

**Problem 4.12)** We begin by computing the Fourier transform of the uniformly-polarized spherical particle, as follows:

$$\begin{aligned}
\mathbf{P}(\mathbf{k}) &= \iiint_{-\infty}^{\infty} \mathbf{P}(\mathbf{r}) \exp(-i\mathbf{k} \cdot \mathbf{r}) \, d\mathbf{r} = \int_{r=0}^R \int_{\theta=0}^{\pi} P_0 \hat{\mathbf{z}} \exp(-ikr \cos \theta) 2\pi r^2 \sin \theta \, dr d\theta \\
&= 2\pi P_0 \hat{\mathbf{z}} \int_{r=0}^R r^2 \left[ \int_{\theta=0}^{\pi} \sin \theta \exp(-ikr \cos \theta) \, d\theta \right] dr \\
&= 2\pi P_0 \hat{\mathbf{z}} \int_0^R r^2 \left[ \frac{\exp(-ikr) - \exp(ikr)}{-ikr} \right] dr = (4\pi P_0 \hat{\mathbf{z}}/k) \int_0^R r \sin(kr) \, dr \quad \leftarrow \text{Change of variable: } x = kr \\
&\xrightarrow{\text{Integration by parts}} = (4\pi P_0 \hat{\mathbf{z}}/k^3) \int_0^{kR} x \sin(x) \, dx = (4\pi P_0 \hat{\mathbf{z}}/k^3) \left[ -x \cos x \Big|_0^{kR} + \int_0^{kR} \cos x \, dx \right] \\
&= 4\pi P_0 \hat{\mathbf{z}} [\sin(kR) - kR \cos(kR)]/k^3. \tag{1}
\end{aligned}$$

Given that  $\rho_{\text{bound}}^{(e)}(\mathbf{r}, t) = -\nabla \cdot \mathbf{P}(\mathbf{r}, t)$ , in the Fourier domain we will have

$$\rho_{\text{bound}}^{(e)}(\mathbf{k}) = -i\mathbf{k} \cdot \mathbf{P}(\mathbf{k}) = -i4\pi P_0 (\hat{\mathbf{z}} \cdot \hat{\mathbf{k}}) [\sin(kR) - kR \cos(kR)]/k^2. \tag{2}$$

Now, the scalar potential in the Fourier domain is known to be

$$\psi(\mathbf{k}) = \rho_{\text{total}}^{(e)}(\mathbf{k})/(\varepsilon_0 k^2). \tag{3}$$

Therefore, inverse Fourier transformation of  $\psi(\mathbf{k})$  yields

$$\begin{aligned}
\psi(\mathbf{r}) &= -\frac{i4\pi P_0}{(2\pi)^3 \varepsilon_0} \hat{\mathbf{z}} \cdot \int_{-\infty}^{\infty} \hat{\mathbf{k}} \left[ \frac{\sin(kR) - kR \cos(kR)}{k^4} \right] \exp(i\mathbf{k} \cdot \mathbf{r}) \, d\mathbf{k} \\
&= -\frac{iP_0}{2\pi^2 \varepsilon_0} (\hat{\mathbf{z}} \cdot \hat{\mathbf{r}}) \int_{k=0}^{\infty} \int_{\theta=0}^{\pi} \cos \theta \left[ \frac{\sin(kR) - kR \cos(kR)}{k^4} \right] \exp(ikr \cos \theta) 2\pi k^2 \sin \theta \, d\theta dk \\
&= -\frac{iP_0}{\pi \varepsilon_0} (\hat{\mathbf{z}} \cdot \hat{\mathbf{r}}) \int_{k=0}^{\infty} \left[ \frac{\sin(kR) - kR \cos(kR)}{k^2} \right] \int_{\theta=0}^{\pi} \sin \theta \cos \theta \exp(ikr \cos \theta) \, d\theta dk \\
&= \frac{2P_0}{\pi \varepsilon_0 r^2} (\hat{\mathbf{z}} \cdot \hat{\mathbf{r}}) \int_0^{\infty} \frac{\sin(kR) - kR \cos(kR)}{k^2} \times \frac{\sin(kr) - kr \cos(kr)}{k^2} \, dk \quad \leftarrow \text{G\&R 3.715-11} \\
&= \frac{2P_0}{\pi \varepsilon_0 r^2} (\hat{\mathbf{z}} \cdot \hat{\mathbf{r}}) \int_0^{\infty} \frac{d}{dk} [\sin(kR)/k] \times \frac{d}{dk} [\sin(kr)/k] \, dk. \tag{4}
\end{aligned}$$

