Problem 2) As in Problem 1, the volume of interest is $\Delta d\Delta s \cos\theta$, of which one half is to the left of Δs and one half to the right. Here Δd is the change of d(t) over a time interval Δt , that is, $\Delta d = d(t + \Delta t) - d(t)$. If a dipole is located within this volume, one of its charges (either +q or -q) will cross the surface Δs during the time interval Δt . Irrespective of which charge crosses, the current flowing from left to right will be $q/\Delta t$. The current crossing Δs due to all dipoles that are oriented at an angle θ to Δs is thus given by

 $\Delta I(t) = (q/\Delta t)(\Delta d\Delta s\cos\theta)N(\theta)d\theta = [\partial \boldsymbol{p}(t)/\partial t] \cdot \Delta s N(\theta)d\theta.$

The corresponding current density, $\Delta J(t) = [\partial p(t)/\partial t]N(\theta)d\theta$, integrated over all possible angles θ , then yields

$$\boldsymbol{J}(t) = \frac{\partial}{\partial t} \int_0^{\pi} \boldsymbol{p}(t) N(\theta) \,\mathrm{d}\theta = \frac{\partial \boldsymbol{P}(t)}{\partial t},$$

where P(t) is the local polarization density. We may thus write, quite generally, $J(r,t) = \partial P(r,t) / \partial t$.

Note that elongation of dipoles, i.e., change of dipole moment p through changes in d, is but one mechanism by which the polarization P of a given medium can become time-dependent. Other possible mechanisms include: charge transfer in time (i.e., q is time-dependent), rotation of dipoles (i.e., θ becomes time-dependent), and motion of entire dipoles through the surface element Δs . It can be shown that all these mechanisms, individually or in combination, make a contribution to the current-density J that can be expressed in the aforementioned form, namely, $J(\mathbf{r},t) = \partial P(\mathbf{r},t)/\partial t$.