

Problem 20)

$$\vec{\nabla}_0 \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0 \Rightarrow \frac{\partial \rho}{\partial t} = -\frac{\partial J_z}{\partial z} = -J_{s_0} \delta(y) e^{\omega'' t} \kappa \sin(\omega t - \kappa z) \Rightarrow$$

$$\frac{\partial \rho}{\partial t} = -\kappa J_{s_0} \delta(y) \frac{e^{-i\kappa z} e^{+i(\omega' - i\omega'')t} - e^{+i\kappa z} e^{-i(\omega' + i\omega'')t}}{2i} \Rightarrow$$

$$\rho(\vec{r}, t) = \frac{1}{2} \kappa J_{s_0} \delta(y) \left[\frac{e^{i\kappa z} e^{-i\omega t}}{\omega_0} + \frac{e^{-i\kappa z} e^{i\omega_0^* t}}{\omega_0^*} \right]$$

$$\rho(\vec{k}, t) = \int_{-\infty}^{\infty} \rho(\vec{r}, t) e^{-i\vec{k} \cdot \vec{r}} d\vec{r} = 2\pi^2 \kappa J_{s_0} \delta(k_x) \left[\frac{\delta(k_z - \kappa) e^{-i\omega t}}{\omega_0} + \frac{\delta(k_z + \kappa) e^{+i\omega_0^* t}}{\omega_0^*} \right]$$

a) We find the scalar potential $\psi(\vec{k}, t)$ and its inverse Fourier transform $\psi(\vec{r}, t)$ first for the term containing ω_0 , then for the term containing ω_0^* , and finally add

→ them up after allowing $\omega'' \rightarrow 0$. (See Problem 40 for a discussion of the choice of sign.)

$$\psi_1(\vec{r}, t) = \frac{1}{(2\pi)^3 \epsilon_0} \int_{-\infty}^{\infty} \frac{\rho(\vec{k}, t)}{k^2 - (\omega/c)^2} e^{+i\vec{k} \cdot \vec{r}} d\vec{k} = \frac{\kappa J_{s_0}}{4\pi \epsilon_0 \omega_0} e^{-i\omega_0 t} e^{i\kappa z} \int_{-\infty}^{\infty} \frac{e^{ik_y y}}{k_y^2 + \kappa^2 - (\omega_0/c)^2} dk_y$$

$$= \begin{cases} \frac{\kappa J_{s_0}}{4\epsilon_0 \omega_0} \frac{\exp(-i\omega_0 t + i\kappa z + i\sqrt{(\omega_0/c)^2 - \kappa^2} |y|)}{-i\sqrt{(\omega_0/c)^2 - \kappa^2}}; & \kappa < \omega_0/c \\ \frac{\kappa J_{s_0}}{4\epsilon_0 \omega_0} \frac{\exp(-i\omega_0 t + i\kappa z) \exp(-\sqrt{\kappa^2 - (\omega_0/c)^2} |y|)}{\sqrt{\kappa^2 - (\omega_0/c)^2}}; & \kappa > \omega_0/c \end{cases}$$

b) Similarly, we find $\psi_2(\vec{r}, t)$ for the term containing ω_0^* and allow $\omega'' \rightarrow 0$.

$$\psi_2(\vec{r}, t) = \begin{cases} \frac{\kappa J_{s_0}}{4\epsilon_0 \omega_0} \frac{\exp(i\omega_0 t - i\kappa z - i\sqrt{(\omega_0/c)^2 - \kappa^2} |y|)}{i\sqrt{(\omega_0/c)^2 - \kappa^2}}; & \kappa < \omega_0/c \\ \frac{\kappa J_{s_0}}{4\epsilon_0 \omega_0} \frac{\exp(i\omega_0 t - i\kappa z) \exp(-\sqrt{\kappa^2 - (\omega_0/c)^2} |y|)}{\sqrt{\kappa^2 - (\omega_0/c)^2}}; & \kappa > \omega_0/c \end{cases}$$

Adding the above expressions yields the final form of the scalar potential $\psi(\vec{r}, t)$.

