

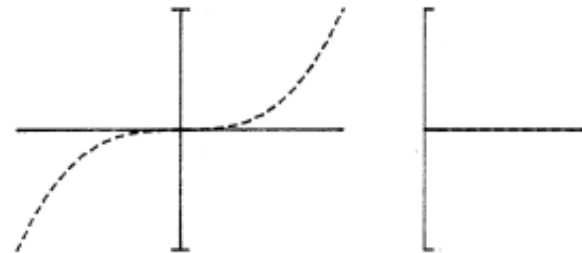
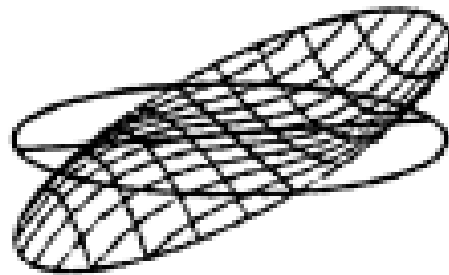
Coma aberration

Lens Design OPTI 517

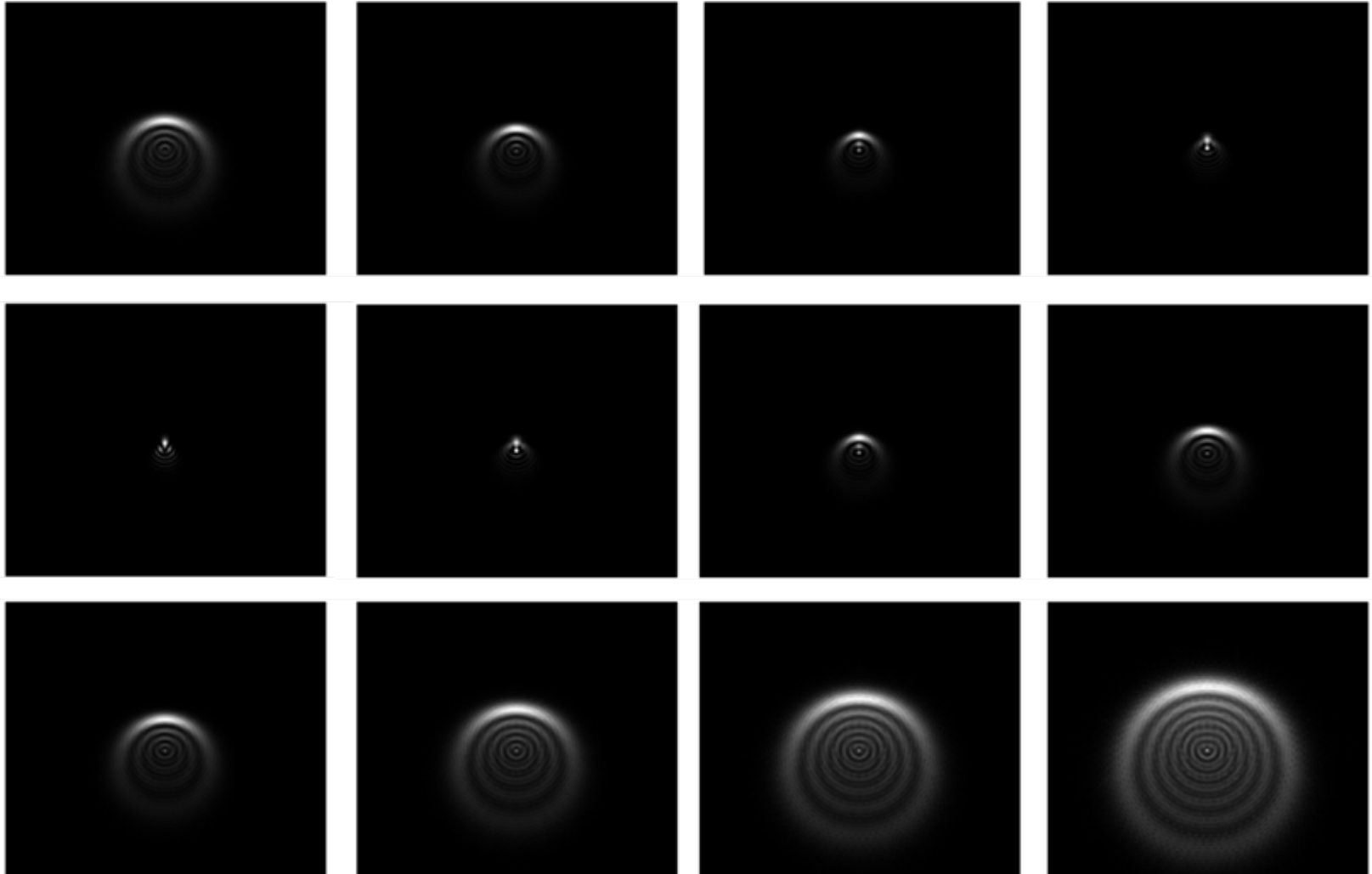
Coma

0.25 wave	1.0 wave	2.0 waves	4.0 waves	Spot diagram

$$\begin{aligned}
 W(H, \rho, \theta) = & W_{200}H^2 + W_{020}\rho^2 + W_{111}H\rho \cos \theta + \\
 & + W_{040}\rho^4 + W_{131}H\rho^3 \cos \theta + W_{222}H^2\rho^2 \cos^2 \theta + \\
 & + W_{220}H^2\rho^2 + W_{311}H^3\rho \cos \theta + W_{400}H^4 + \\
 & + \dots
 \end{aligned}$$



Coma though focus

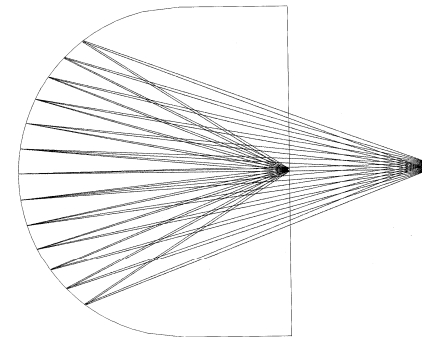
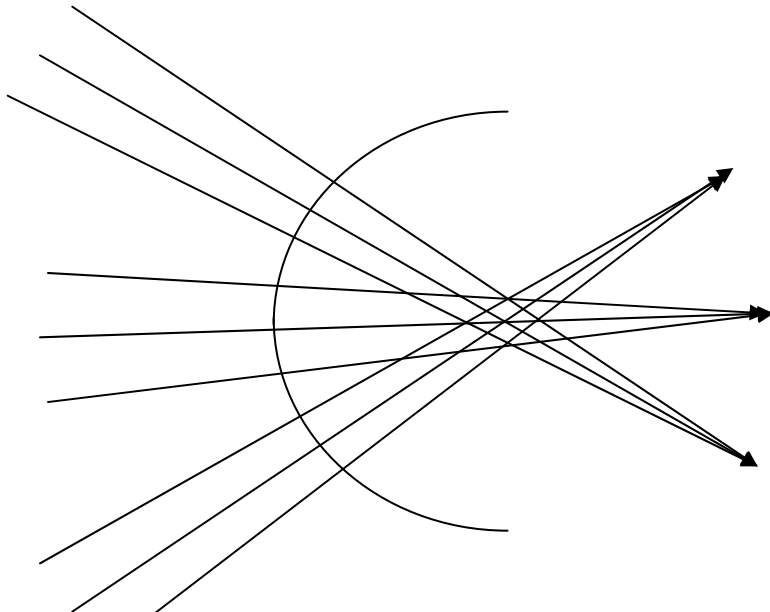
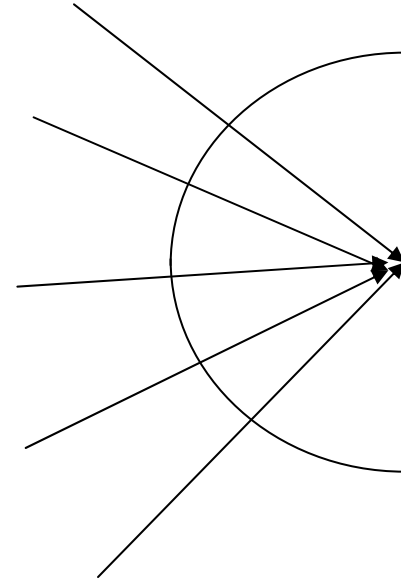
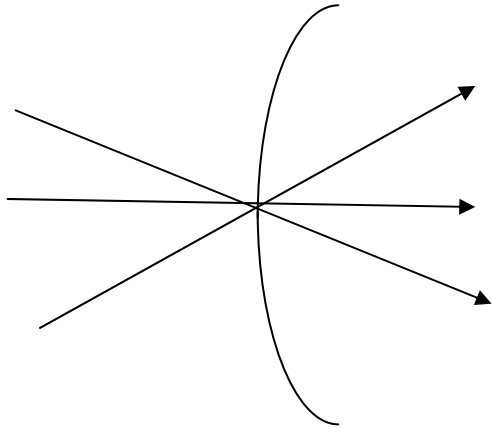


