

## Problem 14)

Point-dipole oscillator  $\vec{P}(\vec{r}, t) = P_0 \hat{\vec{z}} \delta(\vec{r}) e^{i\omega'' t} \cos(\omega' t) \Rightarrow$

$$\vec{P}(\vec{r}, t) = \frac{1}{2} P_0 \hat{\vec{z}} \delta(\vec{r}) e^{i\omega'' t} (e^{i\omega' t} + e^{-i\omega' t}) = \frac{1}{2} P_0 \hat{\vec{z}} \delta(\vec{r}) e^{-i(\omega' + \omega'') t} + \text{c.c.}$$

Defining  $\omega = \omega' + i\omega''$ , where  $\omega'' \geq 0$ , we will find the scalar and vector potentials for  $\vec{P}(\vec{r}, t) = P_0 \hat{\vec{z}} \delta(\vec{r}) e^{-i\omega t}$ , whose spatial Fourier transform is given by

$\vec{P}(\vec{k}, t) = P_0 \hat{\vec{z}} e^{-i\omega t}$ . The spatial parts of  $\psi(\vec{k}, t)$  and  $\vec{A}(\vec{k}, t)$  are thus given by

$$\psi(\vec{k}) = -\frac{iP_0}{\epsilon_0} \frac{\vec{k} \cdot \hat{\vec{z}}}{k^2 - \omega'^2/c^2}; \quad \vec{A}(\vec{k}) = -i\omega \mu_0 P_0 \frac{\hat{\vec{z}}}{k^2 - \omega'^2/c^2}; \quad (\text{See Problem 16}).$$

$$\begin{aligned} \text{Inverse Fourier Transformation: } \psi(\vec{r}) &= -\frac{iP_0}{\epsilon_0} \frac{1}{(2\pi)^3} \int_{-\infty}^{\infty} \frac{k_3}{k^2 - \omega'^2/c^2} e^{+ik \cdot \vec{r}} dk \\ &= -\frac{P_0}{\epsilon_0} \frac{1}{(2\pi)^3} \frac{\partial}{\partial \vec{z}} \int_{-\infty}^{\infty} \frac{e^{+ik \cdot \vec{r}}}{k^2 - \omega'^2/c^2} dk = -\frac{P_0}{(2\pi)^3 \epsilon_0} \frac{\partial}{\partial \vec{z}} \int_{k=0}^{\infty} \int_{\theta=0}^{\pi} \frac{e^{ikr \cos \theta}}{k^2 - \omega'^2/c^2} 2\pi k^2 \sin \theta dk d\theta \\ &= \frac{-iP_0}{4\pi^2 \epsilon_0} \frac{\partial}{\partial \vec{z}} \left\{ \frac{1}{r} \int_{k=0}^{\infty} \frac{k}{k^2 - \omega'^2/c^2} e^{ikr \cos \theta} \Big|_{\theta=0}^{\pi} dk \right\} = -\frac{2P_0}{4\pi^2 \epsilon_0} \frac{\partial}{\partial \vec{z}} \left\{ \frac{1}{r} \int_{k=0}^{\infty} \frac{k \sin(kr)}{k^2 - \omega'^2/c^2} dk \right\} \\ &= -\frac{P_0}{2\pi^2 \epsilon_0} \frac{\partial}{\partial \vec{z}} \left\{ \frac{1}{r} \int_0^{\infty} \frac{k \sin(kr)}{k^2 + (-i\omega/c)^2} dk \right\} = -\underbrace{\frac{P_0}{4\pi \epsilon_0} \frac{\partial}{\partial \vec{z}} \left( \frac{e^{i\omega r/c}}{r} \right)}_{\sim G.R. 3.723-3} \end{aligned}$$

