

Problem 1)

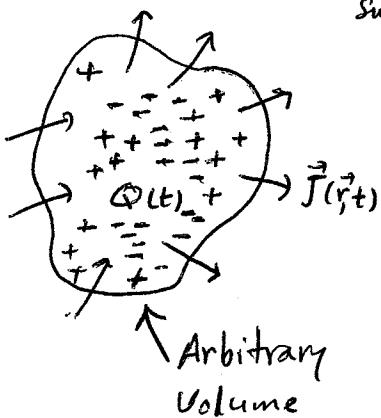
$$a) \vec{J}_{\text{free}}(\vec{r}, t) = \rho_{\text{free}}(\vec{r}, t) \vec{V}(\vec{r}, t) \quad \checkmark$$

$$b) \text{units of } \rho = \text{Coulomb/m}^3; \text{ units of } \vec{V} = \text{m/s} \Rightarrow \text{units of } \vec{J} = \frac{\text{Coulomb}}{\text{m}^3} \cdot \frac{\text{m}}{\text{s}} = \frac{\text{Amp.}}{\text{m}^2} \quad \checkmark$$

$$c) \nabla \cdot \vec{J}_{\text{free}}(\vec{r}, t) + \frac{\partial}{\partial t} \rho_{\text{free}}(\vec{r}, t) = 0 \quad \leftarrow \text{differential form} \quad \checkmark$$

$$\text{Apply Gauss' Theorem} \Rightarrow \int_{\text{Volume}} \nabla \cdot \vec{J}_{\text{free}}(\vec{r}, t) d\vec{r} + \int_{\text{Surface}} \frac{\partial}{\partial t} \rho_{\text{free}}(\vec{r}, t) d\vec{s} = 0$$

$$\Rightarrow \oint_{\text{Surface}} \vec{J}_{\text{free}}(\vec{r}, t) \cdot d\vec{s} + \frac{\partial}{\partial t} Q(t) = 0 \quad \leftarrow \text{Integral form} \quad \checkmark$$



The integral of $\vec{J}_{\text{free}}(\vec{r}, t)$ over the closed surface of the volume of interest yields the net flux of charge out of the volume. (outward flux is because $d\vec{s}$ is defined to point from the inside of the volume to the outside.)

Therefore, during a short time interval Δt , the net charge leaving the volume is equal to $\Delta t \oint_{\text{Surface}} \vec{J}_{\text{free}}(\vec{r}, t) \cdot d\vec{s}$. This must be equal to the net

drop in the total charge $Q(t)$ contained within the volume. Therefore:

$$\Delta t \oint_{\text{Surface}} \vec{J}_{\text{free}}(\vec{r}, t) \cdot d\vec{s} = -\Delta Q \Rightarrow \oint_{\text{Surface}} \vec{J}_{\text{free}}(\vec{r}, t) \cdot d\vec{s} + \frac{\Delta Q}{\Delta t} = 0 \quad \checkmark$$

$$d) \nabla \cdot \vec{J}_{\text{free}}(\vec{r}, t) + \frac{\partial}{\partial t} \rho_{\text{free}}(\vec{r}, t) = 0$$

Fourier transforming both sides of the above equation yields:

$$\int_{-\infty}^{\infty} \nabla \cdot \vec{J}_{\text{free}}(\vec{r}, t) e^{-i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{r} dt + \int_{-\infty}^{\infty} \frac{\partial}{\partial t} \rho_{\text{free}}(\vec{r}, t) e^{-i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{r} dt = 0$$

Using integration by parts, each term in the above equation is simplified as follows:

$$\iiint_{-\infty}^{\infty} \frac{\partial}{\partial x} J_x(\vec{r}, t) e^{-i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{r} dt = \iiint_{-\infty}^{\infty} \left[J_x(\vec{r}, t) e^{-i(\vec{k} \cdot \vec{r} - \omega t)} \right] dy dz dt$$

