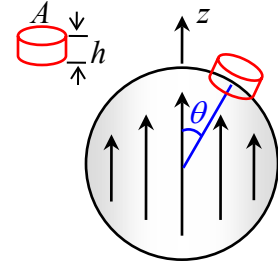


**Problem 2.42)** From Maxwell's 4<sup>th</sup> equation, we find the bound magnetic charge-density to be given by  $\rho_{\text{bound}}^{(m)} = -\nabla \cdot \mathbf{M}(\mathbf{r})$ . Take a small pillbox and place it anywhere inside the sphere. The magnetization entering the box will be equal to that leaving the box and, therefore, the divergence of  $\mathbf{M}(\mathbf{r}) = M_0 \hat{\mathbf{z}}$  will be zero everywhere inside the sphere. The only points where the divergence will be non-zero are at the surface of the sphere. The figure shows a small, thin pillbox placed at  $(r=R, \theta, \phi)$ . Let  $A$  and  $h$  denote the base area and height of the pillbox, respectively; both  $A$  and  $h$  could be as small as desired. The flux of  $\mathbf{M}$  entering from the bottom of the pillbox is  $M_0 A \cos \theta$ , and this is the only contribution to the integral of  $\mathbf{M}(\mathbf{r})$  over the pillbox surface, provided that  $h$  is much smaller than the pillbox diameter. The divergence of  $\mathbf{M}$  at  $(r=R, \theta, \phi)$  is thus given by  $-M_0 A \cos \theta / (Ah)$  in the limit of small  $A$  and  $h$ . Therefore,  $\rho_{\text{bound}}^{(m)}(R, \theta, \phi) = M_0 \cos \theta / h$ . Since the charges are confined to the surface, we should use the surface-charge-density  $\sigma_{\text{bound}}^{(m)} = h \rho_{\text{bound}}^{(m)}$  instead of the volume charge-density. Consequently,  $\sigma_{\text{bound}}^{(m)}(R, \theta, \phi) = M_0 \cos \theta$ .



To determine the bound electric current-density  $\mathbf{J}_{\text{bound}}^{(e)} = \mu_0^{-1} \nabla \times \mathbf{M}(\mathbf{r})$ , we use a small rectangular loop (length =  $\ell$ , width =  $w$ ) at various locations within and on the surface of the sphere in order to calculate the curl of  $\mathbf{M}$ . For all locations within the sphere and for all orientations of the loop, the integral of  $\mathbf{M}(\mathbf{r})$  around the loop turns out to be zero. When the loop is placed on the surface at  $(r=R, \theta, \phi)$  and oriented perpendicular to  $\hat{\phi}$ , as shown, the line integral on the lower leg of the loop will be nonzero ( $\ell M_0 \sin \theta$ ). The curl of  $\mathbf{M}(\mathbf{r})$  will then be nonzero, as the other legs do not contribute to the integral, provided that  $w \ll \ell$ . The curl will then be given by  $[\ell M_0 \sin \theta / (\ell w)] \hat{\phi}$  in the limit when  $\ell$  and  $w$  both tend to zero. The bound current-density at  $(r=R, \theta, \phi)$  is thus given by  $\mathbf{J}_{\text{bound}}^{(e)} = \mu_0^{-1} (M_0 \sin \theta / w) \hat{\phi}$ . Since the current is confined to a thin layer on the surface, we could use the surface-current-density  $\mathbf{J}_{\text{s-bound}}^{(e)} = w \mathbf{J}_{\text{bound}}^{(e)}$  instead of the bulk current-density. Consequently,  $\mathbf{J}_{\text{s-bound}}^{(e)} = \mu_0^{-1} M_0 \sin \theta \hat{\phi}$ .

