

**Problem 4)**

From Maxwell's 3<sup>rd</sup> equation:  $\nabla \times \mathbf{E}(\mathbf{r}) = 0 \rightarrow \mathbf{E}(\mathbf{r}) = -\nabla\psi(\mathbf{r})$  because  $\nabla \times \nabla\psi(\mathbf{r}) = 0$ .

From Maxwell's 1<sup>st</sup> equation:  $\nabla \cdot \mathbf{D}(\mathbf{r}) = \rho_{\text{free}}(\mathbf{r}) \rightarrow \epsilon_0 \nabla \cdot \mathbf{E}(\mathbf{r}) = \rho_{\text{free}}(\mathbf{r}) - \nabla \cdot \mathbf{P}(\mathbf{r})$   
 $\rightarrow \epsilon_0 \nabla \cdot \mathbf{E}(\mathbf{r}) = \rho_{\text{free}}(\mathbf{r}) + \rho_{\text{bound}}(\mathbf{r}) = \rho_{\text{total}}(\mathbf{r})$ .

Combining the above equations then yields  $\nabla^2\psi(\mathbf{r}) = -\rho_{\text{total}}(\mathbf{r})/\epsilon_0$ , whose Fourier transform is  $\psi(\mathbf{k}) = \rho_{\text{total}}(\mathbf{k})/(\epsilon_0 k^2)$ . In Chapter 3, Problem 4(a), it was found that the 3D Fourier transform of  $f(\mathbf{r}) = 1/|\mathbf{r}|$  is  $F(\mathbf{k}) = 4\pi/k^2$ . We may thus write

$$\begin{aligned} \psi(\mathbf{r}) &= \mathcal{F}^{-1}\{\psi(\mathbf{k})\} = (2\pi)^{-3} \int_{-\infty}^{\infty} \psi(\mathbf{k}) \exp(i\mathbf{k} \cdot \mathbf{r}) d\mathbf{k} \\ &= (2\pi)^{-3} \int_{-\infty}^{\infty} \frac{\rho_{\text{total}}(\mathbf{k}) \exp(i\mathbf{k} \cdot \mathbf{r})}{\epsilon_0 k^2} d\mathbf{k} \quad \boxed{\mathcal{F}\{1/|\mathbf{r}|\}} \\ &= (2\pi)^{-3} (4\pi\epsilon_0)^{-1} \int_{-\infty}^{\infty} \rho_{\text{total}}(\mathbf{k}) \exp(i\mathbf{k} \cdot \mathbf{r}) \int_{-\infty}^{\infty} \frac{\exp(-i\mathbf{k} \cdot \mathbf{r}'')}{|\mathbf{r}''|} d\mathbf{r}'' d\mathbf{k} \\ &= (2\pi)^{-3} (4\pi\epsilon_0)^{-1} \int_{-\infty}^{\infty} \frac{1}{|\mathbf{r}''|} \int_{-\infty}^{\infty} \rho_{\text{total}}(\mathbf{k}) \exp[i\mathbf{k} \cdot (\mathbf{r} - \mathbf{r}'')] d\mathbf{k} d\mathbf{r}'' \\ &= (4\pi\epsilon_0)^{-1} \int_{-\infty}^{\infty} \frac{\rho_{\text{total}}(\mathbf{r} - \mathbf{r}'')}{|\mathbf{r}''|} d\mathbf{r}'' \\ &= (4\pi\epsilon_0)^{-1} \int_{-\infty}^{\infty} \frac{\rho_{\text{total}}(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}'. \quad \leftarrow \boxed{\text{Change of variable: } \mathbf{r}' = \mathbf{r} - \mathbf{r}''} \end{aligned}$$

The  $E$ -field is now obtained from the above scalar potential using  $\mathbf{E}(\mathbf{r}) = -\nabla\psi(\mathbf{r})$ , as follows:

$$\mathbf{E}(\mathbf{r}) = -(4\pi\epsilon_0)^{-1} \int_{-\infty}^{\infty} \rho_{\text{total}}(\mathbf{r}') \nabla(|\mathbf{r} - \mathbf{r}'|^{-1}) d\mathbf{r}' = (4\pi\epsilon_0)^{-1} \int_{-\infty}^{\infty} \frac{\rho_{\text{total}}(\mathbf{r}')(\mathbf{r} - \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} d\mathbf{r}'. \quad \leftarrow \boxed{\text{Under the integral, the } \nabla \text{ operator acts on } \mathbf{r}, \text{ treating } \mathbf{r}' \text{ as a constant.}}$$

