

**Problem 2.61)**

$$\begin{aligned}
 \text{a)} \quad \nabla \cdot \mathbf{D}(\mathbf{r}, t) &= \rho_{\text{free}}(\mathbf{r}, t), \\
 \nabla \times \mathbf{H}(\mathbf{r}, t) &= \mathbf{J}_{\text{free}}(\mathbf{r}, t) + \frac{\partial \mathbf{D}(\mathbf{r}, t)}{\partial t}, \\
 \nabla \times \mathbf{E}(\mathbf{r}, t) &= -\frac{\partial \mathbf{B}(\mathbf{r}, t)}{\partial t}, \\
 \nabla \cdot \mathbf{B}(\mathbf{r}, t) &= 0.
 \end{aligned}$$

In the above equations,  $\mathbf{r} = x\hat{\mathbf{x}} + y\hat{\mathbf{y}} + z\hat{\mathbf{z}}$  is an arbitrary point in space, while  $t$  is an arbitrary instant in time.  $\mathbf{E}$  is the electric field,  $\mathbf{H}$  is the magnetic field,  $\mathbf{D}$  is the displacement, and  $\mathbf{B}$  is the magnetic induction. The fields are related to each other, to the permittivity and permeability of free space,  $\epsilon_0$  and  $\mu_0$ , and to polarization  $\mathbf{P}$  and magnetization  $\mathbf{M}$  as follows

$$\begin{aligned}
 \mathbf{D}(\mathbf{r}, t) &= \epsilon_0 \mathbf{E}(\mathbf{r}, t) + \mathbf{P}(\mathbf{r}, t), \\
 \mathbf{B}(\mathbf{r}, t) &= \mu_0 \mathbf{H}(\mathbf{r}, t) + \mathbf{M}(\mathbf{r}, t).
 \end{aligned}$$

The sources of the electromagnetic fields (namely,  $\mathbf{E}$  and  $\mathbf{H}$ ) are the free charge density  $\rho_{\text{free}}$ , free current density  $\mathbf{J}_{\text{free}}$ , polarization  $\mathbf{P}$  (which is the density of electric dipole moments), and magnetization  $\mathbf{M}$  (which is the density of magnetic dipole moments). The operator  $\partial/\partial t$  represents partial differentiation with respect to time,  $\nabla \cdot$  is the divergence operator, and  $\nabla \times$  is the curl operator. The divergence of a vector field such as  $\mathbf{D}(\mathbf{r}, t)$ , which turns out to be a scalar field, is defined as the integral of  $\mathbf{D}(\mathbf{r}, t)$  over a small closed surface, normalized by the enclosed volume. The curl of a vector field such as  $\mathbf{E}(\mathbf{r}, t)$ , which turns out to be another vector field, when projected onto the surface normal of a small surface element, yields the line integral of  $\mathbf{E}(\mathbf{r}, t)$  around the boundary of the small surface element, normalized by the surface area of the element.

b) To derive the charge-current continuity equation from Maxwell's equations, apply the divergence operator to both sides of the second (Maxwell-Ampere) equation. The divergence of curl is always equal to zero and, therefore, the left-hand-side of the equation becomes  $\nabla \cdot (\nabla \times \mathbf{H}) = 0$ . The right-hand side,  $\nabla \cdot \mathbf{J}_{\text{free}} + \partial(\nabla \cdot \mathbf{D})/\partial t$ , thus becomes zero. Maxwell's first equation (Gauss's law) now allows one to replace  $\nabla \cdot \mathbf{D}$  with  $\rho_{\text{free}}$ , yielding the continuity equation as  $\nabla \cdot \mathbf{J}_{\text{free}} + \partial \rho_{\text{free}}/\partial t = 0$ . This equation informs that the integrated free current over any closed surface is precisely balanced by changes in the electrical charge contained within the closed surface. If there is a net outflow of the current, the charge within the closed surface must be decreasing, and if there is a net inflow of current, the charge within must be increasing.

c) In the first of Maxwell's equations, we substitute  $\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P}$  and obtain

$$\nabla \cdot (\epsilon_0 \mathbf{E} + \mathbf{P}) = \rho_{\text{free}} \quad \rightarrow \quad \epsilon_0 \nabla \cdot \mathbf{E} = \rho_{\text{free}} - \nabla \cdot \mathbf{P} \quad \rightarrow \quad \epsilon_0 \nabla \cdot \mathbf{E} = \rho_{\text{free}} + \rho_{\text{bound}}^{(e)}.$$

The bound-charge density is thus seen to be  $\rho_{\text{bound}}^{(e)}(\mathbf{r}, t) = -\nabla \cdot \mathbf{P}(\mathbf{r}, t)$ .

In the second Maxwell equation, we multiply both sides by  $\mu_0$ , then add  $\nabla \times \mathbf{M}$  to both sides, in order to replace  $\mathbf{H}$  with  $\mathbf{B}$  through the identity  $\mathbf{B} = \mu_0 \mathbf{H} + \mathbf{M}$ . We also use  $\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P}$  on the right-hand side of the equation to get rid of  $\mathbf{D}$ . We will have

$$\begin{aligned} \mu_0 \nabla \times \mathbf{H} + \nabla \times \mathbf{M} &= \mu_0 \mathbf{J}_{\text{free}} + \mu_0 \frac{\partial(\epsilon_0 \mathbf{E} + \mathbf{P})}{\partial t} + \nabla \times \mathbf{M} \\ \rightarrow \quad \nabla \times \mathbf{B} &= \mu_0 (\mathbf{J}_{\text{free}} + \partial \mathbf{P} / \partial t + \mu_0^{-1} \nabla \times \mathbf{M}) + \mu_0 \epsilon_0 \partial \mathbf{E} / \partial t \\ \rightarrow \quad \nabla \times \mathbf{B} &= \mu_0 (\mathbf{J}_{\text{free}} + \mathbf{J}_{\text{bound}}^{(e)}) + \mu_0 \epsilon_0 \partial \mathbf{E} / \partial t. \end{aligned}$$

The bound electric current density is thus found to be  $\mathbf{J}_{\text{bound}}^{(e)} = \partial \mathbf{P} / \partial t + \mu_0^{-1} \nabla \times \mathbf{M}$ . Since the remaining Maxwell equations do not contain  $\mathbf{D}$  and  $\mathbf{H}$ , they remain unchanged.

d) The divergence of  $\mathbf{J}_{\text{bound}}^{(e)}$  is readily obtained as follows:

$$\nabla \cdot \mathbf{J}_{\text{bound}}^{(e)} = \partial(\nabla \cdot \mathbf{P}) / \partial t + \mu_0^{-1} \nabla \cdot (\nabla \times \mathbf{M}).$$

On the right-hand side of the above equation, the divergence of curl is always zero. Also the divergence of  $\mathbf{P}(\mathbf{r}, t)$  is, by definition,  $-\rho_{\text{bound}}^{(e)}$ . Therefore,  $\nabla \cdot \mathbf{J}_{\text{bound}}^{(e)} + \partial \rho_{\text{bound}}^{(e)} / \partial t = 0$ . This is the charge-current continuity equation for the bound electrical charge and current defined in part (c).

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