## **Opti 501 Solutions Solutions 1/1**

**Problem 12**) Denoting by  $\sigma$  the normalized *k*-vector  $k/k_0 = ck/\omega_0 = \lambda_0 k/2\pi$ , we write

 $\mathbf{E}(\mathbf{r},t) = \mathbf{E}_0 \exp[i(2\pi/\lambda_0)(\boldsymbol{\sigma}\cdot\mathbf{r}-ct)], \quad \mathbf{H}(\mathbf{r},t) = \mathbf{H}_0 \exp[i(2\pi/\lambda_0)(\boldsymbol{\sigma}\cdot\mathbf{r}-ct)].$ 

Maxwell's first and fourth equations then yield

 $\nabla \cdot \mathbf{E} = 0 \rightarrow \sigma \cdot \mathbf{E}_0 = 0$  and  $\nabla \cdot \mathbf{B} = \mu_0 \nabla \cdot \mathbf{H} = 0 \rightarrow \sigma \cdot \mathbf{H}_0 = 0.$ Invoking Maxwell's second and third equations, we now find

 $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \rightarrow i(2\pi/\lambda_0)\vec{\sigma} \times \vec{E}_0 = -(-i2\pi c/\lambda_0)\mu_0\vec{H}_0 \rightarrow \vec{\sigma} \times \vec{E}_0 = Z_0\vec{H}_0$  $\nabla \times H = \partial D/\partial t \rightarrow i(2\pi/\lambda_0)\sigma \times H_0 = (-i2\pi c/\lambda_0)\varepsilon_0 \varepsilon E_0 \rightarrow \sigma \times H_0 = -(\varepsilon/Z_0)E_0$ . Combining the above equations, we will have  $\boldsymbol{0}$ 

 $\sigma \times (\sigma \times E_{o}) = -\varepsilon E_{o} \quad \rightarrow \quad (\sigma \times E_{o})\sigma - (\sigma \cdot \sigma)E_{o} = -\varepsilon E_{o} \quad \rightarrow \quad \sigma \cdot \sigma = \varepsilon = n^{2}.$ a) In free space:  $n = 1$   $\rightarrow$   $\sigma = \hat{z} \rightarrow Z_0 H_i = \hat{z} \times E_i \hat{x} \rightarrow Z_0 H_i = E_i \hat{y}$ . b) Inside the slab:  $\sigma = n\hat{z} \rightarrow Z_0H_t = \sigma \times E_t = n\hat{z} \times E_t \rightarrow Z_0H_t = nE_t\hat{y}$ . c) Rate of energy flow in free space:  $\langle S_i \rangle = \frac{1}{2} E_i \times H_i = (E_i^2 / 2Z_o) \hat{z}$ . Rate of energy flow in the transparent slab:  $\langle S_{\rm t} \rangle = \frac{1}{2} E_{\rm t} \times H_{\rm t} = (n E_{\rm t}^2 / 2 Z_{\rm o}) \hat{z}$ .  $\langle S_i \rangle = \langle S_t \rangle \rightarrow E_i^2/2Z_0 = nE_t^2/2Z_0 \rightarrow E_t = E_i/\sqrt{n} \rightarrow H_t = nE_t/Z_0 = \sqrt{n}E_i/Z_0.$ 

d) E-field energy-density inside dispersionless medium:  $\frac{1}{4\epsilon_0 \epsilon E_t^2} = \frac{1}{4\epsilon_0 n^2 (E_i/\sqrt{n})^2} = \frac{1}{4\epsilon_0 n E_i^2}$ . *H*-field energy-density inside dispersionless medium:  $\frac{1}{4}\mu_0 H_t^2 = \frac{1}{4}\mu_0 (\sqrt{n} E_i/Z_0)^2 = \frac{1}{4}\varepsilon_0 n E_i^2$ .

Thus, the  $E$ - and  $H$ -field energy-densities within the transparent medium of the slab are equal.

Let the pulse duration and cross-sectional area be  $T$  and  $A$ , respectively. In the free-space region, the length of the pulse is  $cT$ , its volume is  $cTA$ , and its energy-density is  $\frac{1}{4}\epsilon_0 E_1^2 +$  $\frac{1}{4}\mu_0 H_1^2 = \frac{1}{4}\varepsilon_0 E_1^2 + \frac{1}{4}\mu_0 (E_1/Z_0)^2 = \frac{1}{2}\varepsilon_0 E_1^2$ . Consequently, the total energy of the light pulse in the free-space region is  $\frac{1}{2} \varepsilon_0 E_i^2 cTA$ .

Inside the glass medium of the dielectric slab, the length of the pulse is  $cT/n$ , its volume is  $cTA/n$ , and its energy-density is  $\frac{1}{4}\epsilon_0 nE_1^2 + \frac{1}{4}\epsilon_0 nE_1^2 = \frac{1}{2}\epsilon_0 nE_1^2$ . Therefore, the total energy of the pulse propagating within the slab is also  $\frac{1}{2} \varepsilon_0 E_i^2 c T A$ . The pulse energy is thus seen to be preserved.