$$\psi(\vec{r}, t) = \psi_1 + \psi_2 = \begin{cases} \frac{\kappa J_{s_0}}{2\epsilon_0 \omega_0 \sqrt{(\omega_0/c)^2 - \kappa^2}} \sin[\omega_0 t - \kappa z - \sqrt{(\omega_0/c)^2 - \kappa^2} |y|]; & \kappa < \omega_0/c \\ \frac{\kappa J_{s_0}}{2\epsilon_0 \omega_0 \sqrt{\kappa^2 - (\omega_0/c)^2}} \exp(-\sqrt{\kappa^2 - (\omega_0/c)^2} |y|) \cos(\omega_0 t - \kappa z); & \kappa > \omega_0/c \end{cases}$$

A similar procedure is followed for determining the vector potential $\vec{A}(\vec{r}, t)$.

$$\vec{J}(\vec{r}, t) = \frac{1}{2} J_{s_0} \delta(y) \hat{z} [\exp(-i\omega_0 t + \kappa z) + \exp(i\omega_0^* t - \kappa z)]$$

$$\vec{J}(\vec{k}, t) = \int_{-\infty}^{\infty} \vec{J}(\vec{r}, t) e^{-i\vec{k}\cdot\vec{r}} d\vec{r} = 2\pi^2 J_{s_0} \delta(k_x) \hat{z} [\delta(k_z - \kappa) e^{-i\omega_0 t} + \delta(k_z + \kappa) e^{+i\omega_0^* t}]$$

$$\vec{A}_1(\vec{r}, t) = \frac{\mu_0}{(2\pi)^3} \int_{-\infty}^{\infty} \frac{\vec{J}_1(\vec{k}, t)}{k^2 - (\omega_0/c)^2} e^{+i\vec{k}\cdot\vec{r}} dk = \frac{\mu_0 J_{s_0} \hat{z}}{4\pi} e^{-i\omega_0 t} e^{i\kappa z} \int_{-\infty}^{\infty} \frac{e^{ik_y y}}{k_y^2 + \kappa^2 - (\omega_0/c)^2} dk_y$$

$$= \begin{cases} \frac{1}{4} \mu_0 J_{s_0} \hat{z} \frac{\exp(-i\omega_0 t + i\kappa z + i\sqrt{(\omega_0/c)^2 - \kappa^2} |y|)}{-i\sqrt{(\omega_0/c)^2 - \kappa^2}}; & \kappa < \omega_0/c \\ \frac{1}{4} \mu_0 J_{s_0} \hat{z} \frac{\exp(-i\omega_0 t + i\kappa z) \exp(-\sqrt{\kappa^2 - (\omega_0/c)^2} |y|)}{\sqrt{\kappa^2 - (\omega_0/c)^2}}; & \kappa > \frac{\omega_0}{c} \end{cases}$$

A similar expression is found for the term containing ω_0^* . When the two expressions are added together, we find:

$$\vec{A}(\vec{r}, t) = \vec{A}_1(\vec{r}, t) + \vec{A}_2(\vec{r}, t) = \begin{cases} \frac{\mu_0 J_{s_0} \hat{z}}{2\sqrt{(\omega_0/c)^2 - \kappa^2}} \sin[\omega_0 t - \kappa z - \sqrt{(\omega_0/c)^2 - \kappa^2} |y|]; & \kappa < \frac{\omega_0}{c} \\ \frac{\mu_0 J_{s_0} \hat{z}}{2\sqrt{\kappa^2 - (\omega_0/c)^2}} \exp(-\sqrt{\kappa^2 - (\omega_0/c)^2} |y|) \cos(\omega_0 t - \kappa z); & \kappa > \frac{\omega_0}{c} \end{cases}$$

$$\vec{E}(\vec{r}, t) = -\vec{\nabla}\psi - \frac{\partial \vec{A}}{\partial t} = -\frac{\partial \psi}{\partial y} \hat{y} - \left(\frac{\partial \psi}{\partial z} + \frac{\partial A_z}{\partial t}\right) \hat{z} \Rightarrow$$

$$E_y(\vec{r}, t) = -\frac{\partial \psi}{\partial y} = \begin{cases} \frac{\kappa J_{s_0} \text{sign}(y)}{2\epsilon_0 \omega_0} \cos[\omega_0 t - \kappa z - \sqrt{(\omega_0/c)^2 - \kappa^2} |y|]; & \kappa < \omega_0/c \\ \frac{\kappa J_{s_0} \text{sign}(y)}{2\epsilon_0 \omega_0} \exp(-\sqrt{\kappa^2 - (\omega_0/c)^2} |y|) \cos(\omega_0 t - \kappa z); & \kappa > \omega_0/c. \end{cases}$$