Cases of zero coma

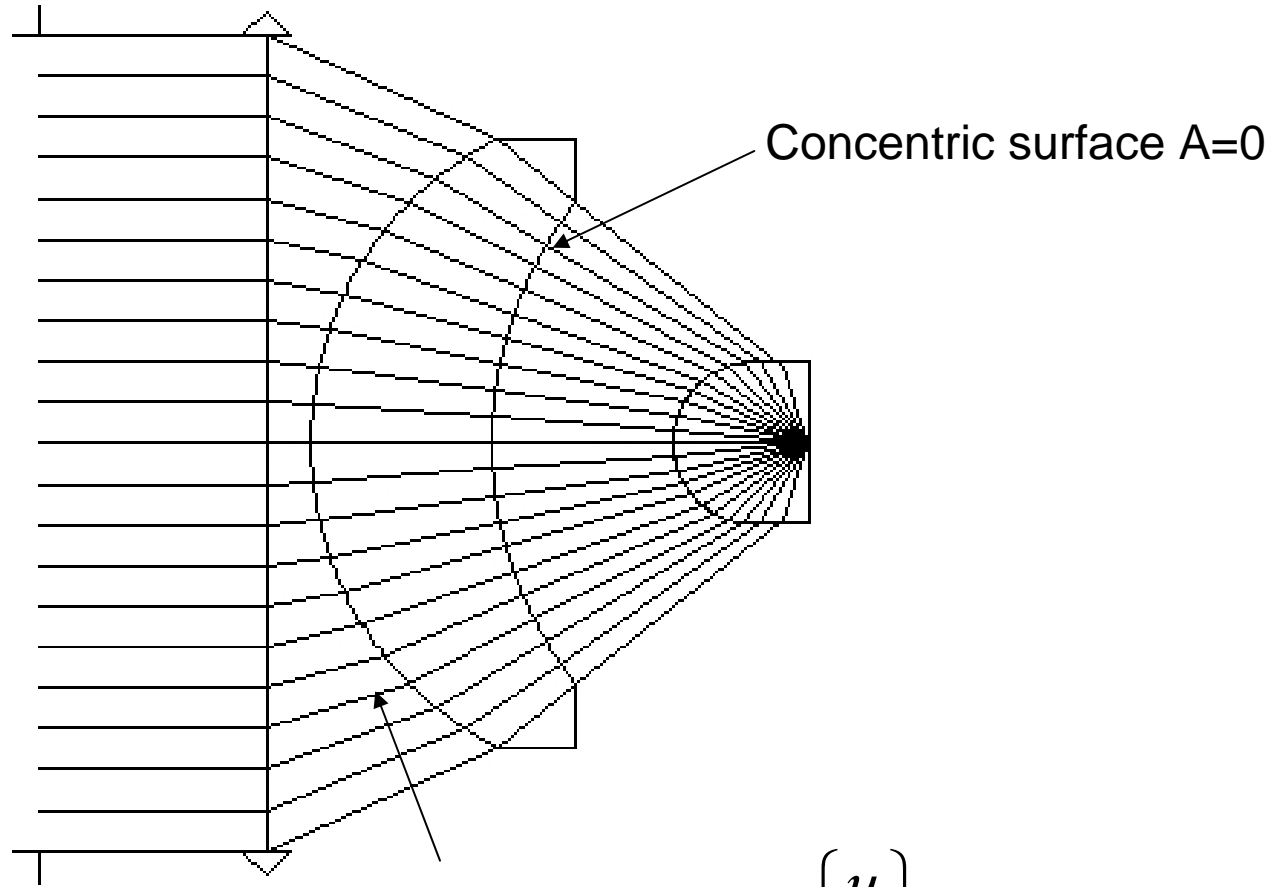
$$W_{131} = -\frac{1}{2} A \bar{A} \Delta \left\{ \frac{u}{n} \right\} y$$

- At $y=0$, surface is at an image
- $A=0$, On axis beam concentric with center of curvature
- $\bar{A}=0$, Off-axis beam concentric, chief ray goes through the center of curvature
- Aplanatic points

Cases of zero coma



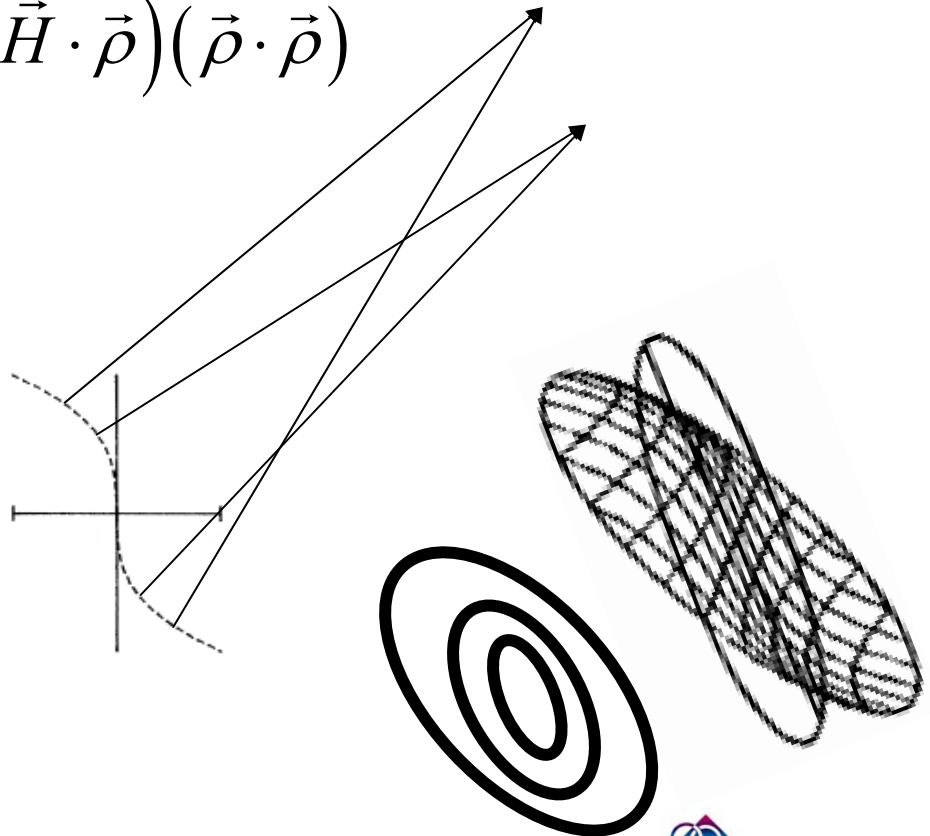
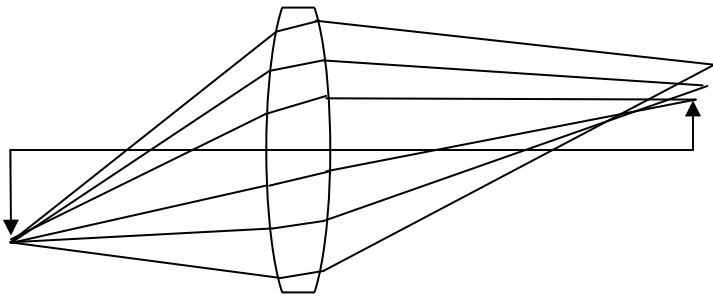
Aplanatic-concentric



