

**Problem 4.40)**

$$\text{a) } \mathbf{E}(\mathbf{r}) = -\nabla\psi(\mathbf{r}) = -\frac{\partial\psi}{\partial r}\hat{\mathbf{r}} - \frac{\partial\psi}{r\partial\theta}\hat{\boldsymbol{\theta}} = \begin{cases} -(P_0/3\epsilon_0)(\cos\theta\hat{\mathbf{r}} - \sin\theta\hat{\boldsymbol{\theta}}); & r < R, \\ (P_0R^3/3\epsilon_0)(2\cos\theta\hat{\mathbf{r}} + \sin\theta\hat{\boldsymbol{\theta}})/r^3; & r \geq R. \end{cases}$$

The field inside the sphere may be further simplified and written as  $\mathbf{E}(\mathbf{r}) = -(P_0/3\epsilon_0)\hat{\mathbf{z}}$ .

$$\begin{aligned} \text{b) } \rho_{\text{bound}}^{(e)}(\mathbf{r}) &= -\nabla \cdot \mathbf{P}(\mathbf{r}) = -\nabla \cdot [P_0 \text{Sphere}(r/R)(\cos\theta\hat{\mathbf{r}} - \sin\theta\hat{\boldsymbol{\theta}})] \\ &= -\frac{\partial(r^2 P_r)}{r^2 \partial r} - \frac{\partial(\sin\theta P_\theta)}{r \sin\theta \partial\theta} \\ &= -\left(\frac{P_0 \cos\theta}{r^2}\right) \frac{\partial[r^2 \text{Sphere}(r/R)]}{\partial r} + \left[\frac{P_0 \text{Sphere}(r/R)}{r \sin\theta}\right] \frac{\partial \sin^2\theta}{\partial\theta} \\ &= -\frac{P_0 \cos\theta [2r \text{Sphere}(r/R) - r^2 \delta(r-R)]}{r^2} + \frac{2P_0 \text{Sphere}(r/R) \cos\theta}{r} = P_0 \delta(r-R) \cos\theta. \end{aligned}$$

Because of the  $\delta$ -function appearing in the above charge-density profile, we can state that the sphere has a bound *surface-charge-density*  $\sigma_s(r=R, \theta, \phi) = P_0 \cos\theta$ . Note that the surface-charge-density is positive on the upper hemisphere and negative on the lower hemisphere.

c) The parallel (or tangential) component of the  $E$ -field at the surface of the sphere is  $E_\theta$ , which is found in (a) to be equal to  $(P_0/3\epsilon_0) \sin\theta$  immediately inside as well as immediately outside the sphere. The continuity requirement for the tangential  $E$ -field is, therefore, satisfied.

The perpendicular component of the  $D$ -field inside the sphere is given by

$$D_\perp = \epsilon_0 E_r + P_r = -\frac{1}{3}P_0 \cos\theta + P_0 \cos\theta = \frac{2}{3}P_0 \cos\theta.$$

Outside the sphere and immediately above the surface, we have  $D_\perp = \epsilon_0 E_r = \frac{2}{3}P_0 \cos\theta$ . Therefore, in the absence of free surface-charge-density, the continuity of  $D_\perp$  is confirmed.

If, instead of  $D_\perp$ , we examine the perpendicular component of the  $E$ -field, we find, at the surface of the sphere, the following discontinuity in  $E_\perp$ :

$$\begin{aligned} E_r(r=R^+, \theta, \phi) - E_r(r=R^-, \theta, \phi) &= (2P_0/3\epsilon_0) \cos\theta - (-P_0/3\epsilon_0) \cos\theta \\ &= (P_0/\epsilon_0) \cos\theta. \end{aligned}$$

This, however, is precisely equal to the bound surface-charge-density  $\sigma_s = P_0 \cos\theta$  found in part (b), divided by  $\epsilon_0$ , which is, once again, consistent with the boundary condition derived from Maxwell's 1<sup>st</sup> equation,  $\nabla \cdot \mathbf{E}(\mathbf{r}, t) = \epsilon_0^{-1} \rho_{\text{total}}^{(e)}(\mathbf{r}, t)$ .