$$E_z(\vec{r}, t) = -\left(\frac{\partial \psi}{\partial z} + \frac{\partial A_z}{\partial t}\right) = \begin{cases} -\frac{J_{s_0} \sqrt{(\omega_0/c)^2 - \kappa^2}}{2\epsilon_0 \omega_0} \cos[\omega_0 t - \kappa z - \sqrt{(\omega_0/c)^2 - \kappa^2} |y|]; & \kappa < \omega_0/c \\ -\frac{J_{s_0} \sqrt{\kappa^2 - (\omega_0/c)^2}}{2\epsilon_0 \omega_0} \exp(-\sqrt{\kappa^2 - (\omega_0/c)^2} |y|) \sin(\omega_0 t - \kappa z); & \kappa > \omega_0/c \end{cases}$$

$$\vec{H}(\vec{r}, t) = \frac{1}{\mu_0} \vec{\nabla} \times \vec{A}(\vec{r}, t) = \frac{1}{\mu_0} \frac{\partial A_z}{\partial y} \hat{x} \Rightarrow$$

$$H_x(\vec{r}, t) = \begin{cases} -\frac{1}{2} J_{s_0} \text{sign}(y) \cos[\omega_0 t - \kappa z - \sqrt{(\omega_0/c)^2 - \kappa^2} |y|]; & \kappa < \omega_0/c \\ -\frac{1}{2} J_{s_0} \text{sign}(y) \exp(-\sqrt{\kappa^2 - (\omega_0/c)^2} |y|) \cos(\omega_0 t - \kappa z); & \kappa > \omega_0/c \end{cases}$$

Note that E_y and H_x satisfy the boundary conditions at $y=0$.

$$\vec{S}(\vec{r}, t) = \vec{E}(\vec{r}, t) \times \vec{H}(\vec{r}, t) = (E_z H_x) \hat{y} - (E_y H_x) \hat{z} \Rightarrow$$

$$S_y(\vec{r}, t) = \begin{cases} \frac{1}{4} \epsilon_0 J_{s_0}^2 \sin(y) \sqrt{1 - (ck/\omega)^2} \cos^2[\omega_0 t - kz - \sqrt{(\omega_0/c)^2 - k^2} |y|]; & k < \omega_0/c \\ \frac{1}{4} \epsilon_0 J_{s_0}^2 \sin(y) \sqrt{(ck/\omega)^2 - 1} \exp[-2\sqrt{k^2 - (\omega_0/c)^2} |y|] \sin(\omega_0 t - kz) \cos(\omega_0 t - kz); & k > \omega_0/c \end{cases}$$

$$\Rightarrow \langle S_y(\vec{r}, t) \rangle = \begin{cases} \frac{1}{8} \epsilon_0 J_{s_0}^2 \sin(y) \sqrt{1 - (ck/\omega)^2}; & k < \omega_0/c \\ 0; & k > \omega_0/c \end{cases}$$

$$S_z(\vec{r}, t) = -E_y H_x = \begin{cases} \frac{1}{4} \epsilon_0 J_{s_0}^2 (ck/\omega) \cos^2[\omega_0 t - kz - \sqrt{(\omega_0/c)^2 - k^2} |y|]; & k < \omega_0/c \\ \frac{1}{4} \epsilon_0 J_{s_0}^2 (ck/\omega) \exp(-2\sqrt{k^2 - (\omega_0/c)^2} |y|) \cos^2(\omega_0 t - kz); & k > \omega_0/c \end{cases}$$

$$\Rightarrow \langle S_z(\vec{r}, t) \rangle = \begin{cases} \frac{1}{8} \epsilon_0 J_{s_0}^2 (ck/\omega); & k < \omega_0/c \\ \frac{1}{8} \epsilon_0 J_{s_0}^2 (ck/\omega) \exp(-2\sqrt{k^2 - (\omega_0/c)^2} |y|); & k > \omega_0/c. \end{cases}$$

