Problem 2) As in Problem 1, the volume of interest is Δ*d*Δ*s*cosθ, of which one half is to the left of Δs and one half to the right. Here Δd is the change of $d(t)$ over a time interval Δt , that is, $\Delta d = d(t + \Delta t) - d(t)$. If a dipole is located within this volume, one of its charges (either +*q* or –*q*) will cross the surface Δ*s* during the time interval Δ*t*. Irrespective of which charge crosses, the current flowing from left to right will be *q*/Δ*t*. The current crossing Δ*s* due to all dipoles that are oriented at an angle θ to Δs is thus given by

 $\Delta I(t) = (q/\Delta t)(\Delta d \Delta s \cos \theta) N(\theta) d\theta = [\partial p(t)/\partial t] \cdot \Delta s N(\theta) d\theta$.

The corresponding current density, $\Delta \mathbf{J}(t) = [\partial \mathbf{p}(t) / \partial t] N(\theta) d\theta$, integrated over all possible angles θ , then yields

$$
\boldsymbol{J}(t) = \frac{\partial}{\partial t} \int_0^{\pi} \boldsymbol{p}(t) N(\theta) \, \mathrm{d}\theta = \frac{\partial \boldsymbol{P}(t)}{\partial t},
$$

where $P(t)$ is the local polarization density. We may thus write, quite generally, *J*(*r*,*t*) =∂*P*(*r*,*t*)/∂*t*.

Note that elongation of dipoles, i.e., change of dipole moment *p* through changes in *d*, is but one mechanism by which the polarization *P* of a given medium can become time-dependent. Other possible mechanisms include: charge transfer in time (i.e., *q* is time-dependent), rotation of dipoles (i.e., θ becomes time-dependent), and motion of entire dipoles through the surface element Δ*s*. It can be shown that all these mechanisms, individually or in combination, make a contribution to the current-density *J* that can be expressed in the aforementioned form, namely, *J*(\bm{r}, t) =∂ $\bm{P}(\bm{r}, t)$ /∂ t .