Problem 18)

Solutions



From symmetry, it is clear that the vector potential A is everywhere in the $\hat{\phi}$ direction, and that the contribution of the volume element shown in the figure is $J_{so} \sin \phi' R \Delta \phi' \Delta z'$. Therefore,

$$A_{\phi}(r) = \frac{\mu_{o}}{4\pi} \int_{\phi'=0}^{2\pi} \int_{z'=-\infty}^{\infty} \frac{J_{so}R\sin\phi'}{\sqrt{R^{2} + r^{2} - 2rR\sin\phi' + z'^{2}}} d\phi' dz'$$
$$= \frac{\mu_{o}J_{so}R}{2\pi} \int_{\phi'=0}^{2\pi} \sin\phi' d\phi' \int_{z'=0}^{\infty} \frac{dz'}{\sqrt{R^{2} + r^{2} - 2rR\sin\phi' + z'^{2}}}$$

Depending on the value of ϕ' , the expression $(R^2 + r^2 - 2rR\sin\phi')$ is somewhere between $(R-r)^2$ and $(R+r)^2$, thus ensuring that it is always non-negative.

Now, from the Table of Integrals (Gradshteyn & Ryzhik, 2.261) we find

$$\int_{z'=0}^{z_{o}} \frac{dz'}{\sqrt{\alpha + z'^{2}}} = \ln\left(z' + \sqrt{\alpha + z'^{2}}\right)\Big|_{z'=0}^{z_{o}} = \ln\left(z_{o} + \sqrt{\alpha + z_{o}^{2}}\right) - \ln\sqrt{\alpha}; \qquad \alpha \ge 0$$

When $z_0 \rightarrow \infty$, the first term on the right-hand-side of the above equation becomes very large, but its variations with α become insignificant. We can, therefore, drop this term, which does not vary with α , even though its magnitude is infinite. We will have

$$A_{\phi}(r) = -\frac{\mu_{o}J_{so}R}{4\pi} \int_{\phi'=0}^{2\pi} \sin\phi' \ln(R^{2} + r^{2} - 2rR\sin\phi') d\phi'$$

= $-\frac{\mu_{o}J_{so}R}{4\pi} \Big\{ \int_{\phi'=0}^{2\pi} \sin\phi' \ln(R^{2}) d\phi' + \int_{\phi'=0}^{2\pi} \sin\phi' \ln[1 - 2(r/R)\sin\phi' + (r/R)^{2}] d\phi' \Big\}.$

The first integral on the right-hand-side of the above equation is zero. As for the second integral, when ϕ' is replaced by $(\frac{1}{2}\pi - \phi')$, $\sin \phi'$ changes to $\cos \phi'$, but nothing else changes because the integral covers one full period of the sine and cosine functions. We will have

$$A_{\phi}(r) = -\frac{\mu_{o}J_{so}R}{2\pi} \int_{\phi'=0}^{\pi} \cos\phi' \ln\left[1 - 2(r/R)\cos\phi' + (r/R)^{2}\right] d\phi' = -\frac{\mu_{o}J_{so}R}{2\pi} \begin{cases} -\pi(r/R); & r \le R, \\ -\pi(R/r); & r > R. \end{cases}$$

Therefore,

$$A_{\phi}(r) = \begin{cases} \frac{1}{2} \mu_{o} J_{so} r; & r \leq R, \\ \frac{1}{2} \mu_{o} J_{so} R^{2} / r; & r > R. \end{cases}$$