$$\text{Therefore, } \psi(\vec{r}, t) = -\frac{P_0}{4\pi \epsilon_0} \frac{\partial}{\partial \vec{z}} \left[ \frac{e^{-i\omega t} e^{i\omega r/c} + e^{i\omega t} e^{-i\omega r/c}}{2r} \right] = -\frac{P_0}{4\pi \epsilon_0} \frac{\partial}{\partial \vec{z}} \left\{ \frac{1}{r} \cos[\omega(t - \frac{r}{c})] \right\}$$

$$= -\frac{P_0}{4\pi \epsilon_0} \left\{ -\frac{1}{r^2} \left( \frac{\partial r}{\partial \vec{z}} \right) \cos[\omega(t - \frac{r}{c})] + \frac{1}{r} \left( \frac{\omega}{c} \right) \left( \frac{\partial r}{\partial \vec{z}} \right) \sin[\omega(t - \frac{r}{c})] \right\}$$

$$= \frac{P_0}{4\pi \epsilon_0} \left\{ \frac{3}{r^3} \cos[\omega(t - \frac{r}{c})] - \left( \frac{\omega}{c} \right) \left( \frac{3}{r^2} \right) \sin[\omega(t - \frac{r}{c})] \right\} \Rightarrow$$

$$\boxed{\psi(\vec{r}, t) = \frac{P_0 \cos \theta}{4\pi \epsilon_0 r} \left\{ \frac{1}{r} \cos[\omega(t - \frac{r}{c})] - \left( \frac{\omega}{c} \right) \sin[\omega(t - \frac{r}{c})] \right\}}$$

$$\begin{aligned} r &= \sqrt{x^2 + y^2 + z^2} \\ \frac{\partial r}{\partial \vec{z}} &= \frac{\vec{z}}{r} = \cos \theta \end{aligned}$$

$$\text{Vector potential: Inverse Fourier Transform: } \vec{A}(\vec{r}) = -\frac{i\omega \mu_0 P_0 \hat{\vec{z}}}{(2\pi)^3} \int_{-\infty}^{\infty} \frac{e^{+ik \cdot \vec{r}}}{k^2 - \omega'^2/c^2} dk$$

$$= -\frac{i\omega \mu_0 P_0 \hat{\vec{z}}}{(2\pi)^3} \int_{k=0}^{\infty} \int_{\theta=0}^{\pi} \frac{e^{ikr \cos \theta}}{k^2 - \omega'^2/c^2} 2\pi k^2 \sin \theta dk d\theta = \frac{\omega \mu_0 P_0 \hat{\vec{z}}}{4\pi^2 r} \int_{k=0}^{\infty} \frac{k}{k^2 - \omega'^2/c^2} e^{ikr \cos \theta} \Big|_{\theta=0}^{\pi} dk$$

$$= -\frac{i\omega \mu_0 P_0 \hat{\vec{z}}}{2\pi^2 r} \int_{k=0}^{\infty} \frac{k \sin(kr)}{k^2 - \omega'^2/c^2} dk = -\underbrace{\frac{i\omega \mu_0 P_0 \hat{\vec{z}}}{4\pi r} e^{i\omega r/c}}$$

$$\text{Therefore, } \vec{A}(\vec{r}, t) = \frac{\mu_0 P_0 \hat{z}}{4\pi r} \left\{ -i\omega e^{i\omega r/c - i\omega t} + i\omega e^{-i\omega r/c} e^{i\omega t} \right\} \Rightarrow$$

$$\boxed{\vec{A}(\vec{r}, t) = -\frac{\mu_0 P_0 \omega \hat{z}}{4\pi r} \sin[\omega(t - \frac{r}{c})]}$$

$$\vec{B}(\vec{r}, t) = \vec{\nabla} \times \vec{A}(\vec{r}, t) = \vec{\nabla} \times \left\{ -\frac{\mu_0 P_0 \omega}{4\pi r} \sin[\omega(t - \frac{r}{c})] (\hat{r} \cos\theta - \hat{\theta} \sin\theta) \right\}$$

