

## Problem 9)

In problem 7, we calculated the potential function  $\Psi(r)$  for a uniformly-charged hollow sphere. We now use that function to obtain the potential function for a uniformly charged solid sphere of radius  $R$  and charge density  $\rho_0$ . We must integrate the contributions of nested spherical shells having radii from  $r' = 0$  to  $r' = R$ . The charge content of each such shell is  $Q(r') = (4\pi\rho_0 r'^2) dr'$ .

At a point  $\vec{r}$  outside the sphere (i.e.,  $r \geq R$ ), we'll have:

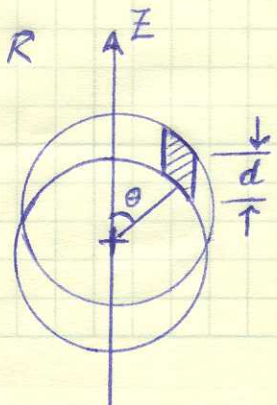
$$\Psi(r) = \frac{1}{4\pi\epsilon_0 r} \int_{r'=0}^R (4\pi\rho_0 r'^2) dr' = \frac{\rho_0 R^3}{3\epsilon_0 r}; \quad r \geq R$$

At a point  $\vec{r}$  inside the sphere (i.e.,  $r < R$ ), we'll have:

$$\begin{aligned} \Psi(r) &= \frac{1}{4\pi\epsilon_0 r} \int_{r'=0}^r (4\pi\rho_0 r'^2) dr' + \frac{1}{4\pi\epsilon_0} \int_{r'=r}^R \frac{(4\pi\rho_0 r'^2)}{r'} dr' \\ &= \frac{\rho_0 r^2}{3\epsilon_0} + \frac{\rho_0 (R^2 - r^2)}{2\epsilon_0} \Rightarrow \Psi(r) = \frac{\rho_0}{2\epsilon_0} \left( R^2 - \frac{r^2}{3} \right); \quad r \leq R \end{aligned}$$

$$\vec{E}(r) = -\vec{\nabla}\Psi(r) = -\frac{\partial\Psi}{\partial r} \hat{r} = \begin{cases} +\frac{\rho_0 r}{3\epsilon_0} \hat{r} & r \leq R \\ +\frac{\rho_0 R^3}{3\epsilon_0 r^2} \hat{r} & r \geq R \end{cases}$$

Consider two spheres of radius  $R$ , shifted by a short distance  $d$  along the  $z$ -axis ( $d \ll R$ ).



Next, Consider a point on the surface at polar coordinates  $(\theta, \phi)$ , and let a small patch of the surface area around this point be defined by small increments  $\Delta\theta, \Delta\phi$ . The area of the small patch is thus given by  $\Delta S = R^2 \sin\theta \Delta\theta \Delta\phi$ .

With reference to the diagram at the bottom of the previous page, the area of the "charged" parallelogram is given by

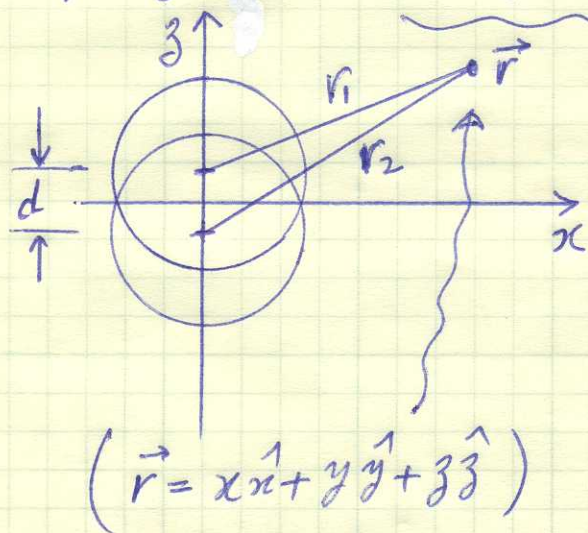
$(R \Delta\theta) d \cos\theta$ . In the direction perpendicular to the plane of the paper, the thickness of this parallelogram is  $R \sin\theta \Delta\phi$ .