$\curvearrowright \rightarrow 0 \text{ as } x \rightarrow \pm\infty$

$$- \iiint_{-\infty}^{\infty} J_x(\vec{r}, t) \frac{\partial}{\partial x} e^{-i(\vec{k} \cdot \vec{r} - \omega t)} dx dy dz dt = 0 + i k_x \iiint_{-\infty}^{\infty} J_x(\vec{r}, t) e^{-i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{r} dt$$

$$= i k_x J_x(\vec{k}, \omega) \quad \checkmark$$

Similarly, one can show that

$$\int_{-\infty}^{\infty} \frac{\partial}{\partial y} J_y(\vec{r}, t) e^{-i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{r} dt = i k_y J_y(\vec{k}, \omega) \quad \checkmark$$

$$\int_{-\infty}^{\infty} \frac{\partial}{\partial z} J_z(\vec{r}, t) e^{-i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{r} dt = i k_z J_z(\vec{k}, \omega) \quad \checkmark$$

$$\int_{-\infty}^{\infty} \frac{\partial}{\partial t} \rho(\vec{r}, t) e^{-i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{r} dt = -i\omega \rho(\vec{k}, \omega) \quad \checkmark$$

$$\text{Therefore, } i k_x J_x(\vec{k}, \omega) + i k_y J_y(\vec{k}, \omega) + i k_z J_z(\vec{k}, \omega) - i\omega \rho(\vec{k}, \omega) = 0 \Rightarrow$$

$\checkmark \quad \boxed{\vec{k} \cdot \vec{J}(\vec{k}, \omega) = \omega \rho_{\text{free}}(\vec{k}, \omega)}$ ← Continuity equation in the Fourier domain.

Alternative method for Part(d):

$$\vec{D} \cdot \vec{J}_{\text{free}}(\vec{r}, t) + \frac{\partial}{\partial t} \rho_{\text{free}}(\vec{r}, t) = 0 \Rightarrow \vec{D} \cdot \left\{ \frac{1}{(2\pi)^4} \int_{-\infty}^{\infty} \vec{J}_{\text{free}}(\vec{k}, \omega) e^{i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{k} d\omega \right\} +$$

$$\frac{\partial}{\partial t} \left\{ \frac{1}{(2\pi)^4} \int_{-\infty}^{\infty} \rho_{\text{free}}(\vec{k}, \omega) e^{+i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{k} d\omega \right\} = 0 \Rightarrow$$

$$\int_{-\infty}^{\infty} J_x(\vec{k}, \omega) \frac{\partial}{\partial x} e^{i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{k} d\omega + \int_{-\infty}^{\infty} J_y(\vec{k}, \omega) \frac{\partial}{\partial y} e^{i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{k} d\omega + \int_{-\infty}^{\infty} J_z(\vec{k}, \omega) \frac{\partial}{\partial z} e^{i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{k} d\omega + \int_{-\infty}^{\infty} \rho_{\text{free}}(\vec{k}, \omega) \frac{\partial}{\partial t} e^{i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{k} d\omega = 0 \Rightarrow$$

$$+ \int_{-\infty}^{\infty} \rho_{\text{free}}(\vec{k}, \omega) \frac{\partial}{\partial t} e^{i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{k} d\omega = 0 \Rightarrow$$

$$\int_{-\infty}^{\infty} (k_x J_x + k_y J_y + k_z J_z) e^{i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{k} d\omega - i \int_{-\infty}^{\infty} \rho_{\text{free}}(\vec{k}, \omega) e^{i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{k} d\omega = 0 \Rightarrow$$

$$\int_{-\infty}^{\infty} [\vec{k} \cdot \vec{J}_{\text{free}}(\vec{k}, \omega) - \omega \rho_{\text{free}}(\vec{k}, \omega)] e^{i(\vec{k} \cdot \vec{r} - \omega t)} d\vec{k} d\omega = 0 \Rightarrow \boxed{\vec{k} \cdot \vec{J}_{\text{free}}(\vec{k}, \omega) - \omega \rho_{\text{free}}(\vec{k}, \omega) = 0}$$