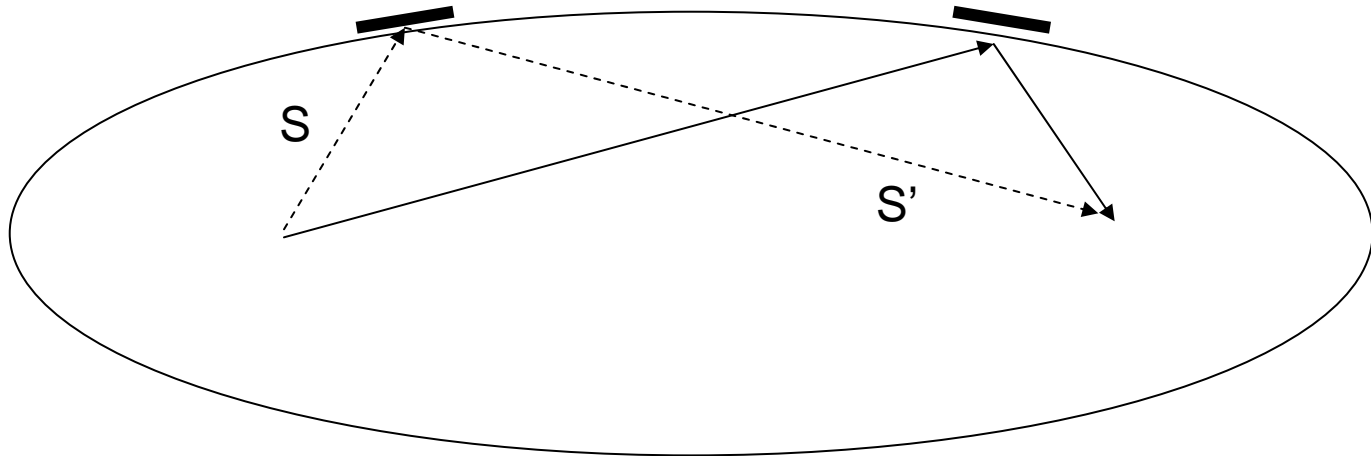
Aplanatic surface $\Delta \left\{ \frac{u}{n} \right\} = 0$

Coma as a variation of magnification with aperture

$$W_{131}(\vec{H}, \vec{\rho}) = -\frac{1}{2} A \bar{A} \Delta \left\{ \frac{u}{n} \right\} y(\vec{H} \cdot \vec{\rho})(\vec{\rho} \cdot \vec{\rho})$$



Coma as a variation of magnification with aperture II

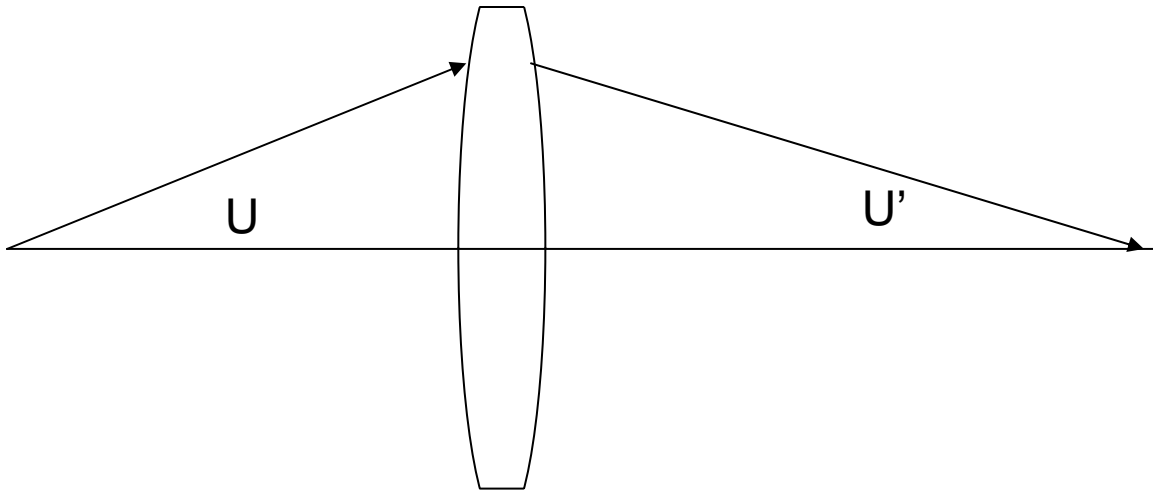


$$m = s' / s$$

Sine condition

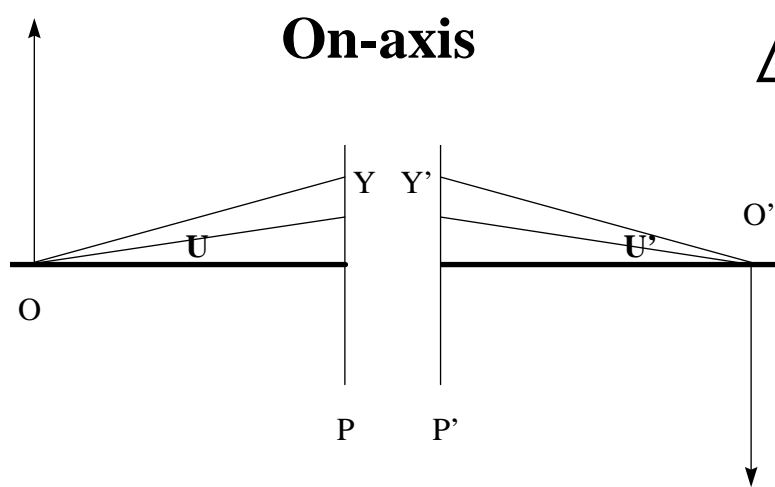
Coma aberration can be considered as a variation of magnification with respect to the aperture. If the paraxial magnification is equal to the real ray marginal magnification, then an optical system would be free of coma.

Spherical aberration can be considered as a variation of the focal length with the aperture.



$$\frac{u}{u'} = \frac{\sin(U)}{\sin(U')}$$

Sine condition



Optical path length between O and O' is L_{axis} and does not depend on Y or Y'

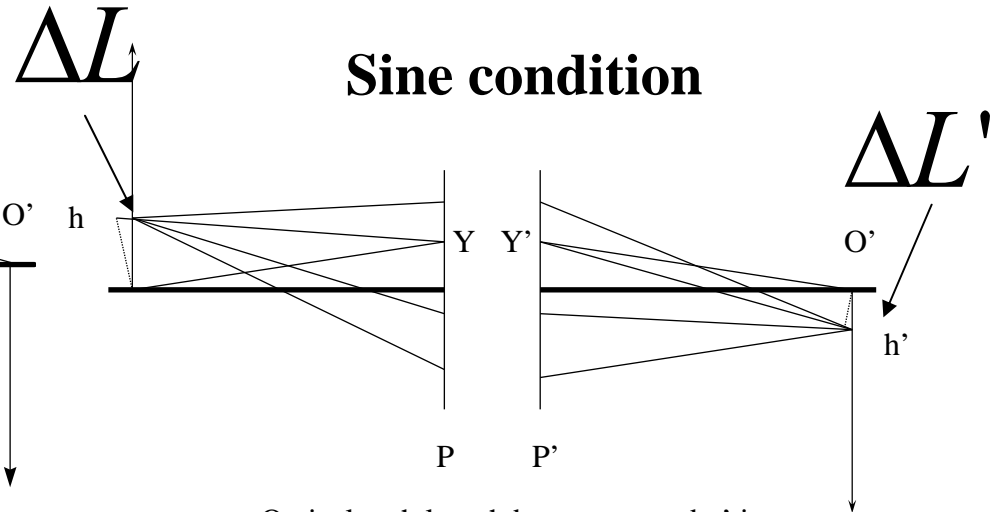
$$\Delta L = h \sin(U)$$

$$\Delta L' = h' \sin(U'),$$

$$h' n' \sin(U') = h n \sin(U)$$

That is: OPD has no linear phase errors as a function of field of view!

