

Problem 19) Maxwell's equations in free space:

$$\begin{aligned}
 1) \quad \vec{\nabla} \cdot \vec{E} = 0 &\Rightarrow \vec{k} \cdot \vec{E}_0 = 0 \\
 2) \quad \vec{\nabla} \times \vec{H} = \epsilon_0 \partial \vec{E} / \partial t &\Rightarrow \vec{k} \times \vec{H}_0 = -\epsilon_0 \omega \vec{E}_0 \\
 3) \quad \vec{\nabla} \times \vec{E} = -\mu_0 \partial \vec{H} / \partial t &\Rightarrow \vec{k} \times \vec{E}_0 = +\mu_0 \omega \vec{H}_0 \\
 4) \quad \vec{\nabla} \cdot \vec{H} = 0 &\Rightarrow \vec{k} \cdot \vec{H}_0 = 0
 \end{aligned}
 \left. \begin{array}{l} \\ \\ \\ \end{array} \right\} \Rightarrow \vec{k} \times (\vec{k} \times \vec{E}_0) = -\mu_0 \epsilon_0 \omega^2 \vec{E}_0 \Rightarrow$$

$$(\vec{k} \cdot \vec{E}_0) \vec{k} - k^2 \vec{E}_0 = -(\omega/c)^2 \vec{E}_0 \Rightarrow k^2 = (\omega/c)^2$$

= 0 from Eq. (1))

$$a) \quad \vec{k} = k_x \hat{x} + i k_z \hat{z} \Rightarrow k^2 = \vec{k} \cdot \vec{k} = (k_x \hat{x} + i k_z \hat{z}) \cdot (k_x \hat{x} + i k_z \hat{z}) = k_x^2 - k_z^2 = (\omega/c)^2$$

$$b) \quad \text{From Maxwell's 1st equation: } \vec{k} \cdot \vec{E}_0 = 0 \Rightarrow (k_x \hat{x} + i k_z \hat{z}) \cdot (E_{x0} \hat{x} + i E_{z0} \hat{z}) = 0 \Rightarrow$$

$$k_x E_{x0} - k_z E_{z0} = 0 \Rightarrow k_x E_{x0} = k_z E_{z0}$$

$$c) \quad \text{From Maxwell's 3rd equation: } \vec{k} \times \vec{E}_0 = \mu_0 \omega \vec{H}_0 \Rightarrow (k_x \hat{x} + i k_z \hat{z}) \times (E_{x0} \hat{x} + i E_{z0} \hat{z})$$

$$= i \mu_0 \omega H_{y0} \hat{y} \Rightarrow i k_x E_{z0} (\hat{x} \times \hat{z}) + i k_z E_{x0} (\hat{z} \times \hat{x}) = i (k_z E_{x0} - k_x E_{z0}) \hat{y} = i \mu_0 \omega H_{y0} \hat{y}$$

$$\Rightarrow H_{y0} = \frac{k_z E_{x0} - k_x E_{z0}}{\mu_0 \omega} = \frac{k_z - (k_x^2/k_z)}{\mu_0 \omega} E_{x0} = \frac{k_z^2 - k_x^2}{\mu_0 \omega k_z} E_{x0} = -\frac{(\omega/c)^2}{\mu_0 \omega \sqrt{k_x^2 - \omega^2/c^2}} E_{x0}$$

$$\Rightarrow H_{y0} = -\frac{\omega/c}{\mu_0 \omega \sqrt{(ck_x/\omega)^2 - 1}} E_{x0} \Rightarrow H_{y0} = -\frac{E_{x0}}{z_0 \sqrt{(ck_x/\omega)^2 - 1}}$$

In the above derivation we chose the positive sign for $k_z = \sqrt{k_x^2 - (\omega/c)^2}$, the reason being that, the field-amplitudes are proportional to $\exp(i \vec{k} \cdot \vec{r}) = \exp[i(k_x \hat{x} + i k_z \hat{z}) \cdot \vec{r}] = \exp(-k_z z) \exp(i k_x x)$. When $z \rightarrow +\infty$ the fields can't become indefinitely large; therefore, $k_z \geq 0$.

$$\begin{aligned}
 d) \quad \langle \vec{S}(\vec{r}, t) \rangle &= \frac{1}{2} \operatorname{Re} \{ \vec{E}(\vec{r}) \times \vec{H}^*(\vec{r}) \} = \frac{1}{2} \operatorname{Re} \{ (E_{x_0} \hat{x} + i E_{z_0} \hat{z}) e^{i(k_x \hat{x} + k_z \hat{z}) \cdot \vec{r}} \times \\
 &\quad (-i H_{y_0} \hat{y}) e^{-i(k_x \hat{x} - k_z \hat{z}) \cdot \vec{r}} \} = \frac{1}{2} e^{-2k_z z} \operatorname{Re} \{ (E_{x_0} \hat{x} + i E_{z_0} \hat{z}) \times (-i H_{y_0} \hat{y}) \} \\
 &= \frac{1}{2} e^{-2k_z z} \operatorname{Re} \{ -i E_{x_0} H_{y_0} \hat{z} - E_{z_0} H_{y_0} \hat{x} \} = \frac{-1}{2} \exp(-2k_z z) E_{z_0} H_{y_0} \hat{x}
 \end{aligned}$$

$$\Rightarrow \langle \vec{S}(\vec{r}, t) \rangle = \frac{1}{2} \exp(-2k_z z) \frac{E_{x_0} E_{z_0} \hat{x}}{\epsilon_0 \sqrt{(ck_x/\omega)^2 - 1}} = \frac{k_x E_{x_0}^2 \exp(-2k_z z)}{2 \epsilon_0 k_z \sqrt{(ck_x/\omega)^2 - 1}} \hat{x} \Rightarrow$$

$$\langle \vec{S}(\vec{r}, t) \rangle = \frac{(ck_x/\omega) \exp[-2(\omega/c) \sqrt{(ck_x/\omega)^2 - 1} z]}{2 \epsilon_0 [(ck_x/\omega)^2 - 1]} E_{x_0}^2 \hat{x}.$$

The energy flows along the positive x -axis when $k_x > 0$.