The volume of the charged parallelepiped, therefore, is

$R^2 d \sin\theta \cos\theta \Delta\theta \Delta\phi$ . This volume has a charge density of  $\rho_0$ , resulting in a charge content  $\Delta Q = R^2 d \rho_0 \sin\theta \cos\theta \Delta\theta \Delta\phi$ .

Dividing the charge  $\Delta Q$  by the surface area  $\Delta S$  of the patch yields the surface charge density  $\sigma(\theta) = \rho_0 d \cos\theta$ .

Comparing with problem 4, we find that  $\sigma_0 = \rho_0 d$ .



$$r_1^2 = x^2 + y^2 + (z - \frac{1}{2}d)^2$$

$$= x^2 + y^2 + z^2 - 3d + \frac{1}{4}d^2$$

$$= r^2 - 3d + \frac{1}{4}d^2 = r^2 \left[ 1 - \left(\frac{3d}{r^2}\right) + \left(\frac{d}{2r}\right)^2 \right]$$

$$\Rightarrow r_1 = r \sqrt{1 - \left(\frac{3d}{r^2}\right) + \left(\frac{d}{2r}\right)^2} \Rightarrow$$

$$r_1 \approx r \left[ 1 - \frac{1}{2} \left(\frac{3d}{r^2}\right) + \frac{1}{2} \left(\frac{d}{2r}\right)^2 \right] \approx r - \frac{3d}{2r}$$

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Similarly,  $r_2^2 = r^2 + 3d + \frac{1}{4}d^2$  and  $r_2 \approx r + \frac{3d}{2r}$ .

For the overlapping spheres one must add the corresponding potential functions, thus:  $\hat{\Psi}(\vec{r}) = \Psi_1(\vec{r}) + \Psi_2(\vec{r}) \Rightarrow \hat{\Psi}(\vec{r}) = \Psi(r_1) - \Psi(r_2)$ . Here  $\hat{\Psi}(\vec{r})$  is the potential function for the overlapping spheres, and the minus sign before  $\Psi(r_2)$  is due to the fact that the lower sphere is negatively charged.

Therefore, for  $r \leq R$ , we'll have:

$$\hat{\Psi}(\vec{r}) = \frac{\rho_0}{2\epsilon_0} \left( R^2 - \frac{r_1^2}{3} \right) - \frac{\rho_0}{2\epsilon_0} \left( R^2 - \frac{r_2^2}{3} \right) = \frac{\rho_0}{6\epsilon_0} (r_2^2 - r_1^2) \Rightarrow$$

$$\hat{\Psi}(\vec{r}) = \frac{\rho_0}{3\epsilon_0} 3d = \frac{\rho_0 d}{3\epsilon_0} z = \frac{\sigma_0}{3\epsilon_0} z = \frac{\sigma_0}{3\epsilon_0} r \cos \theta; \quad r \leq R$$

For  $r \geq R$  we have:

$$\hat{\Psi}(\vec{r}) = \frac{\rho_0 R^3}{3\epsilon_0 r_1} - \frac{\rho_0 R^3}{3\epsilon_0 r_2} = \frac{\rho_0 R^3}{3\epsilon_0} \left( \frac{1}{r_1} - \frac{1}{r_2} \right) = \frac{\rho_0 R^3}{3\epsilon_0} \frac{(r_2 - r_1)}{r_1 r_2}$$

$$\approx \frac{\rho_0 R^3}{3\epsilon_0} \frac{3d/r}{r^2} = \frac{\rho_0 d R^3}{3\epsilon_0} \frac{3/r}{r^2} = \frac{\sigma_0 R^3}{3\epsilon_0} \frac{\cos \theta}{r^2}; \quad r \geq R$$

The above results, which are valid when  $d \ll R$ , are consistent with the results of problem 4.