Prof. Jose Sasian



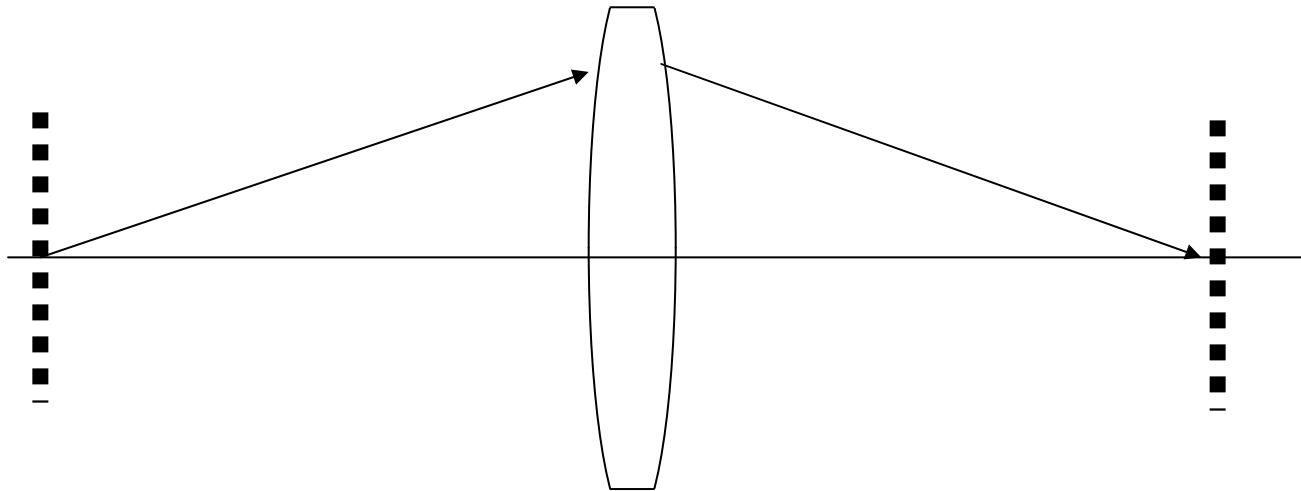
Optical path length between y and y' is

$$L_{off-axis} = L_{axis} + \Delta L' - \Delta L$$

$$L_{off-axis} = L_{axis} + h' n' \sin(U') - h n \sin(U)$$

$$\frac{u}{u'} = \frac{\sin(U)}{\sin(U')}$$

Imaging a grating

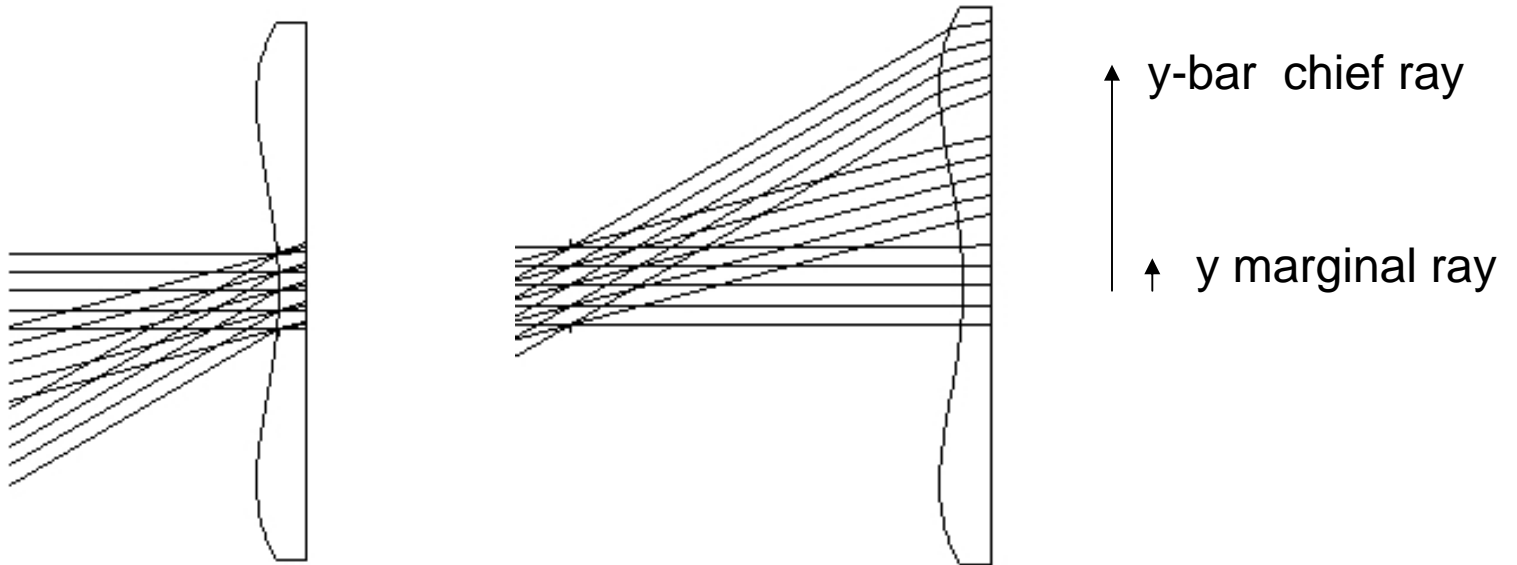


$$\sin(U) = \frac{m \cdot \lambda}{d}$$

$$d \cdot \sin(U) = m \cdot \lambda = d' \sin(U')$$

Contribution from an aspheric surface

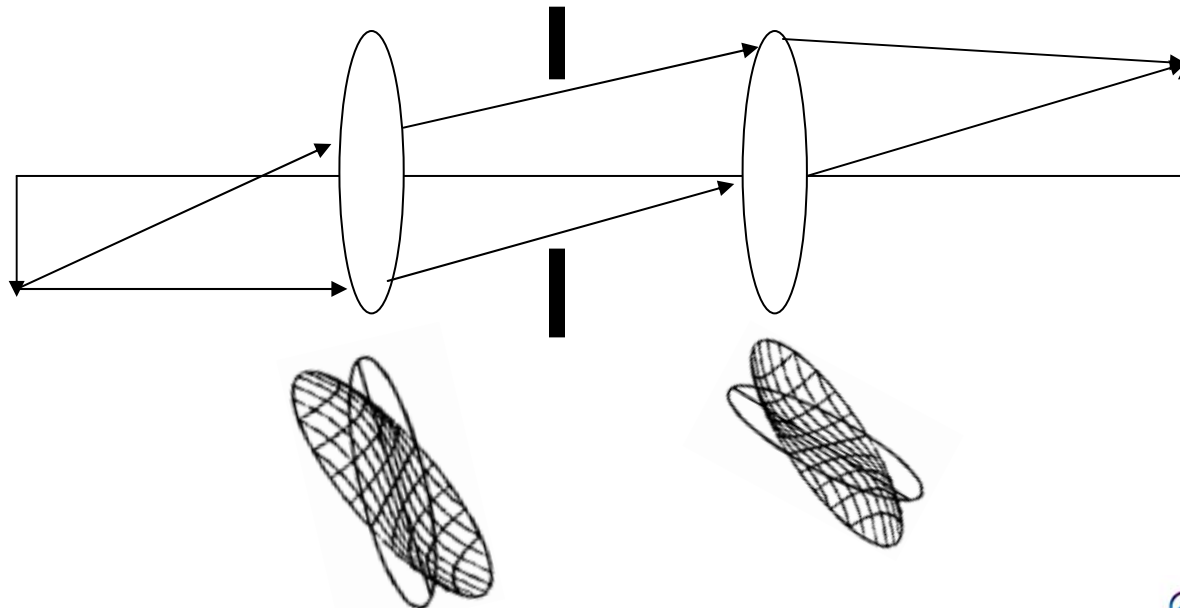
$$\Delta W_{131} = -\frac{1}{2} \frac{\bar{y}}{y} \left\{ 8 A_4 y^4 \Delta(n) \right\}$$



Aberrations and symmetry

$$\begin{aligned} W(H, \rho, \theta) = & W_{200}H^2 + W_{020}\rho^2 + W_{111}H\rho \cos \theta + \\ & + W_{040}\rho^4 + W_{131}H\rho^3 \cos \theta + W_{222}H^2\rho^2 \cos^2 \theta + \\ & + W_{220}H^2\rho^2 + W_{311}H^3\rho \cos \theta + W_{400}H^4 + \\ & + \dots \end{aligned}$$