Alternatively, one could observe that Maxwell's first equation,  $\epsilon_0 \nabla \cdot \mathbf{E}(\mathbf{r}) = \rho_{\text{total}}(\mathbf{r})$ , yields the longitudinal component of the  $E$ -field, as follows:

$$i\mathbf{k} \cdot \mathbf{E}(\mathbf{k}) = \rho_{\text{total}}(\mathbf{k})/\epsilon_0 \rightarrow \mathbf{E}_{\parallel}(\mathbf{k}) = -\frac{i\rho_{\text{total}}(\mathbf{k})\hat{\mathbf{k}}}{\epsilon_0 k}$$

However, Maxwell's third equation,  $\nabla \times \mathbf{E}(\mathbf{r}) = 0$ , implies that the transverse component of  $\mathbf{E}(\mathbf{r})$  is zero, and that, therefore,  $\mathbf{E}(\mathbf{r})$  is purely longitudinal. Consequently,  $\mathbf{E}(\mathbf{k}) = \mathbf{E}_{\parallel}(\mathbf{k})$ , and we have

$$\mathbf{E}(\mathbf{r}) = \mathcal{F}^{-1}\{\mathbf{E}(\mathbf{k})\} = (2\pi)^{-3} \int_{-\infty}^{\infty} \mathbf{E}(\mathbf{k}) \exp(i\mathbf{k} \cdot \mathbf{r}) d\mathbf{k} = -i(2\pi)^{-3} \int_{-\infty}^{\infty} \frac{\rho_{\text{total}}(\mathbf{k})\hat{\mathbf{k}}}{\epsilon_0 k} \exp(i\mathbf{k} \cdot \mathbf{r}) d\mathbf{k}.$$

Now, the 3D Fourier transform of  $f(\mathbf{r}) = -\hat{\mathbf{r}}/r^2$  was found in Chapter 3, Problem 4(d), to be  $F(\mathbf{k}) = 4\pi i \hat{\mathbf{k}}/k$ . We may thus write

$$\mathcal{F}\{\hat{\mathbf{r}}/r^2\} = \mathcal{F}\{\mathbf{r}/|\mathbf{r}|^3\}$$



$$\begin{aligned} \mathbf{E}(\mathbf{r}) &= (2\pi)^{-3} (4\pi\epsilon_0)^{-1} \int_{-\infty}^{\infty} \rho_{\text{total}}(\mathbf{k}) \exp(i\mathbf{k} \cdot \mathbf{r}) \int_{-\infty}^{\infty} \frac{\mathbf{r}'' \exp(-i\mathbf{k} \cdot \mathbf{r}'')}{|\mathbf{r}''|^3} d\mathbf{r}'' d\mathbf{k} \\ &= (2\pi)^{-3} (4\pi\epsilon_0)^{-1} \int_{-\infty}^{\infty} \frac{\mathbf{r}''}{|\mathbf{r}''|^3} \int_{-\infty}^{\infty} \rho_{\text{total}}(\mathbf{k}) \exp[i\mathbf{k} \cdot (\mathbf{r} - \mathbf{r}'')] d\mathbf{k} d\mathbf{r}'' \\ &= (4\pi\epsilon_0)^{-1} \int_{-\infty}^{\infty} \frac{\mathbf{r}'' \rho_{\text{total}}(\mathbf{r} - \mathbf{r}'')}{|\mathbf{r}''|^3} d\mathbf{r}'' \\ &= (4\pi\epsilon_0)^{-1} \int_{-\infty}^{\infty} \frac{(\mathbf{r} - \mathbf{r}') \rho_{\text{total}}(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} d\mathbf{r}'. \quad \leftarrow \text{Change of variable: } \mathbf{r}' = \mathbf{r} - \mathbf{r}'' \end{aligned}$$

This is the same result as obtained previously by applying the gradient operator to the scalar potential. Either way, the  $E$ -field is seen to be computed by applying Coulomb's law to individual volume elements of charge, then integrating over the entire space. Finally, the  $D$ -field is obtained by adding  $\mathbf{P}(\mathbf{r})$  to  $\epsilon_0$  times the above  $E$ -field, that is,

$$\mathbf{D}(\mathbf{r}) = \epsilon_0 \mathbf{E}(\mathbf{r}) + \mathbf{P}(\mathbf{r}).$$


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