vec{r}, t) = \vec{E}(\vec{r}, t) \times \vec{H}(\vec{r}, t) = \vec{E}_{0R} \times \vec{H}_{0R} \cos^2(k_0 \vec{\sigma}_R \cdot \vec{r} - \omega t) + \vec{E}_{0I} \times \vec{H}_{0I} \sin^2(k_0 \vec{\sigma}_R \cdot \vec{r} - \omega t) \\ - (\vec{E}_{0R} \times \vec{H}_{0I} + \vec{E}_{0I} \times \vec{H}_{0R}) \sin(k_0 \vec{\sigma}_R \cdot \vec{r} - \omega t) \cos(k_0 \vec{\sigma}_R \cdot \vec{r} - \omega t)$$

Using the vector identity  $\vec{A} \times (\vec{B} \times \vec{C}) = (\vec{A} \cdot \vec{C})\vec{B} - (\vec{A} \cdot \vec{B})\vec{C}$ , we write;

$$\vec{E}_{OR} \times \vec{H}_{OR} = \frac{1}{Z_0} \vec{E}_{OR} \times (\vec{\sigma}_R \times \vec{E}_{OR}) = \frac{1}{Z_0} (\vec{E}_{OR} \cdot \vec{E}_{OR}) \vec{\sigma}_R - \frac{1}{Z_0} (\vec{E}_{OR} \cdot \vec{\sigma}_R) \vec{E}_{OR}$$

$$\vec{E}_{OI} \times \vec{H}_{OI} = \frac{1}{Z_0} \vec{E}_{OI} \times (\vec{\sigma}_R \times \vec{E}_{OI}) = \frac{1}{Z_0} (\vec{E}_{OI} \cdot \vec{E}_{OI}) \vec{\sigma}_R - \frac{1}{Z_0} (\vec{E}_{OI} \cdot \vec{\sigma}_R) \vec{E}_{OI}$$

$$\vec{E}_{OR} \times \vec{H}_{OI} = \frac{1}{Z_0} \vec{E}_{OR} \times (\vec{\sigma}_R \times \vec{E}_{OI}) = \frac{1}{Z_0} (\vec{E}_{OR} \cdot \vec{E}_{OI}) \vec{\sigma}_R - \frac{1}{Z_0} (\vec{E}_{OR} \cdot \vec{\sigma}_R) \vec{E}_{OI}$$

$$\vec{E}_{OI} \times \vec{H}_{OR} = \frac{1}{Z_0} \vec{E}_{OI} \times (\vec{\sigma}_R \times \vec{E}_{OR}) = \frac{1}{Z_0} (\vec{E}_{OI} \cdot \vec{E}_{OR}) \vec{\sigma}_R - \frac{1}{Z_0} (\vec{E}_{OI} \cdot \vec{\sigma}_R) \vec{E}_{OR}$$

Therefore,

$$\vec{S}(\vec{r}, t) = \frac{\vec{\sigma}_R}{Z_0} \left\{ \vec{E}_{OR} \cdot \vec{E}_{OR} \cos^2(k_0 \vec{\sigma}_R \cdot \vec{r} - \omega t) + \vec{E}_{OI} \cdot \vec{E}_{OI} \sin^2(k_0 \vec{\sigma}_R \cdot \vec{r} - \omega t) - \vec{E}_{OR} \cdot \vec{E}_{OI} \sin[2(k_0 \vec{\sigma}_R \cdot \vec{r} - \omega t)] \right\} \Rightarrow$$

$$\vec{S}(\vec{r}, t) = \frac{\vec{\sigma}_R}{2Z_0} \left\{ (|\vec{E}_{OR}|^2 + |\vec{E}_{OI}|^2) + (|\vec{E}_{OR}|^2 - |\vec{E}_{OI}|^2) \cos(2k_0 \vec{\sigma}_R \cdot \vec{r} - 2\omega t) - 2\vec{E}_{OR} \cdot \vec{E}_{OI} \sin(2k_0 \vec{\sigma}_R \cdot \vec{r} - 2\omega t) \right\}$$

Note that, for circularly polarized beams,  $|\vec{E}_{OR}| = |\vec{E}_{OI}|$  and  $\vec{E}_{OR} \cdot \vec{E}_{OI} = 0$ ;

Therefore,  $\vec{S}(\vec{r}, t) = \langle \vec{S}(\vec{r}, t) \rangle$ , that is, the rate of flow of energy is independent of  $\vec{r}$  and  $t$ .