- Coma is an odd aberration with respect to the stop
- Natural stop position to cancel coma by symmetry



Structural coefficients: Thin lens (stop at lens)

$$S_I = \frac{1}{4} y^4 \phi^3 [AX^2 - BXY + CY^2 + D]$$

$$A = \frac{n+2}{n(n-1)^2}$$

$$S_{II} = \frac{1}{2} \mathcal{K} y^2 \phi^2 [EX - FY]$$

$$B = \frac{4(n+1)}{n(n-1)}$$

$$S_{III} = \mathcal{K}^2 \phi$$

$$S_{IV} = \mathcal{K}^2 \phi \frac{1}{n}$$

$$X = \frac{c_1 + c_2}{c_1 - c_2} = \frac{r_2 + r_1}{r_2 - r_1}$$

$$C = \frac{3n+2}{n}$$

$$S_V = 0$$

$$Y = \frac{1+m}{1-m} = \frac{u'+u}{u'-u}$$

$$D = \frac{n^2}{(n-1)^2}$$

$$C_L = y^2 \phi \frac{1}{v}$$

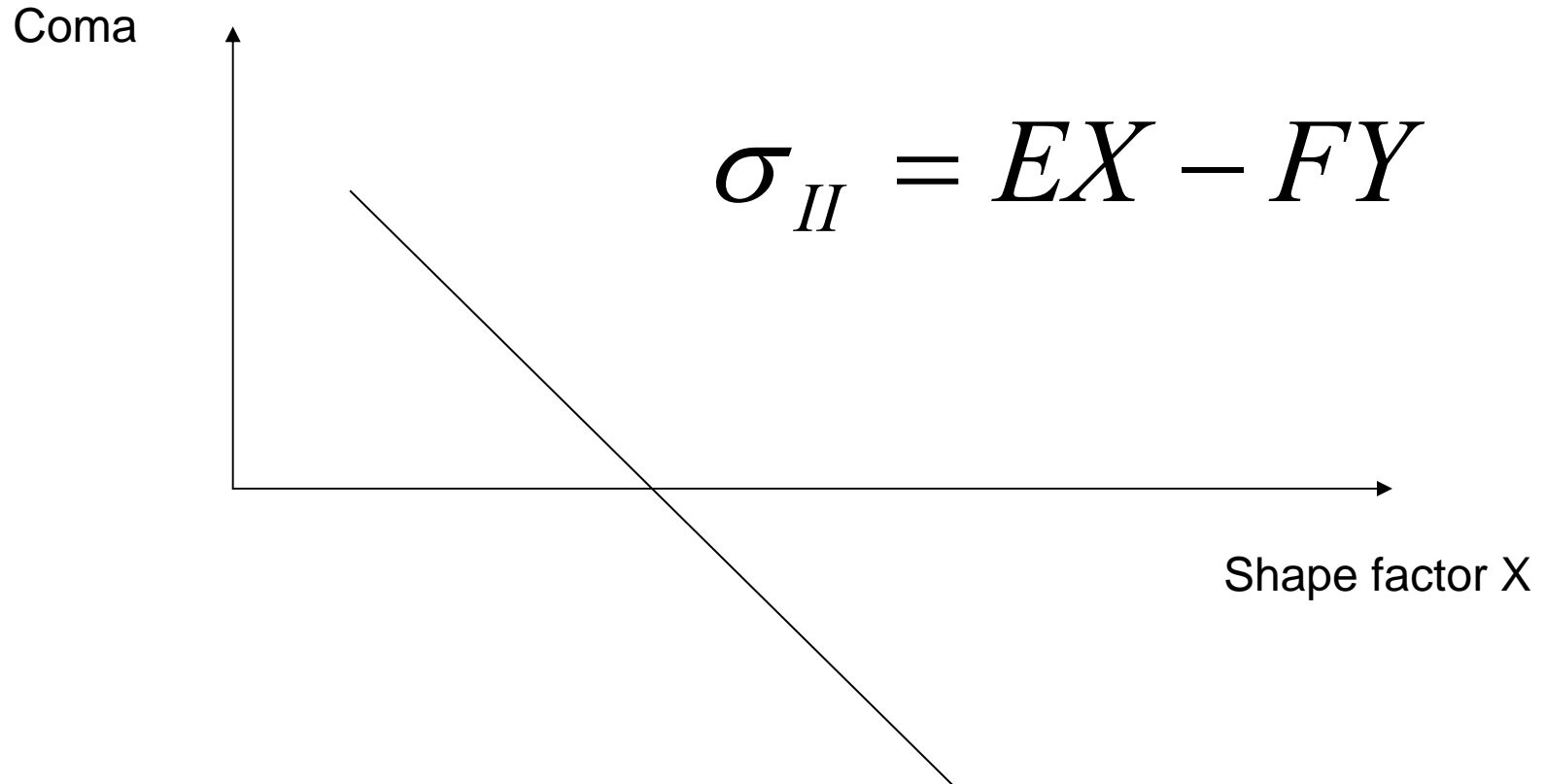
$$E = \frac{n+1}{n(n-1)}$$

$$C_T = 0$$

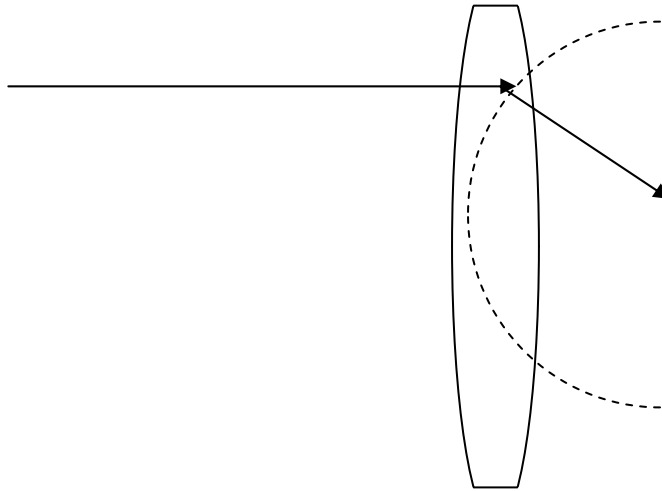
$$\phi = \Delta n \Delta c = (n-1)(c_1 - c_x)$$

$$F = \frac{2n+1}{n}$$

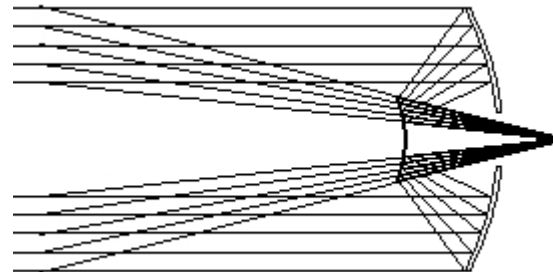
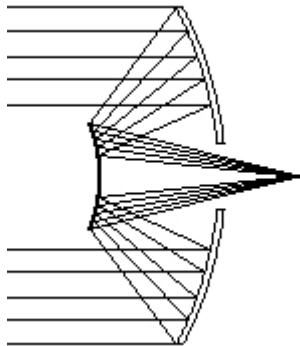
Coma vs Bending



