

Problem 21)

$$a) \quad \mathbf{J}(\mathbf{r}, t) = qV(t) \delta(x) \delta(y) \delta[z - \int_0^t V(t') dt'] \hat{\mathbf{z}}.$$

$$b) \quad \partial \mathcal{E}(\mathbf{r}, t) / \partial t = \mathbf{E}(\mathbf{r}, t) \cdot \mathbf{J}(\mathbf{r}, t) = E_0 q V(t) \delta(x) \delta(y) \delta[z - \int_0^t V(t') dt'].$$

Integrating over the energy density, let $\mathcal{E}(t) = \iiint_{-\infty}^{\infty} \mathcal{E}(\mathbf{r}, t) dx dy dz$ denote the energy of the particle at time t . We will have

$$d\mathcal{E}(t)/dt = (d/dt) \iiint_{-\infty}^{\infty} \mathcal{E}(\mathbf{r}, t) dx dy dz = \iiint_{-\infty}^{\infty} [\partial \mathcal{E}(\mathbf{r}, t) / \partial t] dx dy dz = qE_0 V(t).$$

The change in the kinetic energy of the particle as it moves from $z=0$ to $z=d$ is thus given by

$$\Delta \mathcal{E} = \int_{t_0}^{t_1} qE_0 V(t) dt = qE_0 \int_{t_0}^{t_1} V(t) dt = qE_0 d.$$

c) The E -field of the system is the sum of the constant E -field between the plates and that of the point particle, namely, $\mathbf{E}(\mathbf{r}, t) = E_0 \text{Rect}(z/d) \hat{\mathbf{z}} + (q/4\pi\epsilon_0) [\mathbf{r} - \mathbf{r}_o(t)] / |\mathbf{r} - \mathbf{r}_o(t)|^3$. The E -field energy density is therefore given by

$$\begin{aligned} \mathcal{E}(\mathbf{r}, t) &= \frac{1}{2} \epsilon_0 \mathbf{E}^2(\mathbf{r}, t) = \frac{1}{2} \epsilon_0 \mathbf{E}(\mathbf{r}, t) \cdot \mathbf{E}(\mathbf{r}, t) \\ &= \frac{1}{2} \epsilon_0 \{ E_0^2 \text{Rect}(z/d) + (q/4\pi\epsilon_0)^2 / |\mathbf{r} - \mathbf{r}_o(t)|^4 + 2(qE_0/4\pi\epsilon_0) \text{Rect}(z/d) [\mathbf{r} - \mathbf{r}_o(t)] \cdot \hat{\mathbf{z}} / |\mathbf{r} - \mathbf{r}_o(t)|^3 \} \\ &= \frac{1}{2} \epsilon_0 E_0^2 \text{Rect}(z/d) + (q^2/32\pi^2 \epsilon_0) \{ x^2 + y^2 + [z - z_o(t)]^2 \}^{-2} \\ &\quad + (qE_0/4\pi) \text{Rect}(z/d) [z - z_o(t)] / \{ x^2 + y^2 + [z - z_o(t)]^2 \}^{-3/2}. \end{aligned}$$

The total E -field energy of the system at time t is obtained by integrating the above energy density over the entire space, that is,

$$\begin{aligned} \mathcal{E}(t) &= \iiint_{-\infty}^{\infty} \mathcal{E}(\mathbf{r}, t) dx dy dz = \frac{1}{2} \epsilon_0 E_0^2 \iiint_{-\infty}^{\infty} \text{Rect}(z/d) dx dy dz \\ &\quad + (q^2/32\pi^2 \epsilon_0) \iiint_{-\infty}^{\infty} (x^2 + y^2 + [z - z_o(t)]^2)^{-2} dx dy dz \\ &\quad + (qE_0/4\pi) \iiint_{-\infty}^{\infty} \text{Rect}(z/d) [z - z_o(t)] (x^2 + y^2 + [z - z_o(t)]^2)^{-3/2} dx dy dz \\ &= \frac{1}{2} \epsilon_0 E_0^2 Ad + (q^2/32\pi^2 \epsilon_0) \iiint_{-\infty}^{\infty} (x^2 + y^2 + z^2)^{-2} dx dy dz \quad \leftarrow \text{self energy of point charge} = \infty \\ &\quad + (qE_0/4\pi) \int_{-\infty}^{\infty} \text{Rect}(z/d) [z - z_o(t)] dz \int_0^{\infty} (r^2 + [z - z_o(t)]^2)^{-3/2} 2\pi r dr. \end{aligned}$$

The first term in the above expression is the E -field energy of the plates in the absence of the point particle, which is a constant, independent of the position of the particle in the system. The second term, although infinite in magnitude – because of the zero diameter of the point particle – is constant nonetheless, as it does not depend on the particle's position within the system. The only relevant term, therefore, is the last term, which may be evaluated as follows:

$$\begin{aligned}
\mathcal{E}(t) &= \text{constant} + \frac{1}{2}qE_0 \int_{-\infty}^{\infty} \text{Rect}(z/d)[z-z_0(t)] \left\{ -(r^2 + [z-z_0(t)]^2)^{-1/2} \right\} \Big|_0^{\infty} dz \\
&= \text{constant} + \frac{1}{2}qE_0 \int_{-\infty}^{\infty} \text{Rect}(z/d)[z-z_0(t)] |z-z_0(t)|^{-1} dz \\
&= \text{constant} + \frac{1}{2}qE_0 [d - 2z_0(t)].
\end{aligned}$$

Thus, as the particle moves from $z_0(t_0)=0$ to $z_0(t_1)=d$, the total E -field energy $\mathcal{E}(t)$ of the system, aside from the additive constant, drops from $+\frac{1}{2}qE_0d$ to $-\frac{1}{2}qE_0d$, for a total decline of qE_0d . This, of course, is precisely the kinetic energy gained by the particle when it moves from the bottom plate to the top plate, as evaluated in part (b).