$$= -\frac{\mu_0 P_0 \omega}{4\pi} \left\{ \frac{1}{r} \frac{\partial}{\partial r} (-\sin[\omega(t - \frac{r}{c})] \sin\theta) - \frac{1}{r} \frac{\partial}{\partial \theta} \left( \frac{1}{r} \sin[\omega(t - \frac{r}{c})] \cos\theta \right) \right\} \hat{\phi}$$

$$= -\frac{\mu_0 P_0 \omega}{4\pi} \left\{ \frac{1}{r} \left( \frac{\omega}{c} \right) \cos[\omega(t - \frac{r}{c})] \sin\theta + \frac{1}{r^2} \sin[\omega(t - \frac{r}{c})] \sin\theta \right\} \hat{\phi} \Rightarrow$$

$$\boxed{\vec{H}(\vec{r}, t) = -\frac{P_0 \omega \sin\theta \hat{\phi}}{4\pi r} \left\{ \left( \frac{\omega}{c} \right) \cos[\omega(t - \frac{r}{c})] + \frac{1}{r} \sin[\omega(t - \frac{r}{c})] \right\}}$$

$$\vec{E}(\vec{r}, t) = -\vec{\nabla} \psi - \frac{\partial \vec{A}}{\partial t} = -\frac{\partial \psi}{\partial r} \hat{r} - \frac{1}{r} \frac{\partial \psi}{\partial \theta} \hat{\theta} + \frac{\mu_0 P_0 \omega^2}{4\pi r} \cos[\omega(t - \frac{r}{c})] (\cos\theta \hat{r} - \sin\theta \hat{\theta})$$

$$= -\frac{P_0 \cos\theta \hat{r}}{4\pi \epsilon_0} \left\{ \frac{-2}{r^3} \cos[\omega(t - \frac{r}{c})] + \frac{1}{r^2} \left( \frac{\omega}{c} \right) \sin[\omega(t - \frac{r}{c})] + \frac{1}{r^2} \sin[\omega(t - \frac{r}{c})] \left( \frac{\omega}{c} \right) \right. \\ \left. + \frac{1}{r} \left( \frac{\omega}{c} \right)^2 \cos[\omega(t - \frac{r}{c})] \right\} + \frac{P_0 \sin\theta \hat{\theta}}{4\pi \epsilon_0 r^2} \left\{ \frac{1}{r} \cos[\omega(t - \frac{r}{c})] - \left( \frac{\omega}{c} \right) \sin[\omega(t - \frac{r}{c})] \right\}$$

$$+ \frac{P_0}{4\pi \epsilon_0 r} \left( \frac{\omega}{c} \right)^2 (\cos\theta \hat{r} - \sin\theta \hat{\theta}) \cos[\omega(t - \frac{r}{c})] \Rightarrow$$

$$\vec{E}(\vec{r}, t) = \frac{P_0 \cos\theta \hat{r}}{4\pi \epsilon_0 r} \left\{ \frac{2}{r^2} \cos[\omega(t - \frac{r}{c})] - \frac{2}{r} \left( \frac{\omega}{c} \right) \sin[\omega(t - \frac{r}{c})] \right\}$$

$$+ \frac{P_0 \sin\theta \hat{\theta}}{4\pi \epsilon_0 r} \left\{ \left( \frac{1}{r^2} - \frac{\omega^2}{c^2} \right) \cos[\omega(t - \frac{r}{c})] - \frac{1}{r} \left( \frac{\omega}{c} \right) \sin[\omega(t - \frac{r}{c})] \right\} \Rightarrow$$

$$\boxed{\vec{E}(\vec{r}, t) = -\frac{P_0 \sin\theta \hat{\theta}}{4\pi \epsilon_0 r} \left( \frac{\omega}{c} \right)^2 \cos[\omega(t - \frac{r}{c})] - \frac{P_0}{4\pi \epsilon_0 r^2} \left\{ \left( \frac{\omega}{c} \right) \sin[\omega(t - \frac{r}{c})] - \frac{1}{r} \cos[\omega(t - \frac{r}{c})] \right\}}$$