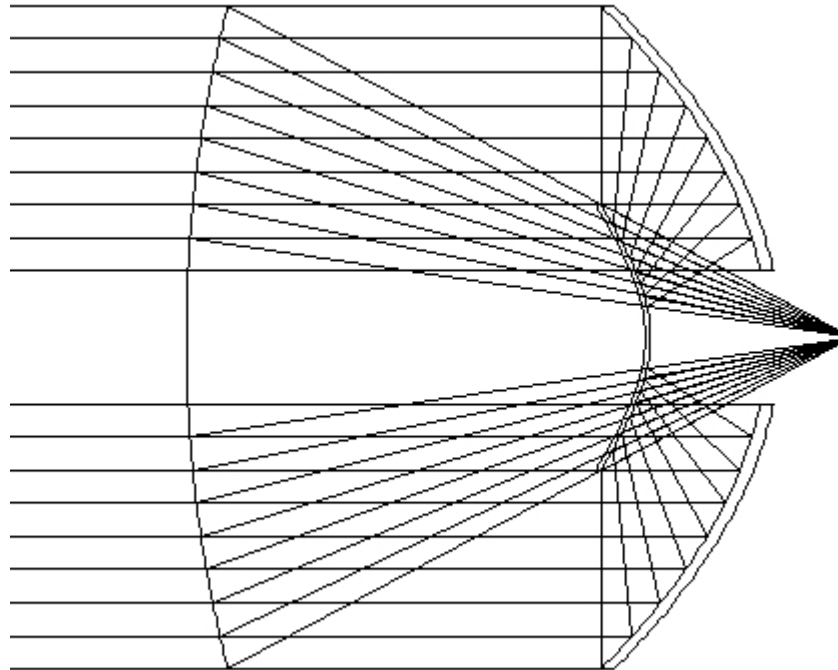
Principal surface



In an aplanat working at $m=0$
the equivalent
refracting surface
is a hemisphere

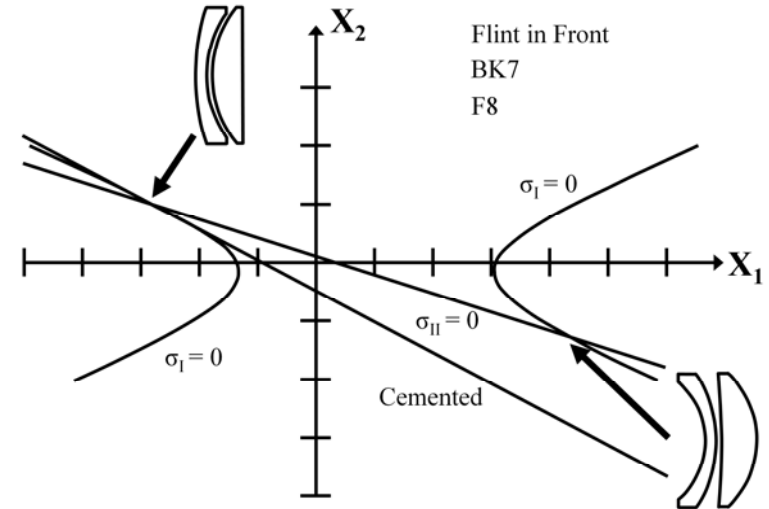
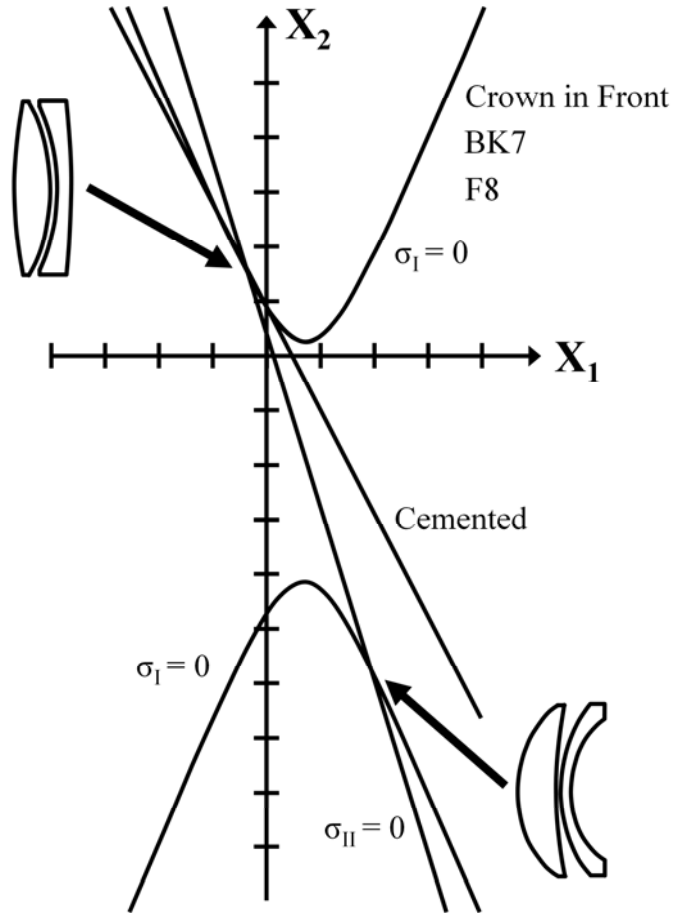


Cassegrain's principal surface



Since the equivalent refracting surface in a Cassegrain telescope is a paraboloid then the coma of that Cassegrain is the same of a paraboloid mirror with the same focal length.

Aplanat doublets

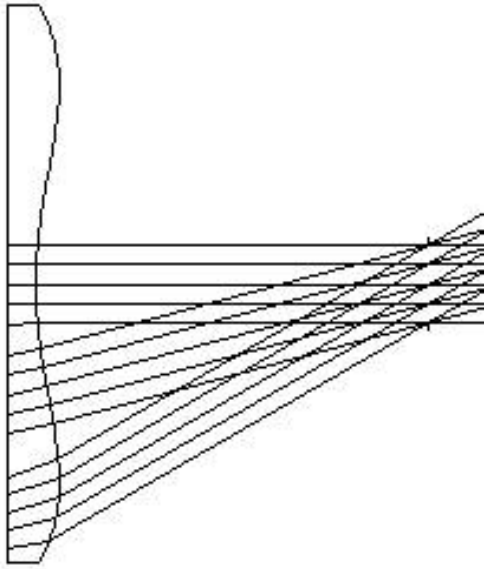


Kingslake's cemented aplanat

Chromatic correction
Spherical aberration correction
Coma correction
Still cemented

$C=1/r$	d	Glass	n_D	V_D
0.1509				
	0.32	SK-11	1.56376	60.75
-0.2246				
	0.15	SF-19	1.66662	33.08
-.052351				

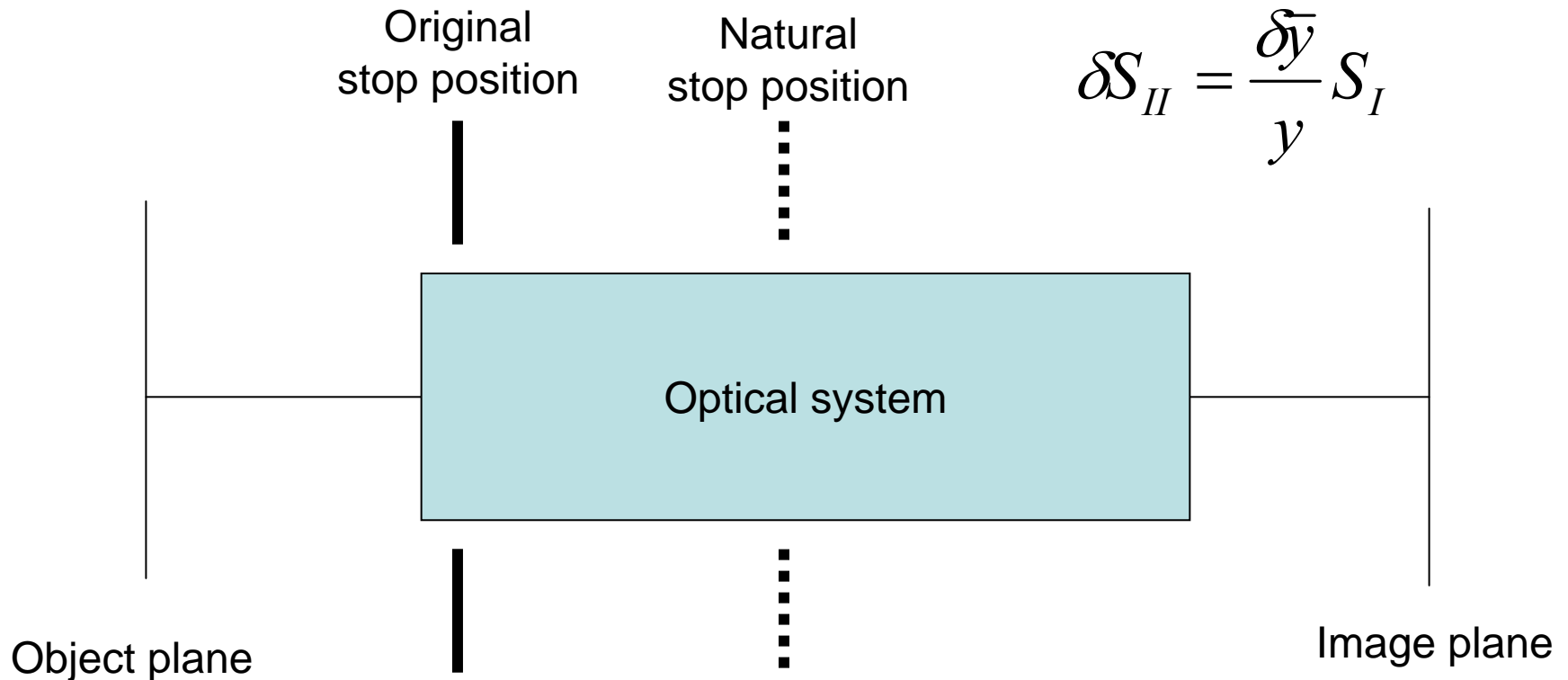
Control of coma in the presence of an aspheric mirror near a pupil



