

Problem 10)

$$\vec{J}(\vec{r}, t) = I_0 \hat{z} \delta(x) \delta(y) \text{Rect}(z/L)$$

$$\begin{aligned} \vec{J}(\vec{k}, \omega) &= \int_{-\infty}^{\infty} \vec{J}(\vec{r}, t) e^{-i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{r} dt = I_0 \hat{z} \int_{-\infty}^{\infty} \delta(x) e^{-ik_x x} dx \int_{-\infty}^{\infty} \delta(y) e^{-ik_y y} dy \\ &\quad \times \int_{-\infty}^{\infty} \text{Rect}(z/L) e^{-ik_z z} dz \int_{-\infty}^{\infty} e^{i\omega t} dt = I_0 \hat{z} \frac{e^{-ik_z L/2} - e^{+ik_z L/2}}{-ik_z} 2\pi \delta(\omega) \Rightarrow \end{aligned}$$

$$\vec{J}(\vec{k}, \omega) = 4\pi I_0 \hat{z} \frac{\sin(Lk_z/2)}{k_z} \delta(\omega)$$

$$\begin{aligned} \vec{A}(\vec{r}, t) &= \frac{1}{(2\pi)^4} \int_{-\infty}^{\infty} \frac{\mu_0 \vec{J}(\vec{k}, \omega)}{k^2 - (\omega/c)^2} e^{+i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{k} d\omega \\ &= \frac{4\pi \mu_0 I_0 \hat{z}}{(2\pi)^4} \int_{-\infty}^{\infty} \frac{\sin(Lk_z/2)}{k_z (k^2 - \frac{\omega^2}{c^2})} \delta(\omega) e^{+i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{k} d\omega \\ &= \frac{\mu_0 I_0 \hat{z}}{4\pi^3} \int_{-\infty}^{\infty} \frac{\sin(Lk_z/2)}{k_z (k_x^2 + k_y^2 + k_z^2)} e^{i(k_x x + k_y y + k_z z)} dk_x dk_y dk_z \end{aligned}$$

Let $\vec{k}_{||} = k_x \hat{x} + k_y \hat{y}$ and $\vec{r}_{||} = x \hat{x} + y \hat{y}$. We'll have:

$$\vec{A}(\vec{r}, t) = \frac{\mu_0 I_0 \hat{z}}{4\pi^3} \int_{-\infty}^{\infty} \frac{\sin(Lk_z/2)}{k_z} e^{ik_z z} \left\{ \int_{k_{||}=0}^{\infty} \int_{\theta=0}^{2\pi} \frac{e^{i\vec{k}_{||} \cdot \vec{r}_{||} \cos\theta}}{k_{||}^2 + k_z^2} k_{||} dk_{||} d\theta \right\} dk_z$$

G.R. 3.915-2

$$\Rightarrow = \frac{\mu_0 I_0 \hat{z}}{4\pi^3} \int_{-\infty}^{\infty} \frac{\sin(Lk_z/2)}{k_z} e^{ik_z z} \left\{ \int_{k_{||}=0}^{\infty} \frac{2\pi k_{||} J_0(k_{||} r_{||})}{k_{||}^2 + k_z^2} dk_{||} \right\} dk_z$$

G.R. 6.532-4

$$\Rightarrow = \frac{\mu_0 I_0 \hat{z}}{2\pi^2} \int_{-\infty}^{\infty} \frac{\sin(Lk_z/2)}{k_z} e^{ik_z z} K_0(|k_z| r_{||}) dk_z$$

$$\begin{aligned} &= \frac{\mu_0 I_0 \hat{z}}{2\pi^2} \left\{ \int_{-\infty}^0 \frac{\sin(Lk_z/2)}{k_z} e^{ik_z z} K_0(|k_z| r_{||}) dk_z + \int_0^{\infty} \frac{\sin(Lk_z/2)}{k_z} e^{ik_z z} K_0(|k_z| r_{||}) dk_z \right\} \\ &= \frac{\mu_0 I_0 \hat{z}}{\pi^2} \int_0^{\infty} \frac{\sin(Lk_z/2) \cos(k_z z)}{k_z} K_0(k_z r_{||}) dk_z \\ &= \frac{\mu_0 I_0 \hat{z}}{2\pi^2} \int_0^{\infty} \frac{\sin[(\frac{L}{2} + z)k_z] + \sin[(\frac{L}{2} - z)k_z]}{k_z} K_0(k_z r_{||}) dk_z \end{aligned}$$

Using G.R. 6.699-3, 9.121-28, and 8.338-2, we'll have

$$\begin{aligned}\vec{A}(\vec{r}, t) &= \frac{\mu_0 I_0 \hat{z}}{4\pi} \left\{ \ln \frac{|\frac{L}{2} + z| + \sqrt{(\frac{L}{2} + z)^2 + r_{||}^2}}{r_{||}} + \ln \frac{|\frac{L}{2} - z| + \sqrt{(\frac{L}{2} - z)^2 + r_{||}^2}}{r_{||}} \right\} \\ &= \frac{\mu_0 I_0 \hat{z}}{4\pi} \left\{ \ln \left[\left(|\frac{L}{2} + z| + \sqrt{(\frac{L}{2} + z)^2 + r_{||}^2} \right) \left(|\frac{L}{2} - z| + \sqrt{(\frac{L}{2} - z)^2 + r_{||}^2} \right) \right] - 2 \ln r_{||} \right\}\end{aligned}$$

In the limit $L \rightarrow \infty$, the logarithmic term becomes infinitely large.

However, its dependence on $r_{||}$ and z diminishes to the point that any variations of this term with $r_{||}$ and z may be ignored. In other words, when computing $\vec{\nabla} \times \vec{A}$, the derivatives with respect to $r_{||}$ and z of the logarithmic term approach zero when $L \rightarrow \infty$. We thus ignore this logarithmic term; even though its magnitude is very large, it effectively behaves like a constant term. We thus have

$$\vec{A}(\vec{r}, t) = - \frac{\mu_0 I_0 \hat{z}}{2\pi} \ln r_{||} + \text{Constant}_{\infty} \leftarrow (\text{when } L \rightarrow \infty)$$

$$\vec{B}(\vec{r}, t) = \vec{\nabla} \times \vec{A}(\vec{r}, t) \Rightarrow \mu_0 \vec{H}(\vec{r}, t) = - \frac{\partial A_z}{\partial r_{||}} \hat{\phi} = \frac{\mu_0 I_0}{2\pi r_{||}} \hat{\phi} \Rightarrow \vec{H}(\vec{r}, t) = \frac{I_0 \hat{\phi}}{2\pi r_{||}}$$