$$\times (2 \cos\theta \hat{r} + \sin\theta \hat{\theta})$$

In the far field the terms that decrease as  $1/r^2$  and  $1/r^3$  may be ignored. The Poynting vector (in the far-field) is thus given by:

$$\vec{S}(\vec{r}, t) = \vec{E} \times \vec{H} = \frac{P_0 \sin\theta \hat{\theta}}{4\pi \epsilon_0 r} \left( \frac{\omega}{c} \right)^2 \cos[\omega(t - \frac{r}{c})] \times \frac{P_0 \omega \sin\theta \hat{\phi}}{4\pi r} \left( \frac{\omega}{c} \right) \cos[\omega(t - \frac{r}{c})] \\ = \frac{P_0^2 \sin^2\theta \omega^4}{16\pi^2 \epsilon_0^2 c^3 r^2} \cos^2[\omega(t - \frac{r}{c})] \hat{r} \Rightarrow \langle \vec{S}(\vec{r}, t) \rangle = \frac{\mu_0 P_0^2 \sin^2\theta \omega^4}{32\pi^2 c r^2} \hat{r}$$

The total radiated power is obtained by integrating over the surface of an arbitrary sphere of radius  $r$ . We find:

$$\begin{aligned}\text{Total radiated power} &= \int_{\theta=0}^{\pi} 2\pi r^2 \sin \theta \langle \vec{s}(\vec{r}, t) \rangle d\theta \\ &= \frac{M_0 P_0^2 \omega^4}{16\pi c} \int_0^{\pi} \sin^3 \theta d\theta = \underbrace{\frac{M_0 P_0^2 \omega^4}{12\pi c}}\end{aligned}$$

Digression: The scalar potential  $\Phi(\vec{r})$  may also be obtained by direct integration, as follows:

$$\begin{aligned}\int_{-\infty}^{\infty} \frac{k_3}{k^2 - \omega_c^2} e^{ik \cdot \vec{r}} dk &= \int_{k=0}^{\infty} \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} \frac{k \cos \theta}{k^2 - \omega_c^2} e^{ik(k \sin \theta \cos \phi \hat{x} + k \sin \theta \sin \phi \hat{y} + k \cos \theta \hat{z})} \\ &\quad k^2 \sin \theta dkd\theta d\phi \\ &= \int_{k=0}^{\infty} \int_{\theta=0}^{\pi} \frac{k^3 \sin \theta \cos \theta e^{ik_3 \sin \theta}}{k^2 - \omega_c^2} \left\{ \int_{\phi=0}^{2\pi} e^{ik \sin \theta (x \sin \phi + y \cos \phi)} d\phi \right\} dk d\theta \quad \nwarrow \text{G.R. 3.437-2}\end{aligned}$$

$$= 2\pi \int_{k=0}^{\infty} \frac{k^3}{k^2 - \omega_c^2} \left\{ \int_{\theta=0}^{\pi} e^{ik_3 \sin \theta} J_0(k \sqrt{x^2 + y^2} \sin \theta) \sin \theta \cos \theta d\theta \right\} dk$$

$$= 2\pi \int_{k=0}^{\infty} \frac{k^3}{k^2 - \omega_c^2} \left\{ 2i \int_0^1 \sin(k_3 \sqrt{1-u^2}) J_0(k \sqrt{x^2 + y^2} u) u du \right\} dk$$

$$= \frac{4\pi i}{x^2 + y^2} \int_{k=0}^{\infty} \frac{k}{k^2 - \omega_c^2} \left\{ \int_0^{k \sqrt{x^2 + y^2}} \sin(k_3 \sqrt{1 - \frac{v^2}{k^2(x^2 + y^2)}}) J_0(v) v dv \right\} dk \quad \nwarrow \text{G.R. 6.738-1}$$