In the presence of a strong aspheric surface near the stop or pupil, coma aberration can be corrected by moving the surface

$$\Delta W_{131} = -\frac{1}{2} \frac{\bar{y}}{y} \left\{ 8A_4 y^4 \Delta(n) \right\}$$

Coma correction by phantom stop position



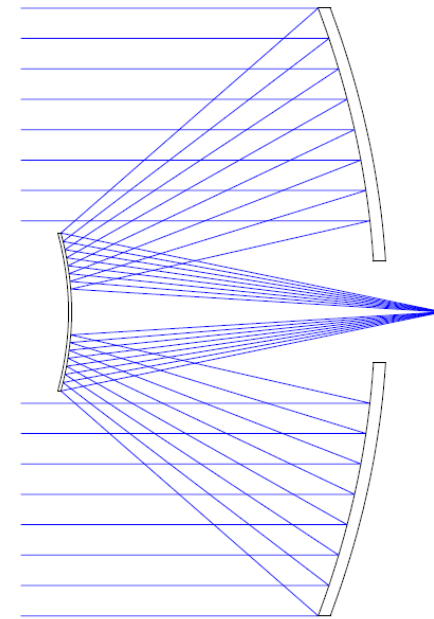
In the presence of spherical aberration there is a stop position for which coma is zero. At that stop position spherical aberration might be corrected. Then the system becomes aplanatic and the stop can be shifted back to its original position.

The aplanatic member(s) in a family of solutions

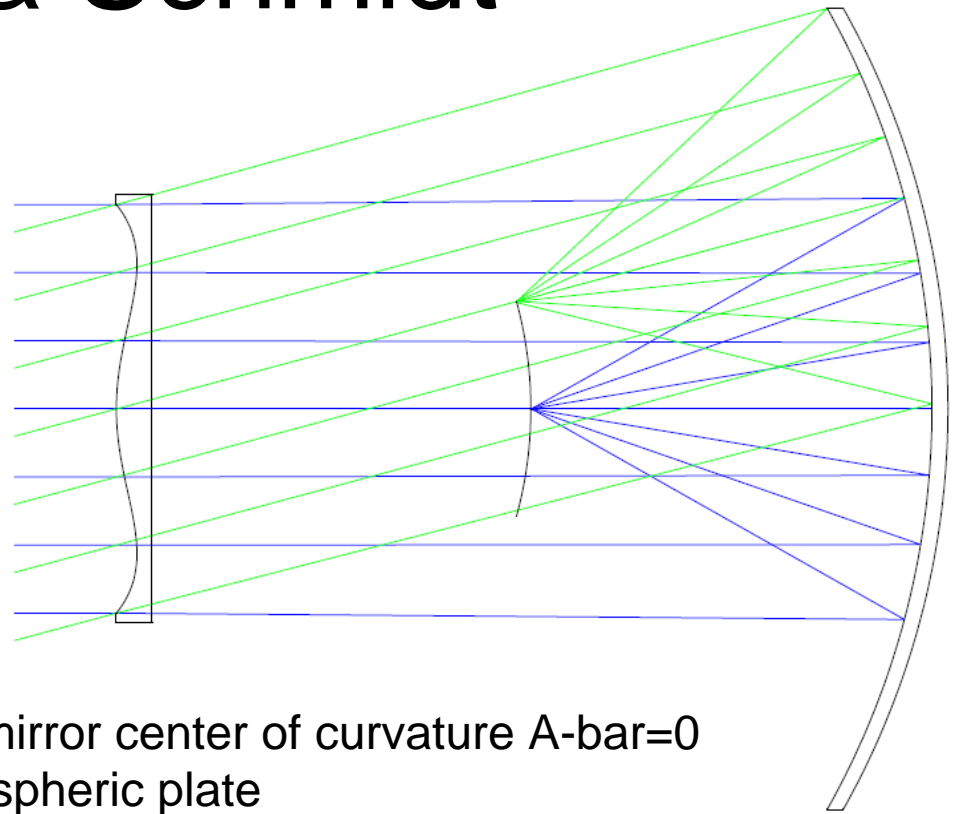
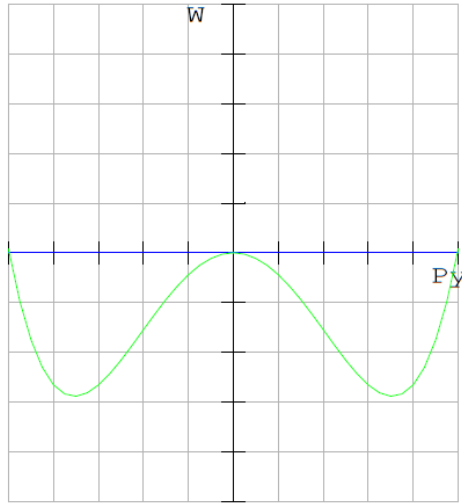
Doublet



Ritchey-Chretien



Camera Schmidt



Aspheric plate at mirror center of curvature $\bar{A}=0$

Stop aperture at aspheric plate

Note symmetry about mirror CC

No spherical aberration

No coma

No astigmatism.

Anastigmatic over a wide field of view!