To evaluate the integral appearing in the above equation, note that  $\sin(\alpha k)/k$  is the Fourier transform of  $\frac{1}{2}\text{Rect}(x/2\alpha)$ , which implies that  $d[\sin(\alpha k)/k]/dk$  is the Fourier transform of  $-(ix/2)\text{Rect}(x/2\alpha)$ . It is not difficult to show that  $\int_{-\infty}^{\infty} F(k)G(k)dk = 2\pi \int_{-\infty}^{\infty} f(x)g(-x)dx$  for any pair of functions  $f(x)$  and  $g(x)$  whose (one-dimensional) Fourier transforms are specified as  $F(k)$  and  $G(k)$ . Finally, taking note of the fact that the integrand in Eq.(4) is an even function of  $k$ , we write

$$\begin{aligned}
\int_0^{\infty} \frac{d}{dk} [\sin(kR)/k] \times \frac{d}{dk} [\sin(kr)/k] \, dk &= \pi \int_{-\infty}^{\infty} (-ix/2)\text{Rect}(x/2R) \times (ix/2)\text{Rect}(-x/2r) \, dx \\
&= (\pi/4) \int_{-\min(r,R)}^{\min(r,R)} x^2 \, dx = (\pi/6) \min(r^3, R^3). \tag{5}
\end{aligned}$$

Upon substitution from Eq.(5) into Eq.(4), and also replacing  $\hat{\mathbf{z}} \cdot \hat{\mathbf{r}}$  with  $\cos \theta$ , we find

$$\psi(\mathbf{r}) = \begin{cases} P_0 r \cos \theta / (3\epsilon_0); & r \leq R, \\ P_0 R^3 \cos \theta / (3\epsilon_0 r^2); & r \geq R. \end{cases} \quad (6)$$

Inside the spherical particle,  $\psi(\mathbf{r}) = (P_0/3\epsilon_0)z$ . Therefore,

$$\mathbf{E}(\mathbf{r}) = -\nabla\psi(\mathbf{r}) = -(\partial\psi/\partial z)\hat{\mathbf{z}} = -(P_0/3\epsilon_0)\hat{\mathbf{z}}. \quad (7)$$

Outside the particle,

$$\mathbf{E}(\mathbf{r}) = -\nabla\psi(\mathbf{r}) = -\frac{\partial\psi}{\partial r}\hat{\mathbf{r}} - \frac{\partial\psi}{r\partial\theta}\hat{\boldsymbol{\theta}} = \frac{(4\pi/3)R^3P_0}{4\pi\epsilon_0r^3}(2\cos\theta\hat{\mathbf{r}} + \sin\theta\hat{\boldsymbol{\theta}}). \quad (8)$$

Thus, from outside the sphere, it appears as though a point-dipole  $p_0\hat{\mathbf{z}} = (4\pi/3)R^3P_0\hat{\mathbf{z}}$  is residing at the origin of the coordinate system. Inside the particle, as can be seen from Eq.(7), the  $E$ -field is uniform and oriented opposite to the direction of polarization.

**Digression:** We prove below that  $\int_{-\infty}^{\infty} F(k)G(k)dk = 2\pi \int_{-\infty}^{\infty} f(x)g(-x)dx$  for any pair of functions  $f(x)$  and  $g(x)$ , whose (one-dimensional) Fourier transforms are specified as  $F(k)$  and  $G(k)$ .

$$\begin{aligned} \int_{-\infty}^{\infty} F(k)G(k)dk &= \int_{-\infty}^{\infty} \left[ \int_{-\infty}^{\infty} f(x) \exp(-ikx) dx \right] G(k)dk \\ &= \int_{-\infty}^{\infty} f(x) \left[ \int_{-\infty}^{\infty} G(k) \exp(-ikx) dk \right] dx = 2\pi \int_{-\infty}^{\infty} f(x)g(-x)dx. \end{aligned} \quad (9)$$

Also, it is easy to show that, if  $F(k)$  is the Fourier transform of  $f(x)$ , then  $dF(k)/dk$  is the Fourier transform of  $-ixf(x)$ . This is because differentiating  $F(k) = \int_{-\infty}^{\infty} f(x) \exp(-ikx) dx$  with respect to  $k$  yields

$$dF(k)/dk = \int_{-\infty}^{\infty} f(x)(-ix) \exp(-ikx) dx = \mathcal{F}\{-ixf(x)\}. \quad (10)$$

We have used both Eq.(9) and Eq.(10) in evaluating the definite integral given by Eq.(5).

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