$$= \frac{4\pi i}{x^2 + y^2} \int_{k=0}^{\infty} \frac{k}{k^2 - \omega_c^2} \left\{ \frac{\sqrt{\pi}}{2} (k \sqrt{x^2 + y^2})^{3/2} \left( \frac{3}{\sqrt{x^2 + y^2}} \right) \left( 1 + \frac{3^2}{x^2 + y^2} \right)^{-3/4} J_{3/2}(k \sqrt{x^2 + y^2 + z^2}) \right\} dk$$

$$= 4\pi i \left( \frac{3}{r^{3/2}} \right) \sqrt{\frac{\pi}{2}} \int_0^{\infty} \frac{k^{5/2}}{k^2 - \omega_c^2} J_0(kr) dk; \quad \leftarrow J_{3/2}(r) = \sqrt{\frac{2}{\pi r}} \left( \frac{\sin r}{r} - \cos r \right); \text{ G.R. 8.464-3$$

$\nwarrow \text{G.R. 6.565-4} \quad [\text{write } k^2 - \omega_c^2 = k^2 + (-i\omega/c)^2; \text{ Re}(-i\omega/c) = \omega_c^2 > 0]$

$$= 4\pi i \left( \frac{3}{r^{3/2}} \right) \sqrt{\frac{\pi}{2}} (-i\omega/c)^{3/2} K_{3/2}(-i \frac{\omega}{c} r); \quad \leftarrow K_{3/2}(r) = \sqrt{\frac{\pi}{2r}} e^{-r} (1 + \frac{1}{r}); \text{ G.R. 8.468}$$

$$\Rightarrow \int_{-\infty}^{\infty} \frac{k_3}{k^2 - \omega_c^2} e^{i\vec{k} \cdot \vec{r}} dk = 4\pi i \left(\frac{3}{r^{3/2}}\right) \sqrt{\pi/2} \left(-i\frac{\omega}{c}\right)^{3/2} \sqrt{\frac{\pi}{-2i\frac{\omega}{c}r}} e^{i\frac{\omega}{c}r} \left(1 + \frac{1}{-i\omega/c r}\right)$$

$$= 2\pi^2 \left(\frac{\omega}{c}\right) \left(\frac{3}{r^2}\right) \left(1 + \frac{i}{(\omega/c)r}\right) e^{i(\omega/c)r}$$

Having completed the integration, we may now assume that  $\omega$  is real-valued (i.e., set  $\omega'' = 0$ ). We'll have:

$$\psi(\vec{r}, t) = -\frac{iP_0}{2\epsilon_0} \frac{1}{(2\pi)^3} \left\{ 2\pi^2 \left(\frac{\omega}{c}\right) \left(\frac{3}{r^2}\right) \left(1 + \frac{i}{(\omega/c)r}\right) e^{i(\omega/c)r - i\omega t} e^{-i(\omega/c)r + i\omega t} + 2\pi^2 \left(-\frac{\omega}{c}\right) \left(\frac{3}{r^2}\right) \left(1 - \frac{i}{(\omega/c)r}\right) e^{-i(\omega/c)r} e^{+i\omega t} \right\}$$

$$= -i \frac{P_0}{8\pi\epsilon_0} \left(\frac{\omega}{c}\right) \left(\frac{3}{r^2}\right) \left\{ \left(1 + \frac{i}{(\omega/c)r}\right) e^{-i\omega t - r/c} - \left(1 - \frac{i}{(\omega/c)r}\right) e^{+i\omega(t - r/c)} \right\}$$

$\approx 3/r$

$$= \frac{P_0 \cos \theta}{4\pi\epsilon_0 r} \left\{ \frac{1}{r} \cos [\omega t - r/c] - \left(\frac{\omega}{c}\right) \sin [\omega(t - r/c)] \right\}$$