Satisfies Conrady's D-d sum

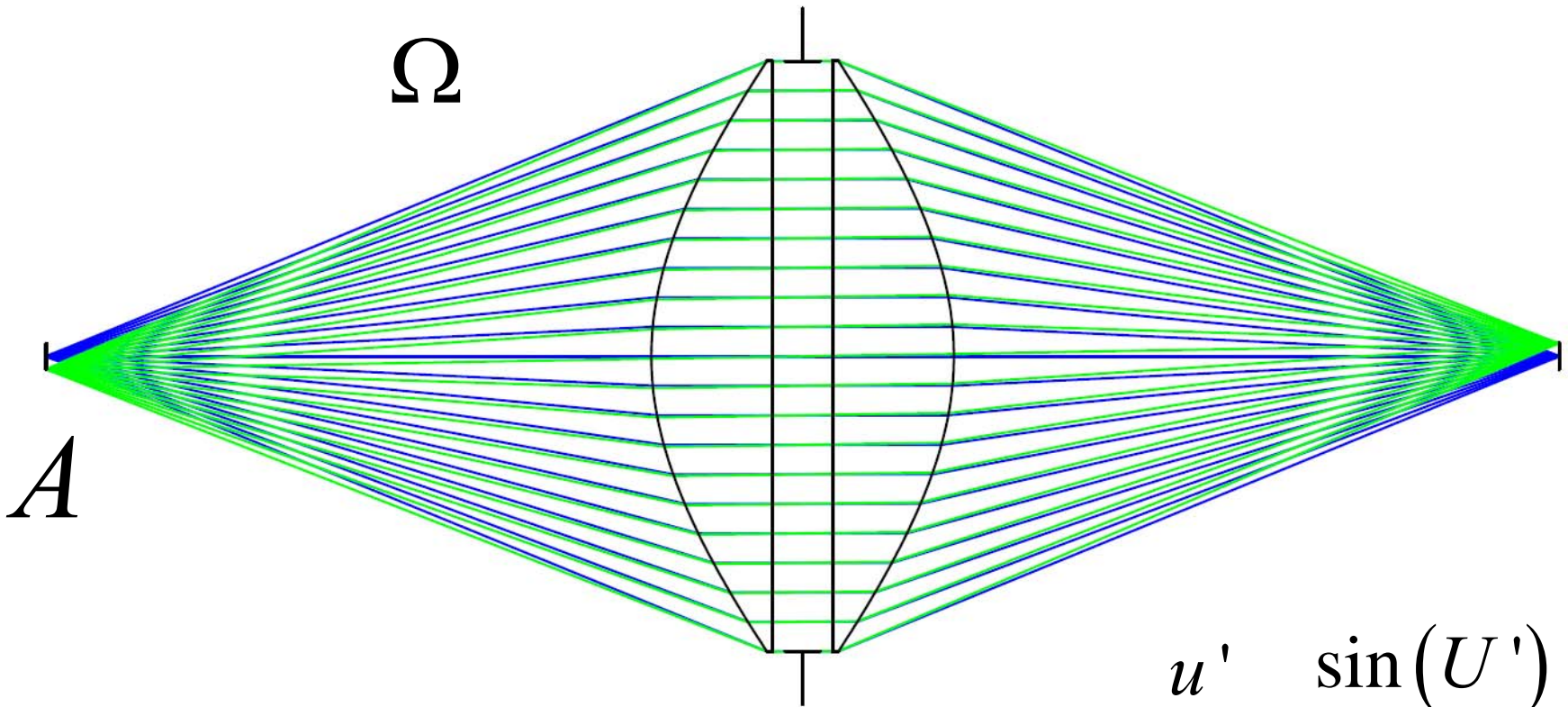
Notes

- Need to make aplanatic zero-field systems (that are fast). The alignment becomes easier.
- Lenses for lasers diodes/optical fibers
- Microscope objectives

Sine condition from optical flux conservation and radiance theorem

Etendu considerations

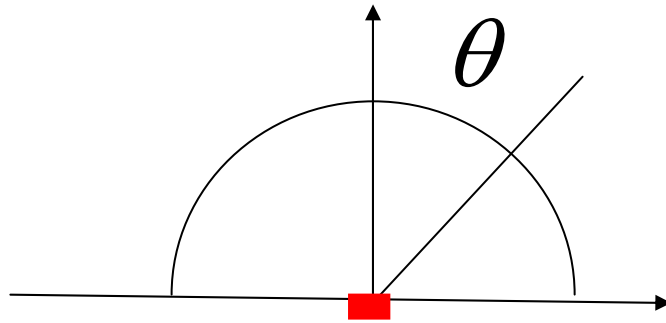
Sine condition



$$\frac{u'}{u} = \frac{\sin(U')}{\sin(U)}$$

Optical flux from a Lambertian source

$$\Phi(\theta) = 2\pi AL_0 \int_0^\theta \cos(\theta) \sin(\theta) d\theta = \pi AL_0 \sin^2(\theta)$$



Compare with homogeneous source

$$\Phi(\theta) = 2\pi AL_0 \int_0^\theta \sin(\theta) d\theta = 2\pi AL_0 (1 - \cos(\theta)) = AL_0 \Omega$$

Optical flux=radiance x throughput

$$\Phi(\theta) = \frac{L_0}{n^2} T \quad \Phi(\theta') = \frac{L_0'}{n'^2} T$$

$$\frac{n^2}{L_0} \Phi(\theta) = \frac{n'^2}{L_0'} \Phi(\theta') = T$$

$$\frac{n^2}{L_0} = \frac{n'^2}{L_0'} \quad \text{Radiance theorem}$$

$$U = \theta$$

$$\frac{n^2}{L_0} \Phi(U) = \frac{n'^2}{L_0'} \Phi(U')$$

Sine condition from optical flux conservation

$$\frac{n^2}{L_0} \Phi(U) = \frac{n'^2}{L_0'} \Phi(U')$$

$$n^2 A \sin^2(U) = n'^2 A' \sin^2(U')$$

$$\pi n^2 h^2 \sin^2(U) = \pi h'^2 n'^2 \sin^2(U')$$

$$n^2 h^2 \sin^2(U) = h'^2 n'^2 \sin^2(U')$$

$$nh \sin(U) = n' h' \sin(U')$$

$$\frac{u'}{u} = \frac{\sin(U')}{\sin(U)}$$

Sine condition

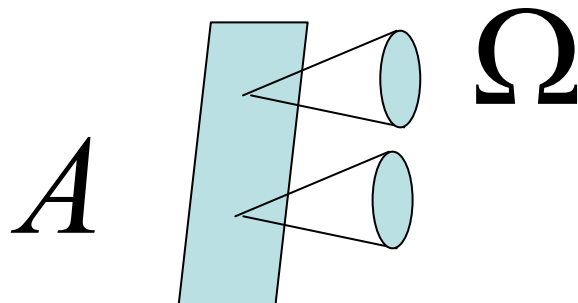
Throughput Etendu

(area-omega product)

$$\mathcal{E} = n^2 A \Omega = n'^2 A' \Omega'$$

$$T = \pi \mathcal{K}^2$$

“Capacity to transfer optical flux”



$$\Omega = 2\pi (1 - \cos(\theta))$$

Even better

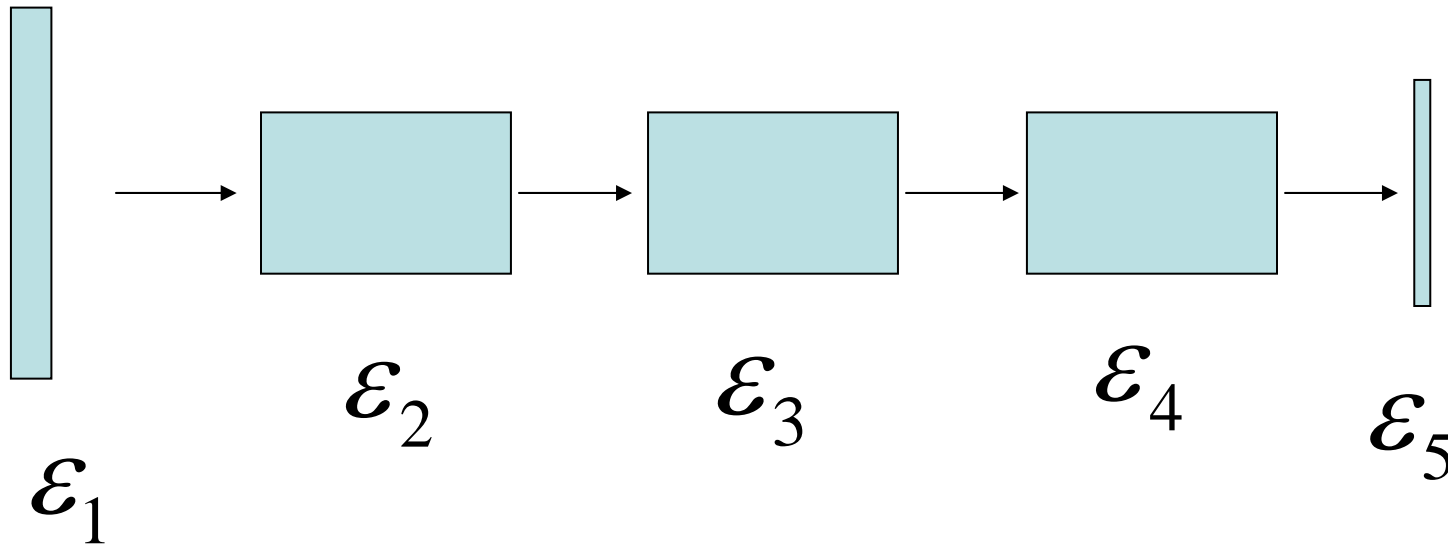
$$\varepsilon = n^2 A \Omega = n'^2 A' \Omega'$$

Homogeneous
source

$$\varepsilon = A (NA)^2 = n^2 A \sin^2(\theta)$$

Lambertian
source

Etendu considerations are key to design an optical system



$$\epsilon_5 \geq \epsilon_4 \geq \epsilon_3 \geq \epsilon_2 \geq \epsilon_1$$

Start with the sensor at the end $\epsilon_5 = A_5 (NA)^2$

Summary

- Coma aberration
- Coma as an odd aberration
- Sine condition
- Natural stop position
- Aplanatic doublets
- Zero-field systems